

Alberto Mariotti



Based on arXiv:2203.16450 (PRL) with S. Blasi and on arXiv:2312.06749 with P. Agrawal, S.Blasi and M. Nee

BIG&C Meeting

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Early Universe Cosmology

★How looks Universe before its first second?



Background of Gravitational waves can probe early universe cosmology





Stochastic Background of GW







★AstroPhysical SGWB

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* Superposition of unresolvable sources

BBH BNS

* Predictable after LIGO/Virgo observations LIGO/Virgo Phys.Rev.D 100 (2019)

! Most likely measured in next few years !



★Cosmological SGWB

* Generated by energetic events during cosmological evolution



★Consider Universe with high reheating temperature ★Consider spontaneously broken symmetry

 ϕ : order parameter of a symmetry



Phase transition during the cooling of the Universe at critical temperature $\sim T_c$



★In the Standard Model

*QCD Phase Transition (T ~ GeV)? In SM No first order

*EW Phase Transition (T~ 100 GeV)? In SM No first order

(If very light Higgs it could have been strongly first order) '81 Witten



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******Probe of Early Universe cosmology*

★In the Standard Model, QCD and EW PT are not first order

First Order Phase Transition would be signal of BSM physics

- * First order EW PT can lead to electroweak baryogenesis
- * New first order PT in dark matter sectors
- * FOPT are powerful sources of stochastic GW signal



Target for current and future GW experiments

> Ligo Virgo Kagra + NANOGrav, PPTA, EPTA

- + LISA
- + Einstein Telescope



+ Cosmic explorer

→ New physics in

Higgs sector

★First Order Phase transition (FOPT) proceeds with bubble nucleation★Nucleation condition in homogeneous Universe



* Phase transition described by effective potential
* Thermal fluctuation induces nucleation of bubbles
* Nucleation rate/volume set by O(3) bounce action

 $\gamma_V(T) \sim T^4 e^{-S_3(T)/T}$

* Nucleation condition sets nucleation temperature

 $\gamma_V(T_n) \sim H(T_n)^4$



+Nucleation rate controlled by the bounce action





fig. from arXiv:1705.01783 D. Weir

★Parameters describing PT properties

$$\alpha_* \simeq \frac{\Delta V}{\rho_{rad}} \Big|_{T_*} \overset{\text{Latent heat}}{\overset{\text{Latent heat}}{\overset{the}}{\overset{th}}}{\overset{th}}{\overset{th}}{\overset{th}}{\overset{th}}{\overset{th}}{\overset{th}}{\overset{th}}{\overset{th}}{\overset{th}}{\overset{th}}{\overset{th}}{\overset{th}}{\overset{th}}{\overset{th}}{\overset{th}}}{\overset{th}}{$$

 $\frac{\beta}{H_*} = T \frac{d}{dT} \frac{S_3}{T} \Big|_{T_*} \operatorname{Tim_{\Theta-SCale}}$

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 $\begin{aligned} f_{peak} &\sim 10^{-3} \, \mathrm{Hz} \left(\frac{\beta/H_*}{100}\right) \left(\frac{H_n}{100}\right) \\ & \text{Figure 4: Kinetic energy 00 in difference of statution snat} \\ & (\text{top right}), 10.8/\beta \text{ (bottom left) and } 20.1/\beta \text{ (bottom right)}, 10.8/\beta \text{ (bottom left) and } 20.1/\beta \text{ (bottom right)}, 10.8/\beta \text{ (bottom left) and } 20.1/\beta \text{ (bottom right)}, 10.8/\beta \text{ (bottom left) and } 20.1/\beta \text{ (bottom right)}, 10.8/\beta \text{ (bottom left) and } 20.1/\beta \text{ (bottom right)}, 10.8/\beta \text{ (bottom left)}, 20.1/\beta \text{ (bottom right)}, 10.8/\beta \text{ (bottom left)}, 20.1/\beta \text{ (bottom right)}, 10.8/\beta \text{ (bottom left)}, 20.1/\beta \text{ (bottom right)}, 20.1/\beta \text{ (bottom$

3 mechanisms to generate SBGW from FO

- + Bubble collisions
- + Sound Waves in the plasma
- + Turbulence

Many subtleties in computation of GW signal (

★GW signal is broken power law



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★Impurities drastically modify the nucleation process

Bubble chamber



Supercool water



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... however ... impurities

★Impurities can play a role also in cosmological phase transitions

MONOPOLE AND VORTEX DISSOCIATION AND DECAY OF THE FALSE VACUUM

Paul Joseph STEINHARDT

Lyman Laboratory of Physics, Harvard University, Cambridge, Massachusetts 02138, USA

Received 17 February 1981

"If monopole (or vortex) solutions exist for a metastable or false vacuum, a finite density of monopoles (or vortices) can act as impurity sites that trigger inhomogeneous nucleation and decay of the false vacuum." Impurities in the early universe

Yutaka Hosotani Department of Physics, University of Pennsylvania, Philadelphia, Pennsylvania 19104 (Received 1 November 1982)

"Now one has to ask the following question: Is the early universe really sufficiently pure in order for supercooling to take place? The aim of this paper is to show that in most cases the early universe is very pure. [...] In this paper we consider ordinary particles as impurities."

Cosmic separation of phases

Edward Witten* Institute for Advanced Study, Princeton, New Jersey 08540 (Received 9 April 1984)

"In particle physics it is often assumed that phase transitions are nucleated by thermal fluctuations. In practice, [...] except in very pure, homogeneous samples, **phase transitions are often nucleated by various forms of impurities and inhomogeneities of nonthermal origin**."

"What if the transition was nucleated by impurities? In this case **the mean spacing between bubbles has nothing to do with free energies** of nucleation and is simply the spacing between the relevant impurities."

★The nature of impurities for cosmological PT



E.g. Affleck, De Luccia '79, --- Selivanov, Voloshin '85, --- Kuznetsov, Tinyakov '97 --- Strumia '23

* Compact objects like BH, gravitational effects

E.g. Hiscock '87, -- Gregory, Moss, Withers '14, --Grinstein, Murphy '15, -- El-Menoufi, Huber, Manuel '20, Balkin et al '21, Strumia '22, Jinno et at. '23

Fig. from Oshita et al.1808.01382



* *Topological defects* (strings, monopoles ...)







★The nature of impurities for cosmological PT





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★What is the origin of the topological defects?

*Remnants of PT depending on vacuum manifold topology [Zel'dovich et al. '74, Kibble '76]





★Topological defects are present if EW symmetry breaking is final step of a multi-step breaking of larger symmetry group

Typical in unified theories

Simplest model for EW FOPT

★Higgs (h) plus Singlet S with a $\mathbb{Z}_2 : S \to -S$ $V(h,S) = -\mu_h^2 |\mathcal{H}|^2 + \lambda |\mathcal{H}|^4 - \frac{\mu_s^2}{2}S^2 + \frac{\eta}{4}S^4 + \kappa |\mathcal{H}|^2S^2$

★The electroweak phase transition occurs in two steps



★Many pheno studies on Higgs Singlet EWPT

- * Simplest new physics scenario with strong EW FOPT [Espinosa, Konstandin, Riva 1107.5441]
- * Minimal mechanism for EW baryogenesis

[Espinosa, Gripaios, Konstandin, Riva 1110.2876]

- * Benchmark for gravitational wave signals [Caprini et al 1512.06239]
- * Singlet is challenging to detect at colliders [Curtin, Meade, Yu 1409.0005]
- ♦ Order O(1000) papers on this model in last 10 years



Fig. adapted from Kurup, Perelstein [1704.03381] PRD

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... but Domain Walls ...



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... but Domain Walls ...



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The rolling (classical instability)

★Domain walls can become classically unstable while the Universe cools down★Developing a region of the true vacuum in their interior, and then dissociate





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Steinhardt '81: Same phenomena for monopoles

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★If classically stable, Domain wall can act as seeds for the EW phase transition

★*Tunneling can occur in two competing processes*



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Cosmological history of Seeded vs homogeneous PT



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Cosmological history of Seeded vs homogeneous PT



Cosmological history of Seeded vs homogeneous PT







How to characterize and compute the seeded tunneling?

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★Look for nucleation probability at the DW location



Geometry of seeded critical bubble



O(2) symmetry on the DW plane

* Nucleation rate/surface set by O(2) bounce action

$$\gamma_S \sim T^3 e^{-S_2/T}$$

* Seeded nucleation condition

$$\gamma_S(T_n^{seed}) \sim \frac{1}{\xi} H^3(T_n^{seed})$$

Remind for comparison homogeneous tunneling $\gamma_V(T) \sim T^4 e^{-S_3(T)/T}$ Nucleation: $\gamma_V(T_n) \sim H(T_n)^4$

★Look for nucleation probability at the DW location





O(2) symmetry on the DW plane

* Nucleation rate/surface set by O(2) bounce action



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Thin wall approximation



★Study 3-dimensional theory on the domain wall

- *** KK spectrum contains massive localized states gapped to a continuum**
- ***** Bound states correspond to localized profiles in the z-direction
- * Scattering states correspond to continuum





✦ Metastability of DW controlled by the 3d h mass

 $\omega_0^2(T) < 0$ Classical instability of DW $\triangleleft \cdots \rightarrow$ The "rolling"

 $\omega_0^2(T) > 0 \ \frac{\text{Classically stable DW}}{\text{Seeded Tunnelling at T<Tc}}$

Kaluza Klein reduction method

★Study 3-dimensional theory on the domain wall

- *** KK spectrum contains massive localized states gapped to a continuum**
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Numerics: mountain pass algorithm

★Numerical solving for the PDE's is a non trivial task

- * Reduced symmetry (only O(2))
- ***** Bounce action is a saddle point



Figure from: Agrawal, Nee '22 SciPost

Bubble profile and bounce action found for full thermal effective potential Employ a numerical algorithm suitable to find saddles Designed to find the "mountain pass" between two points Agrawal, Nee '22 SciPost



Benchmark 1, T = 105 GeV

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Bounce action in 3 methods



EW Phase transition seeded by the DW always dominates !

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★Certain quality of the Z2 symmetry is needed for DW to exist

*** Z2 symmetry explicitly broken by Planck scale suppressed operator is allowed**

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Properties of the seeded PT?



★ Criteria for the seeded phase transition to complete (percolation)
 ★ Special features in latent heat and typical time scale?
 ★ Special features in gravitational waves?

See Blasi, Konstandin, Rubira, Stomberg '23 for GW from sound waves in seeded PT



★ Generic network $\beta = Min(\beta_{DW}, \beta_{\xi})$

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Seeded PT features controlled by O(2) bounce

$$\beta_{DW} = -\frac{dS_{DW}}{dt}\Big|_{T_n}$$

★ Percolation temperature

 $I(T_p) \simeq 0.34 \begin{array}{c} \text{Usual expression for} \\ \text{volume fraction in true} \\ \hline \mathbf{v} \dots & \mathbf{v} \end{array}$

★ Bubble size at percolation

$$R_{\rm p}^{\rm 3D} = -v_w (8\pi)^{1/3} \left(\left. \frac{dB_{\rm DW}}{dt} \right|_{T_{\rm p}} \right)^{-1}$$

Latent Heat and bubble size are different

Sparse DW network



Seeded PT features controlled by number of DW

$$\beta_{\xi} = v_w \xi H \Big|_{T_n}$$

★Percolation temperature



For monopoles: Guth-Weimberg '81

★Bubble size at percolation

$$R_p^{\xi} = \frac{n_c}{2\xi H} \left(1 + \frac{n_c}{2\xi v_w} \right)$$

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★Decorrelation between bubble size and latent heat

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★*How GW signal from networks compare with homogeneos one*







Thanks for the attention

Backup slide

★Percolation temperatures for homogeneous and seeded tunneling



