

# Dark Matter Production via the Vev Flip-Flop

Lukas Mittnacht

THEP  
Johannes Gutenberg University  
Mainz, Germany

TAE, International Summer School on High Energy Physics, Benasque,  
2017

- 1 The Vev Flip-Flop
- 2 Computing the DM Abundance
- 3 Dark Matter Decay Scenario
- 4 Vev Induced Mixing

# The Vev Flip-Flop

- Add a scalar field  $S$  to the SM Lagrangian.
- Coupled via a Higgs Portal coupling

$$-V \supset \mu_H^2 H^\dagger H - \lambda_H (H^\dagger H)^2 + \mu_S^2 S^\dagger S - \lambda_S (S^\dagger S)^2 - \lambda_P (H^\dagger H)(S^\dagger S)$$

# The Vev Flip-Flop

- Add a scalar field  $S$  to the SM Lagrangian.
- Coupled via a Higgs Portal coupling
- Include thermal effects and 1-loop contributions into effective Potential

$$-V \supset \mu_H^2 H^\dagger H - \lambda_H (H^\dagger H)^2 + \mu_S^2 S^\dagger S - \lambda_S (S^\dagger S)^2 - \lambda_P (H^\dagger H)(S^\dagger S) + V^{1\text{-loop}}(T) + V^{1\text{-loop}}$$

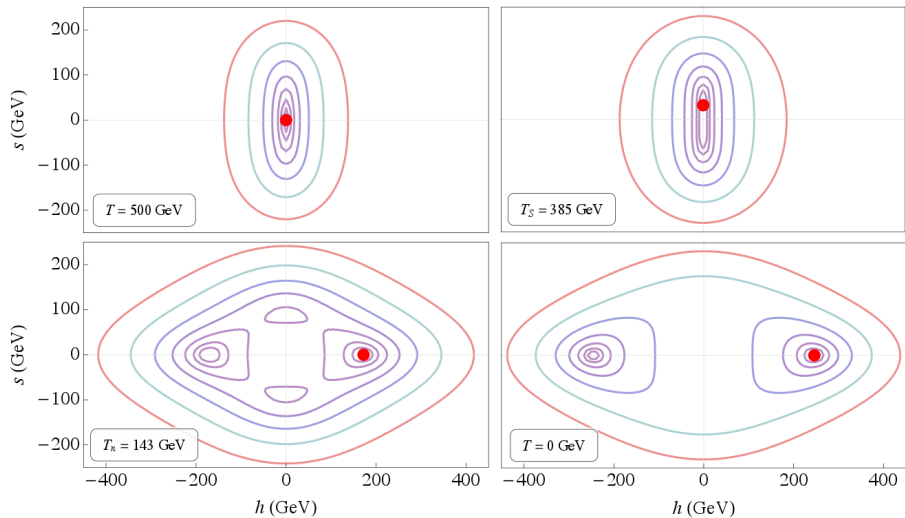
# The Vev Flip-Flop

- Add a scalar field  $S$  to the SM Lagrangian.
- Coupled via a Higgs Portal coupling
- Include thermal effects and 1-loop contributions into effective Potential

$$-V \supset \mu_H^2 H^\dagger H - \lambda_H (H^\dagger H)^2 + \mu_S^2 S^\dagger S - \lambda_S (S^\dagger S)^2 - \lambda_P (H^\dagger H)(S^\dagger S) + V^{1\text{-loop}}(T) + V^{1\text{-loop}}$$

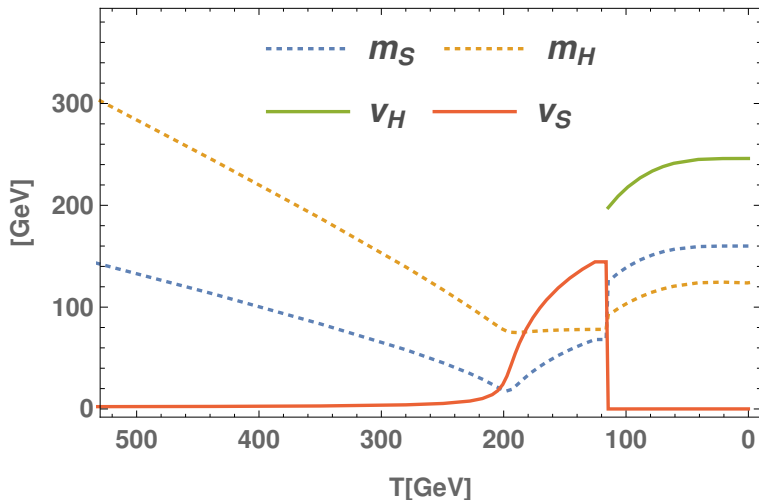
Choose coefficients such that  $S$  undergoes two phase-transitions as thermal corrections become subdominant.

# The Vev Flip-Flop



# The Vev Flip-Flop

- Thermal Evolution tracked with *CosmoTransitions*



# Computing the DM Abundance

- Need to track number density through thermal evolution of universe
- Using the output of *CosmoTransitions*
- Solve Boltzmann Equation(s) with temperature dependent masses and vev's
- Via repeated numerical ODE solving for small temperature(time) steps



# The Boltzmann Equation

- Describes non-equilibrium dynamics
- Depending on the process, different RHS (Collisionterm)
- e.g DM Annihilation into SM (2-to-2)

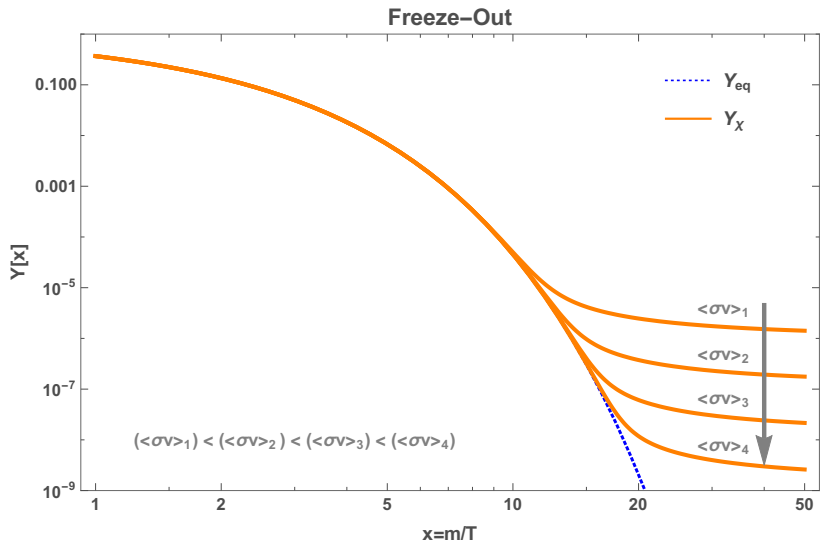
$$\begin{aligned} \frac{dn_\chi}{dt} + 3Hn_\chi = & - \int d\Pi_{\bar{\chi}} d\Pi_\chi \Pi_{\bar{f}} d\Pi_f \\ & (2\pi)^4 \delta(p_{\bar{\chi}} + p_\chi - p_{\bar{f}} - p_f) \\ & [ |M_{\bar{\chi}\chi \rightarrow \bar{f}f}|^2 f_{\bar{\chi}} f_\chi (1 \pm f_{\bar{f}})(1 \pm f_f) \\ & - |M_{\bar{f}f \rightarrow \bar{\chi}\chi}|^2 f_{\bar{f}} f_f (1 \pm f_{\bar{\chi}})(1 \pm f_\chi) ] \end{aligned}$$

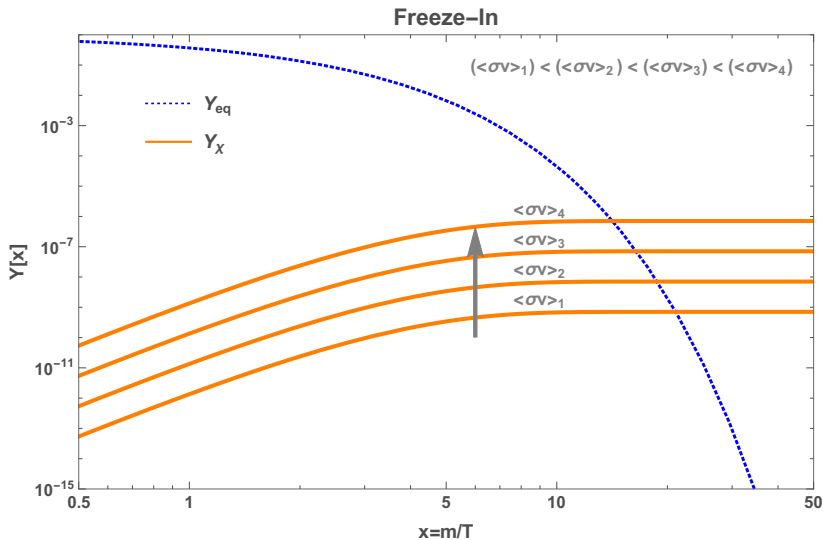
# The Boltzmann Equation

Often approximations and simplifications can be made

- $f_{\bar{f}} = f_f$  &  $f_{\bar{\chi}} = f_{\chi}$
- $|M_{\bar{f}f \rightarrow \bar{\chi}\chi}|^2 = |M_{\bar{\chi}\chi \rightarrow \bar{f}f}|^2$
- neglect Pauli-Blocking/Bose-Condensation  $(1 \pm f_i) \approx 1$
- depending on mass of particles and temperature  $f_i \approx e^{-E_i/T}$   
(neglecting chemical potential  $\mu$ )
- Changing variables from  $t$  to  $x = m/T$  and from  $n$  to  $Y = n/s$  often convenient

$$\Rightarrow \frac{dY_{\chi}}{dx} = -\frac{s \langle \sigma v \rangle}{Hx^2} [Y_{\chi}^2 - (Y_{\chi}^{eq})^2]$$





Field	Spin	Mass	$SU(3)_c \times SU(2)_L \times U(1)_Y$	$\mathbb{Z}_3$
$\chi$	$1/2$	$\mathcal{O}(1 \text{ TeV})$	$(1, 1, 0)$	$+120^\circ$
$\Psi^{(\prime)}$	$1/2$	$\mathcal{O}(1 \text{ TeV})$	$(1, 3, 0)$	$-120^\circ$
$S$	$0$	$\mathcal{O}(100 \text{ GeV})$	$(1, 3, 0)$	$+120^\circ$

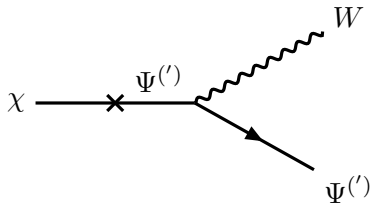
$$\mathcal{L}_{\text{Yuk}} = y_\chi^{(\prime)} S^\dagger \bar{\chi} \Psi^{(\prime)} + y_\Psi \epsilon^{ijk} S^i \bar{\Psi}^j (\Psi^{lk})^c + h.c.$$

# Dark Matter Decay Scenario - arXiv:1608.07578

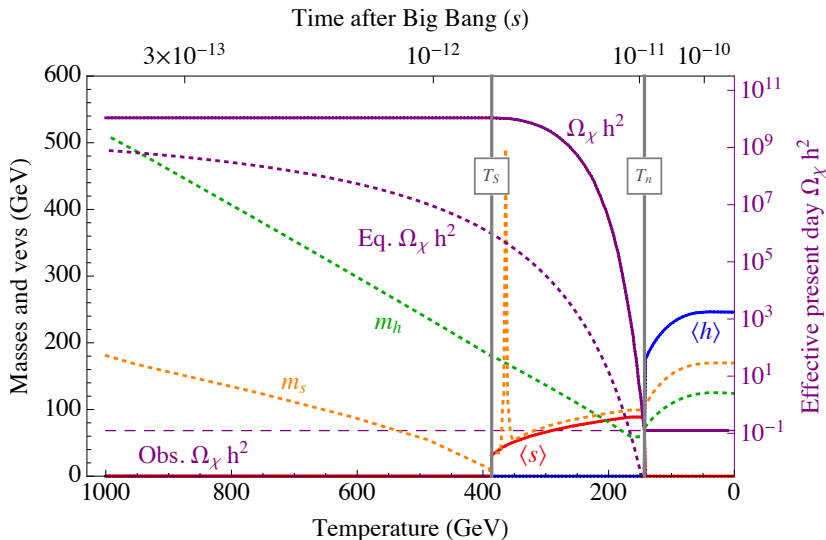
Field	Spin	Mass	$SU(3)_c \times SU(2)_L \times U(1)_Y$	$\mathbb{Z}_3$
$\chi$	$1/2$	$\mathcal{O}(1 \text{ TeV})$	$(1, 1, 0)$	$+120^\circ$
$\Psi^{(\prime)}$	$1/2$	$\mathcal{O}(1 \text{ TeV})$	$(1, 3, 0)$	$-120^\circ$
$S$	$0$	$\mathcal{O}(100 \text{ GeV})$	$(1, 3, 0)$	$+120^\circ$

$$\mathcal{L}_{\text{Yuk}} = y_\chi^{(\prime)} S^\dagger \bar{\chi} \Psi^{(\prime)} + y_\Psi \epsilon^{ijk} S^i \bar{\Psi}^j (\Psi^{tk})^c + h.c.$$

During the  $\langle S \rangle \neq 0$  phase  $\chi$  can mix into  $\Psi^{(\prime)}$  and thus decay via:



# Dark Matter Decay Scenario - arXiv:1608.07578



# Vev Induced Mixing

Field	Mass	SM	$\mathbb{Z}_2$	$\mathbb{Z}'_2$
$\chi$	$\mathcal{O}(100 \text{ GeV})$	(1, 1, 0)	+1	-1
$\psi$	$\mathcal{O}(1 \text{ TeV})$	(1, 1, 0)	-1	-1
$S$	$\mathcal{O}(100 \text{ GeV})$	(1, 1, 0)	-1	+1

$$L_{int} \supset y_\chi S \bar{\chi} \psi + h.c.$$



# Vev Induced Mixing

Field	Mass	SM	$\mathbb{Z}_2$	$\mathbb{Z}'_2$
$\chi$	$\mathcal{O}(100 \text{ GeV})$	(1, 1, 0)	+1	-1
$\psi$	$\mathcal{O}(1 \text{ TeV})$	(1, 1, 0)	-1	-1
$S$	$\mathcal{O}(100 \text{ GeV})$	(1, 1, 0)	-1	+1

$$L_{int} \supset y_\chi S \bar{\chi} \psi + h.c.$$

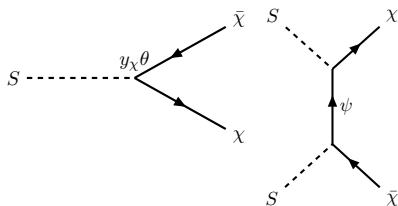
- Reheating Temp.  $< m_\psi$
- Freeze-In Scenario

# Vev Induced Mixing

Field	Mass	SM	$\mathbb{Z}_2$	$\mathbb{Z}'_2$
$\chi$	$\mathcal{O}(100 \text{ GeV})$	(1, 1, 0)	+1	-1
$\psi$	$\mathcal{O}(1 \text{ TeV})$	(1, 1, 0)	-1	-1
$S$	$\mathcal{O}(100 \text{ GeV})$	(1, 1, 0)	-1	+1

$$L_{int} \supset y_\chi S \bar{\chi} \psi + h.c.$$

- Reheating Temp.  $< m_\psi$
- Freeze-In Scenario



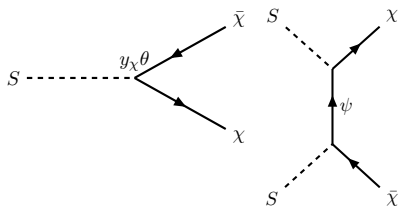
# Vev Induced Mixing

Field	Mass	SM	$\mathbb{Z}_2$	$\mathbb{Z}'_2$
$\chi$	$\mathcal{O}(100 \text{ GeV})$	(1, 1, 0)	+1	-1
$\psi$	$\mathcal{O}(1 \text{ TeV})$	(1, 1, 0)	-1	-1
$S$	$\mathcal{O}(100 \text{ GeV})$	(1, 1, 0)	-1	+1

$$L_{int} \supset y_\chi S \bar{\chi} \psi + h.c.$$

- Work in progress
- Likely only successful in some regions of parameter space

- Reheating Temp.  $< m_\psi$
- Freeze-In Scenario



# Vev Induced Mixing

Field	Mass	SM	$\mathbb{Z}_2$	$\mathbb{Z}'_2$
$\chi$	$\mathcal{O}(100 \text{ GeV})$	(1, 1, 0)	+1	-1
$\psi$	$\mathcal{O}(1 \text{ TeV})$	(1, 1, 0)	-1	-1
$S$	$\mathcal{O}(100 \text{ GeV})$	(1, 1, 0)	-1	+1

$$L_{int} \supset y_\chi S \bar{\chi} \psi + h.c.$$

- Work in progress
- Likely only successful in some regions of parameter space

- Reheating Temp.  $< m_\psi$
- Freeze-In Scenario

