Atomic dark matter [Self-interacting and asymmetric dark matter]

Kallia Petraki Nikhef, Amsterdam



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Outline

- Challenges of collision-less cold DM
- Self-interacting DM
- Self-interacting \cap Asymmetric DM
- Asymmetric DM
- Case study: atomic dark matter

Collisionless ACDM and galactic structure

- Very successful in explaining large-scale structure.
- At galactic and subgalactic scales: discrepancies between simulations and observations.
 - Cusps vs cores: predicted galactic density profiles too steep; observations better fit with cores.
 - Missing satellites: Too many subhaloes predicted for a galaxy of the size of the Milky Way.
 - "Too big to fail"

[Boylan-Kolchin, Bullock, Kaplinghat (2011)]

The largest subhaloes formed in simulations are larger than the brightest satellites, i.e. they have no observed counterpart.

This means that star formation must have failed where it is expected to be very successful.

Collisionless ACDM and galactic structure

Summary: collisionless ACDM predicts too rich structure at small scales

too much matter in central few kpc of typical galaxies

[an overview: Weinberg, Bullock, Governato, Kuzio de Naray, Peter; arXiv: 1306.0913]

Small-scale galactic structure: possible resolutions

- Baryonic physics
- Observations
- Shift in the DM paradigm
 - Warm DM, e.g. keV sterile neutrinos
 - Self-interacting DM

Small-scale galactic structure: shift in the DM paradigm

• At large scales:

no difference from Λ CDM; retain its successes.

• At smaller scales:

suppress structure, in agreement with observations

 \rightarrow how warm, or how self-interacting can DM be?

Self-interacting DM

The energy & momentum exchange between DM particles:

- Heats up the low-entropy material
 - → suppresses overdensities [cusps vs cores]
 - → suppresses star-formation rate [missing satellites, "too big to fail"]
- Isotropises DM halos
 - \rightarrow constrained by observed ellipticity of large haloes.



[Theory: Spergel, Steinhardt (2000); Simulations: Rocha et al. (2012); Peter et al. (2012); Zavala et al (2012)]

Self-interacting DM



- Short-range interactions: σ_{scatt} is velocity independent; available range seems limited.
- Long-range interactions: σ_{scatt} ~ 1 / vⁿ, n > 0; much larger range of possibilities
 - significant effect on small haloes (small velocity dispersion)
 - negligible effect on large haloes (large velocity dispersion)
- → Consider DM self-interactions mediated by light particles.

[Zavala, Vogelsberger, Walker (2012); Zavala, Vogelsberger (2012)]



[Zavala, Vogelsberger, Walker (2012); Zavala, Vogelsberger (2012)]

Density profiles



[Zavala, Vogelsberger, Walker (2012); Zavala, Vogelsberger (2012)]



Bottom right: can be fitted by lower mass subhaloes.

[Zavala, Vogelsberger, Walker (2012); Zavala, Vogelsberger (2012)]



Self-interaction cross-section

Conclusion

Constant cross-section:

some tension between ellipticity constraints and cross-section required to change subhalo kinematics. Nevertheless

 $\sigma_{\text{scatt}} / m_{\text{DM}} \sim 1 \text{ barn } / \text{GeV}$ could work (narrow range).

Velocity-dep cross-section:

satisfies ellipticity constraints and fits subhalo kinematics.

Self-interacting DM

Self-interaction



χ : dark matter **φ** : mediator m_φ << m_χ

Annihilation



- Sizeable self-interactions via light mediators imply minimum contribution to DM annihilation; annihilation cross-section could exceed canonical value for symmetric thermal relic DM
 - → consider **asymmetric DM**

Thermal relic DM



[Review of asymmetric dark matter; KP, Volkas (2013)]

(a little simplified) Venn diagram of stable / long-lived relics



Asymmetric dark matter

- provides a suitable host for DM self-interacting via light species.
- encompasses most of the low-energy parameter space of thermal relic DM \rightarrow study models and low-energy pheno.

a cosmic coincidence



- Unrelated mechanisms → different parameters
 → result expected to differ by orders of magnitude.
- Similarity of abundances hints towards related physics for VM and DM production.

Ordinary matter Stable particles: p e y v

- *p*⁺ make up most of ordinary matter in the universe
- Only p⁺, no p⁻ present today: matter-antimatter asymmetry
 - > observational evidence: negligible antimatter in cosmic rays
 - > theoretical consistency: $p^+ p^-$ annihilation cross-section too large
 - ⇒ they destroy each other too efficiently,
 - \Rightarrow in an expanding universe, very few $p^+ \& p^-$ left over
 - \Rightarrow deficit of antiparticles stops annihilations, excess of particles left.



a non-coincidence



a cosmic coincidence





a persisting coincidence

- Similar relic abundances $\Omega_{DM} \sim \Omega_{VM} \rightarrow asymmetric DM$
- Sub-galactic structure currently explained better by self-interacting DM with

 $\sigma_{scatt} / m_{DM} \sim 1 \text{ barn/GeV} \sim \sigma_{nn} / m_n$ rather than by collisionless DM

 Tentative/unconfirmed direct-detection signal [DAMA], for DM mass

 $m_{DM} \sim few \, GeV \sim m_n = GeV$

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connection between dark and ordinary matter microphysics ?

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 $\sigma_{\rm scatt}$ / m_{DM} ~ 1 barn /GeV ~ σ_{nn} / m_n

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connection between dark and ordinary matter microphysics ?

[Review of asymmetric dark matter; KP, Volkas (2013)]

The hypothesis

- DM relic abundance due to an excess of dark particles over antiparticles (asymmetry).
- Dark asymmetry related to the BAU dynamically, by processes which occurred in the early universe.
- Dark and visible asymmetries conserved separately today.



[Review of asymmetric dark matter; KP, Volkas (2013)]

Ingredients

- Low-energy theory
 - (accidental) global U(1) symmetry, "dark baryon number, B_D", conserved independently of B_V
 - Interaction which annihilates thermal symmetric population of DM: $(\sigma v)_{ann} > 6 \times 10^{-26} \text{ cm}^3/\text{s}$

High-energy theory

Joint violation of (B-L)_V and B_D

[e.g. Nussinov (1985); Kaplan (1992); Foot, Volkas (2003); Farrar, Zaharijas (2004); Hooper, March-Russell, West (2005); Agashe, Servant (2005); Suematsu (2006); Gudnason, Kouvaris, Sannino (2006); Kitano, Low (2006); Kaplan, Luty, Zurek (2009); Davoudiasl et al. (2010); Buckley Randall, (2010); Kaplan, Krnjaic, Rehermann, Wells (2011); Bell, Shoemaker, KP, Volkas (2011); KP, Trodden, Volkas (2011); von Harling, KP, Volkas (2012); Servant, Tulin (2013); Baldes, Bell, KP, Volkas (2014)]

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DM annihilation

• Need $(\sigma v)_{ann} > 6 \times 10^{-26} \text{ cm}^3/\text{s}$

• $\overline{\chi} \chi \rightarrow SM SM$

Annihilation directly into SM particles highly constrained via colliders and direct detection (see bounds on symmetric WIMP DM)

• $\overline{\chi} \chi \rightarrow \phi \phi$

Annihilation into new light states:

- * $\phi \phi \rightarrow SM SM$: metastable mediators decaying into SM
- $\star \phi$ stable light species: extra radiation

e.g. dark photon (possibly massive with kinetic mixing to hypercharge), or a new light scalar.

[Review of asymmetric dark matter; KP, Volkas (2013)]

Structure



Asymmetric dark matter with (long-range) self-interactions

Asymmetric dark matter with long-range self-interactions

- How to go about studying it?
- Many studies of long-range DM self-interactions (in either the symmetric or asymmetric regime) employ a Yukawa potential

 $V_{\chi\chi}(r) = \pm \alpha \exp(-m_{\phi} r) / r$

• However, typically reality is more complicated for asymmetric DM with long-range interactions.

Asymmetric dark matter with long-range self-interactions

Involved cosmology

 \supset Formation of bound states in the early universe

- Rich phenomenology. Could involve
 - Multi-component DM with a variety of intra- and inter-species interactions in haloes today.
 - Direct and indirect detection signals with rich structure.

Necessitates studying the preceding cosmology

 Delineate possibilities (classes of models), study cosmo+pheno self-consistently

atomic dark matter

- Dark interaction: gauged U(1)_D
 - Efficient annihilation in the early universe, into "dark photons"
 - Dark photons mediate DM self-scattering in haloes today
 - Contributes to structural complexity in the dark sector \rightarrow global U(1)_{BD} in the low-energy effective theory

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- Gauge invariance mandates DM be multi-component:
 - massless dark photon [analogous to ordinary matter]:
 U(1)_D charge carried by (dark) protons must be compensated by opposite gauge charge carried by (dark) electrons.
 - mildly broken U(1)_D, light dark photon: similar conclusion in most of the parameter space of interest. [KP, Pearce, Kusenko (2014)]

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fundamental



$$\delta L_{low} = L_{SM} + \bar{p}_{D} (i\not{D} - m_{p})p_{D} + \bar{e}_{D} (i\not{D} - m_{e}) e_{D} + (\varepsilon/2) F_{Y\mu\nu} F_{D}^{\mu\nu}$$

$$\delta L_{high} \Rightarrow (1/M^{8}) (\bar{u}^{c}d \bar{s}^{c} u \bar{d}^{c} s) \bar{e}_{D}^{c} p_{D}$$

preserves $B_{gen} = (B-L)_{V} - B_{D}$
breaks $X = (B-L)_{V} + B_{D}$

$$\begin{bmatrix} X \text{ asymmetry generation: } \Delta (B-L)_{V} = \Delta B_{D} \\ [e.g. via Affleck-Dine mechanism in susy models; von Harling, KP, Volkas (2012)] \end{bmatrix}$$

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Dark relic constituents

- Dark protons p_D (no \overline{p}_D), with mass m_p
- Equal number of dark electrons e_D (no \overline{e}_D), with mass m_e
- Dark Hydrogen atoms $H_D =$ bound states of $p_D \& e_D$, with mass $m_H = m_p + m_e \Delta$, where $\Delta = \alpha_D^2 \mu_D / 2 =$ binding energy, and $\mu_D = m_p m_e / (m_p + m_e)$.

Cosmology

Dark asymmetry generation	T _{asym} > m _p / 25
Freeze-out of annihilations $\overline{p}_{D} p_{D} \rightarrow \gamma_{D} \gamma_{D} \& \overline{e}_{D} e_{D} \rightarrow \gamma_{D} \gamma_{D}$	T _{FO} ≈ m _{p,e} / 30
Dark recombination, $p_{\rm D} + e_{\rm D} \rightarrow H_{\rm D} + \gamma_{\rm D}$	Δ / 50 < T_{recomb} < Δ , Δ = binding energy = $\alpha_D^2 \mu_D / 2$
Residual ionisation fraction	$\boldsymbol{x}_{ion} \equiv \frac{\boldsymbol{n}_{p}}{\boldsymbol{n}_{p} + \boldsymbol{n}_{H}} \sim \min\left[1, 10^{-10} \frac{\boldsymbol{m}_{p} \boldsymbol{m}_{e}}{\boldsymbol{\alpha}_{D}^{4} \mathrm{GeV}^{2}}\right]$
[If dark photon massive] Dark phase transition	T _{PT} ~ m _γ / (8πα _D) ^{1/2}

[Kaplan, Krnjaic, Rehermann, Wells (2009); KP, Trodden, Volkas (2011); Cyr-Racine, Sigurdson (2012); KP, Pearce, Kusenko (2014)]

Atomic DM with a light massive dark photon

- Asymmetry generation in both p_D and e_D if $T_{asym} > T_{PT}$ $m_V < (8\pi\alpha_D)^{1/2} m_p / 25$
- p_D and e_D interact via an attractive Yukawa potential. Bound states exist if $m_{\gamma} < 1 I a_B = \mu_D \alpha_D$. Bound states can form radiatively, $p_D + e_D \rightarrow H_D + \gamma_D$, if $m_{\gamma} < \Delta = (1/2) \mu_D \alpha_D^2$
- Dark recombination happens as for $m_{\gamma} = 0$, if $T_{end rec} > T_{PT}$ $m_{\gamma} < (8\pi\alpha_D)^{1/2} [(1/2) \mu_D \alpha_D^2 / 50]$

[KP, Pearce, Kusenko (2014)]



Atomic DM with a light massive dark photon

Asymmetric DM coupled to a dark photon is atomic in much of the parameter space where the dark photon is light enough to mediate sizeable (long-range) DM self-interactions

(i.e. not easy to avoid bound state formation)

[KP, Pearce, Kusenko (2014)]

Self-interactions of atomic DM

• Multicomponent DM with different inter- and intra-species interactions

 $H_{\rm D} - H_{\rm D}$, $H_{\rm D} - p_{\rm D}$, $H_{\rm D} - e_{\rm D}$, $p_{\rm D} - p_{\rm D}$, $e_{\rm D} - e_{\rm D}$, $p_{\rm D} - e_{\rm D}$

• Strong velocity dependence

$$\sigma_{ion-ion} \propto v^{-4}$$
, screened at $\mu_{ion-ion} v < m_{\gamma}$

$$\sigma_{H-H} \approx (\alpha_D \mu_D)^{-2} \left[b_0 + b_1 \left(\frac{m_H v^2}{4 \mu_D \alpha_D^2} \right) + b_2 \left(\frac{m_H v^2}{4 \mu_D \alpha_D^2} \right)^2 \right]^{-1}$$

(valid away from resonances; b_0 , b_1 , b_2 : fitting parameters, depend mildly on m_p/m_e) [Cline, Liu, Moore, Xue (2013)]

Atomic DM in haloes

Find the parameter space where DM self-scattering

- preserves ellipticity of large haloes; for single-component DM: $\sigma/m_{DM} < 2 \text{ barn / GeV} @ v > 200 \text{ km/s}$
- affects smaller halo dynamics; for single-component DM: $\sigma/m_{DM} > 1 \text{ barn / GeV} @ v \sim 10 \text{ km/s}$
- Atomic DM is multi-component, bounds not directly applicable
 → appropriate average over various components

Binding energy $\Delta = 0.5 \text{ MeV}$ Dark photon mass $m_v = 1 \text{eV}$



- Non-monotonic behaviour in α_D, because of the formation of bound states (→ no upper limit on α_D, or lower limit on m_v).
- Strong velocity dependence of scattering cross-sections allows for ellipticity constraints to be satisfied, while having a sizeable effect on small scales.
- Collisionless CDM limits: large $m_{\rm H} \rightarrow$ small number density large $\alpha_{\rm D} \rightarrow$ tightly bound atoms small $\alpha_{\rm D} \rightarrow$ small interaction large $m_{\gamma} \rightarrow$ no atoms, ion-ion screening small $m_{\gamma} \rightarrow$ atom formation 43

Binding energy $\Delta = 3 \text{ MeV}$ Dark photon mass $m_v = 1 \text{ MeV}$



[KP, Pearce, Kusenko (2014)]

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- DM in bound states: even massless mediators viable (and very interesting: v-dependent scattering)
- If DM mostly ionised, and m_{DM} < 500 GeV → sizable mediator mass needed
- Even if DM mostly ionised, very light/massless mediators still good, if m_{DM} > 500 GeV

[KP, Pearce, Kusenko (2014)]

Extensions

1) Detection prospects of atomic DM: $\delta \mathcal{L} = (\epsilon/2) F_Y F_D$

- Indirect detection signals with rich structure, e.g.
 - bound-state formation in the galaxy today from ionised component

 $p_{\rm D}^+ + e_{\rm D}^- \rightarrow H_{\rm D} + \gamma_{\rm D}$

 $\gamma_{D} \rightarrow e^+ e^-$ (for $m_{\gamma} > 1 \text{ MeV}$)

[KP, Pearce, Kusenko (in preparation)]

level transitions (dark Hydrogen excitations)

 $H_{\rm D} + H_{\rm D} \rightarrow H_{\rm D} + H_{\rm D}^{*}, \quad H_{\rm D}^{*} \rightarrow H_{\rm D} + \gamma_{\rm D}, \quad \gamma_{\rm D} \rightarrow e^+ e^-$

[Along similar lines: Frandsen et al. (2014), Cline et al. (2014)]

Extensions

- 1) Detection prospects of atomic DM: $\delta \mathcal{L} = (\epsilon/2) F_Y F_D$
- Direct detection involving processes with different kinematics
 - * elastic scattering of p_D , e_D , H_D on the target, and inelastic scattering of H_D (excitation or break-up)
 - contact-type and long-range DM-nucleon interactions,
 depending on the incident DM particle, screening scale (dark photon mass / Bohr radius), recoil energy.

[Frandsen, Kouvaris, KP, Shoemaker (in progress)]

Extensions

2) Consider other types of long-range interactions.

- DM interactions mediated by a scalar boson:
 always attractive → large bound states.
- Non-Abelian confining theories

• Non-relativistic freeze-out of symmetric thermal relics:

 $\sigma_{ann} v_{rel} = \sigma_0 = 6 \times 10^{-26} \text{ cm}^3/\text{s}$

• If DM is heavy or mediator is light

(m_{DM} / 2) v_{rel} > m_{mediator}

→ Sommerfeld effect:

$$\sigma_{ann} v_{rel} = \sigma_0 \quad \frac{2\pi\alpha / v_{rel}}{1 - Exp[-2\pi\alpha / v_{rel}]} \longrightarrow \sim \alpha / v_{rel} @ large \alpha / v_{rel}$$

→ Can affect freeze-out, indirect detection [see explanations of Pamela/AMS positrons]

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At freeze-out v_{rel} ~ 0.3; Holds for m_{WIMP} > 1 TeV

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- Long-range attractive interactions \rightarrow bound states. Consider e.g. Dirac fermions X, coupled to *massless* dark photons γ_D annihilation: $X + \overline{X} \rightarrow \gamma_D \gamma_D$
 - Bound-state formation and decay $\begin{cases} X + \overline{X} \rightarrow (X\overline{X})_{bound} + \gamma_{D} \\ (X\overline{X})_{bound} \rightarrow 2\gamma_{D} \text{ or } 3\gamma_{D} \end{cases}$
- Bound-state formation is also Sommerfeld-enhanced, becomes dominant inelastic process at $\alpha / v_{rel} > 0.6$
 - → Implications for freeze-out, indirect detection; hidden-sector and possibly heavy WIMP DM



[[]von Harling, KP (2014)]

Why is this important?



[von Harling, KP (2014)]



Implications for indirect detection



Conclusion

- Symmetric thermal-relic WIMP DM ↔ collisionless CDM Asymmetric (thermal relic) DM ↔ self-interacting DM
- Involved cosmology determines low-energy phenomenology: DM self-scattering in haloes, direct and indirect detection.

Cosmology

near coincidence of dark & visible matter abundances

Observations

clustering in subgalactic scales