Color coherence in multiple antenna medium radiation

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Motivation: jet substructure



- Ideal techniques for heavy ion collisions.
- More direct access to the underlying dynamics:
 - QGP properties.
 - Energy loss.
 - Coherence.

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Color coherence in vacuum

• Is radiation independent?: $q\bar{q}$ antenna as a laboratory.



$$dN = \frac{d\omega}{\omega} \frac{d\Omega}{2\pi} \frac{\alpha_{s} C_{F}}{2\pi} \Big[R_{q} + R_{\bar{q}} - 2\mathcal{J} \Big]$$

- The spectrum is suppressed at large angles due to the presence of destructive inteferences (coherence).
- Angular ordering.

Color coherence in a medium

• How does the medium change this picture?



• A parton can change color through interaction with the medium, breaking the correlation between emitted gluons.

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Particle propagation in matter



$$W(ec{x}) = \mathcal{P} \, exp \Big[ig \int dx_+ A_-(x_+, ec{x}) \Big]$$

- The effect of the medium is to induce color rotation at each scattering center.
- The quark (a high energy quark) loses a negligible amount of energy and propagates in straight lines (*eikonal* propagation).

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In-medium antenna radiation

• To study the degree of coherence we a take a very soft gluon $\omega \rightarrow 0$ (out-out radiation).



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The decoherence parameter

• The interaction of the $q\bar{q}$ pair with the medium is described by the survival probability S.

$$egin{aligned} \mathcal{S} &\equiv rac{1}{N_c^2-1} \Big\langle W(ec{x}_\perp) W^\dagger(ec{y}_\perp) \Big
angle \ &\mathcal{S} &\equiv 1-\Delta_{med}(t) \ &\Delta_{med} &\equiv 1-exp \Big[-rac{1}{4} \hat{q} L(ec{x}_\perp-ec{x}_\perp)^2 \end{aligned}$$

• This factor determines a characteristic time-scale for decoherence of the $q\bar{q}$ pair.

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The resulting spectrum

$$dN = \frac{d\omega}{\omega} \frac{d\Omega}{2\pi} \frac{\alpha_s C_F}{2\pi} \Big[R_q + R_{\bar{q}} - (1 - \Delta_{med}) \ 2\mathcal{J} \Big]$$

$$\Delta_{med}
ightarrow 0: dN \sim R_q + R_{ar{q}} - 2\mathcal{J}$$

 $\boxed{Dilute medium: coherence (angular ordering)}$

$$\Delta_{med}
ightarrow 1: d\mathsf{N} \sim \mathsf{R}_q + \mathsf{R}_{ar{q}}$$

Opaque medium : decoherence (two independent emitters)

[The radiation pattern of a QCD antenna in a diluite/dense medium, Yacine Mehtar-Tani, Carlos A. Salgado and Konrad Tywoniuk]

Main limitations

- We have to deal with more realistic settings:
 - Non-eikonal antenna.
 - Multiple emissions.
 - Finite formation time.

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Multiple emissions

- The antenna provides a simple and intuitive picture.
- Does it hold for more than two emitters?



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Direct terms



 $|\mathcal{M}_1|^2 \propto C_F^2$ $|\mathcal{M}_2|^2 \propto C_F^2$ $|\mathcal{M}_3|^2 \propto N_c \ C_F^2$

• The direct terms are proportional to a color factor.

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Interference terms



Large N_c limit

 $\mathcal{M}_1\otimes \mathcal{M}_3^*\propto \mathcal{S}(t,L)$

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Interference terms



Large N_c limit

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 $\mathcal{M}_2\otimes \mathcal{M}_3^*\propto \mathcal{S}(0,t)\;\mathcal{S}(t,L)$

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Multiple emissions results

- We have considered the case of three emitters.
- The interference terms are proportional to the survival probabilities S in the (0, t) and (t, L) regions: the general result of the antenna is valid for each of the smaller antennas.
- If coherence is not preserved after the in-medium splitting, the antenna won't radiate coherently in the following emission.
- These computations can be generalized to the problem of *n* emitters.

Summary

- Detailed measurements of jet substructure shed light on the intricate nature of the jet interactions with the dense, deconfined QCD matter formed in the collisions.
- Color coherence is essential to understand the jet constituents' energy loss (are they independent or not?).
- In spite of the singlet antenna limitations (eikonal propagation, zero formation time, only one splitting...), it is a very convenient *laboratory*.
- The general result of the singlet antenna is valid for the subsequent antennas in the multiple emissions case.
- These computations go a step forward to obtain a complete description of a QCD cascade.

Thanks for your attention

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