

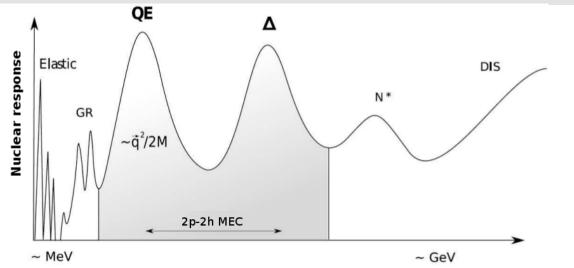
Inelastic neutrino-nucleus scattering in the superscaling model

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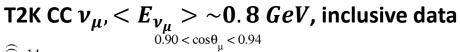
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Introduction

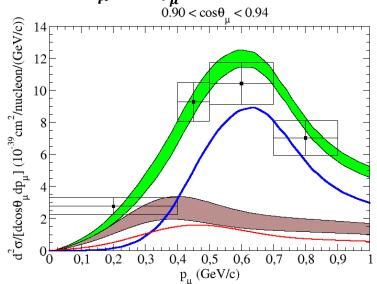


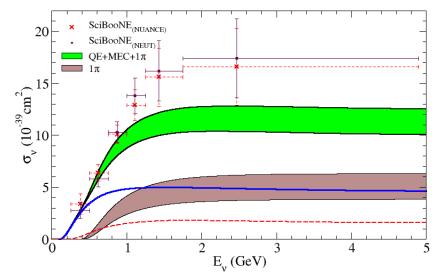


- Quasielastic region.
- 2p-2h excitations.
- Δ resonance, other resonances and DIS.



Total Inclusive Cross Section



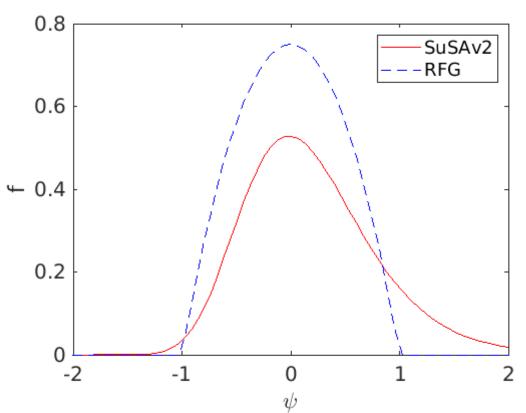


Superscaling model with QE + 2p2h + 1π [M.V Ivanov et al., J. Phys. G 43, 045101 (2016)].

Introduction.



Comparison between SuSAv2 and RFG scaling function.



$$f(\psi) = k_F \frac{\left(\frac{d^2 \sigma}{d\Omega dw}\right)}{\left(\frac{d^2 \sigma}{d\Omega dw}\right)_{s.n}}$$

The scaling function does not depend explicitly on the transferred momentum or the nuclear species [J. E. Amaro et al., J. Phys. G 47, 124001 (2020), G. D. Megias, PhD Thesis (2017)].

SuSAv2 model comes from RMF.

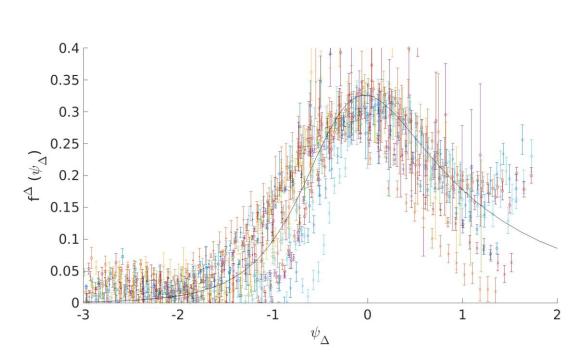
SuSAv2-QE scaling function is going to be implemented in the inelastic regime.

Model: 1π - $f\Delta$



In order to describe Δ resonance region, a Δ pion production model [PRC 71, 015501 (2005)] is used with a phenomenological scaling function.

$$[W^{\mu\nu}]^{\Delta} = \frac{1}{2} \Lambda_0 f^{model} U^{\mu\nu}$$



$$\begin{split} &\left(\frac{d^2\sigma}{d\Omega d\omega}\right)^{\Delta} \\ &= \left(\frac{d^2\sigma}{d\Omega d\omega}\right)^{exp} - \left(\frac{d^2\sigma}{d\Omega d\omega}\right)^{QE} \\ &- \left(\frac{d^2\sigma}{d\Omega d\omega}\right)^{MEC} - \left(\frac{d^2\sigma}{d\Omega d\omega}\right)^{DIS} \end{split}$$

$$f^{\Delta}(\psi_{\Delta}) = k_F \frac{\left(\frac{d^2 \sigma}{d\Omega d\omega}\right)^{\Delta}}{\sigma_{Mott}(v_L G_L^{\Delta} + v_T G_T^{\Delta})}$$

Model:SuSAv2-inelastic



SuSAv2-inelastic model describes the full inelastic spectrum (Δ , other res. And DIS)[G. D. Megias, PhD Thesis (2017), M. B. Barbaro et al. Phys. C 69, 035502 (2004)]. Good agreement with (e,e') data.

$$R_{inel}^{K}(\kappa,\tau) = \frac{N}{\eta_F^2 \kappa} \xi_F \int_{\mu_X^{min}}^{\mu_X^{max}} d\mu_X f^{model}(\psi_X') U^k$$

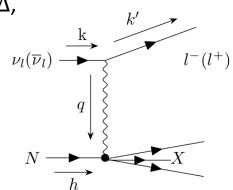
The hadronic response is given by an integration of the single-nucleon tensor over the invariant mass.

The limits of this integral are

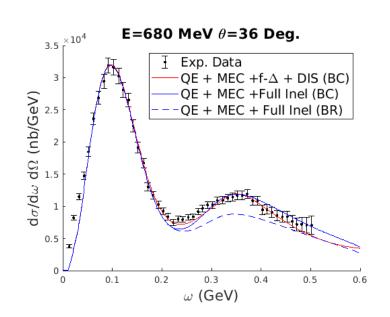
$$\mu_X^{min}=1+rac{m_\pi}{M_N}$$
 , $\mu_X^{max}=1+2\lambda-rac{E_S}{M_N}$

This limits can be changed to work alongside a resonance model.

BR and BC parametrizations (specially BC) work well. PDF gets closer at high ω , but it is not suited to describe Δ region.

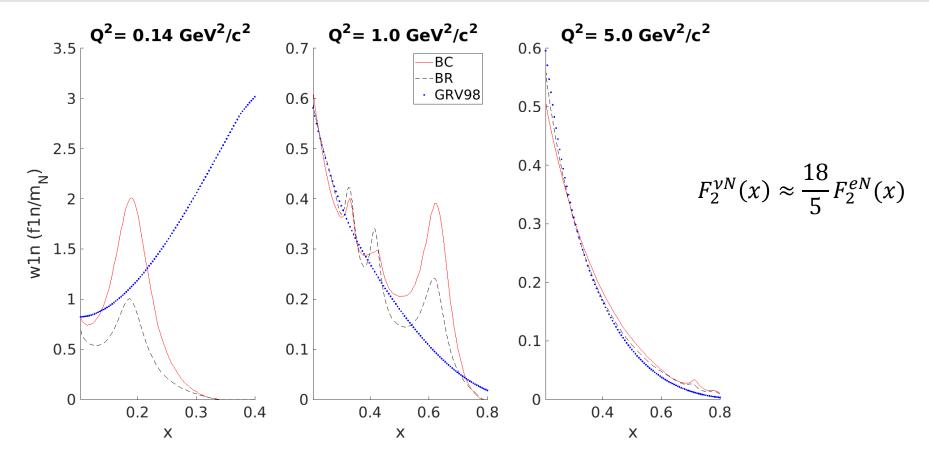


Inelastic Feynmann Diagram



Model:SuSAv2-inelastic





Bodek-Ritchie parametrization [PRD 23, 1070 (1981)].

$$xF_3 = F_2 - 2\bar{Q} \longrightarrow$$

Parton Distribution Function.

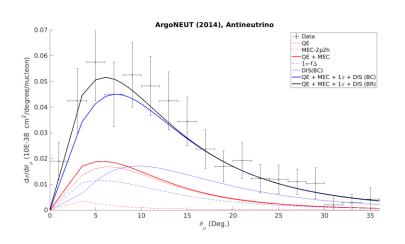
$$xF_3^{vp(n)} = 2x(d(u) + s - \bar{u}(\bar{d}) - \bar{c})$$

Antiquarks

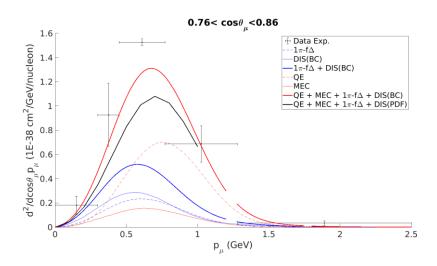
distribution.

Results

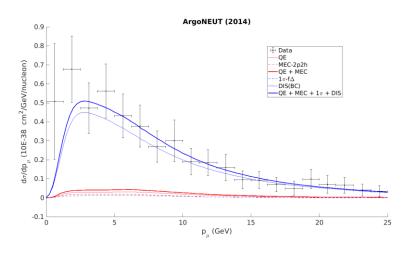




ArgoNEUT CC $\overline{
u}_{\mu}$, < $E_{
u_{\mu}}$ > \sim 3. 6 GeV



Microboone CC u_{μ} , $< E_{
u_{\mu}} > \sim 0.8 \ GeV$



ArgoNEUT CC u_{μ} , < $E_{
u_{\mu}}$ > \sim 9. 6 GeV

Using 1π -f Δ and the SuSA2-inelastc model (DIS). In the case of ArgoNEUT these contribution are needed to explain the experimental data.

Conclusion



- It is necessary to include an analysis of the inelastic scattering to explain the neutrino cross section at certain kinematics.
- The SuSAv2-inelastic works well for electrons and it is expected to performs well for neutrinos.
- In the case of neutrinos, the best option is use a resonance model to reproduce Δ contributions and the SuSAv2-inelastic BC for higher contributions.
- It is necessary better parametrization model at intermediates energies.



Thanks for your attention

Limits of the inelastic region



Kinematically allowed region, recoiling of the daughter nucleus

$$\max[\varepsilon(0), 0] \le \varepsilon \le \varepsilon(\pi)$$

Considering that the mass of daughter nuclei is infinite

$$\varepsilon_{\infty}(\theta) = m_N + \omega - \sqrt{W_X^2 + q^2 + p^2 + 2pq\cos\theta}$$

Considering the limits

$$m_N + m_\pi \le W_X \le m_N + \omega - E_S$$

W1, W2 and W3



