Direct detection of dark matter through molecular excitations

<u>J. Pérez-Ríos</u>¹, H. Ramani², O. Slone³, E. Figueroa⁴ and R. Essig⁵ ¹School of Science and Technology, Universidad del Turabo, USA ²Berkeley Center for Theoretical Physics, Berkeley University, USA ³Princeton Center for Theoretical Science, Princeton University, USA ⁴Department of Physics and Astronomy, Stony Brook University, USA ⁵C. N. Yang Institute for Theoretical Physics, Stony Brook University, USA



Stony Brook

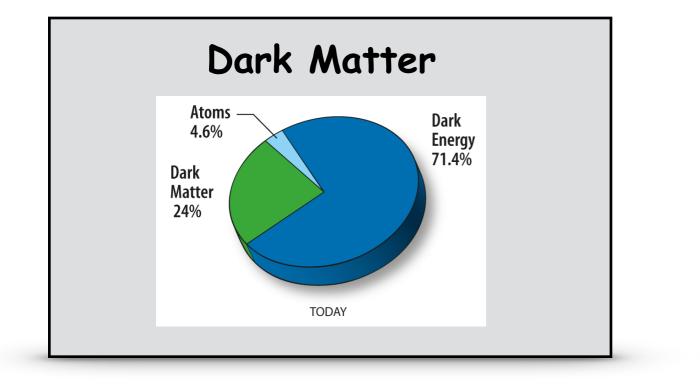
University

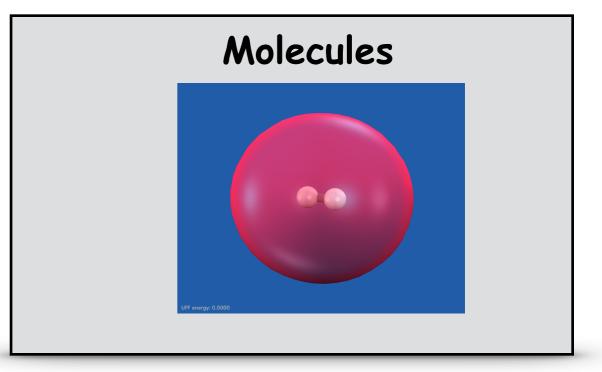


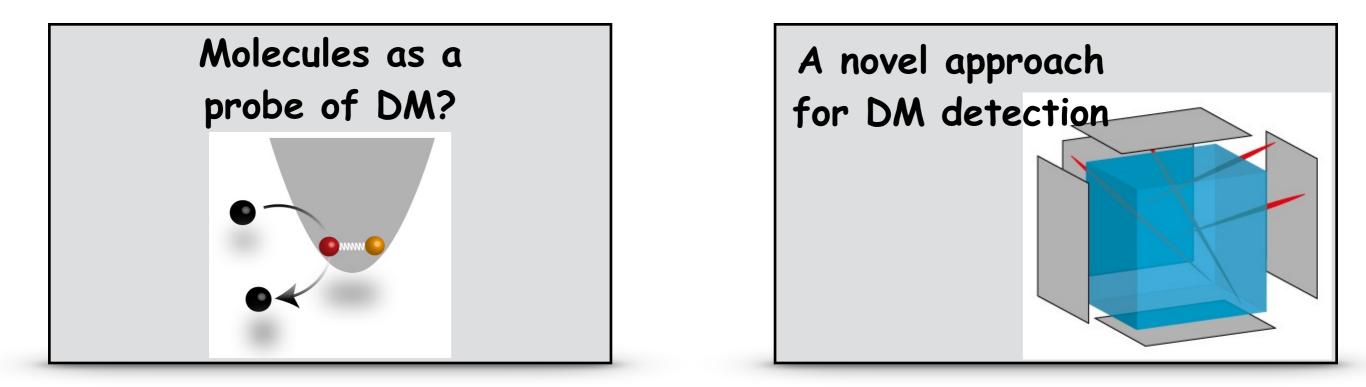


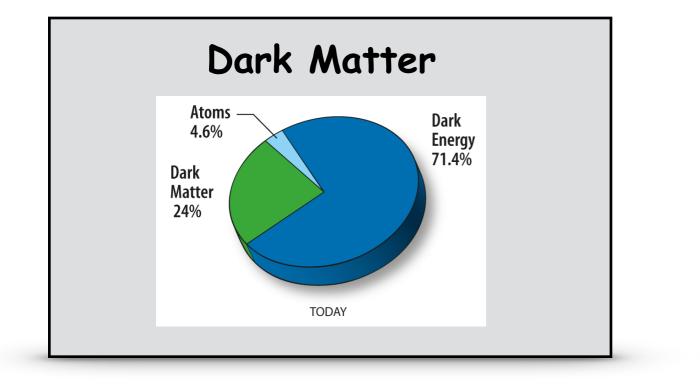
BERKELEY CENTER FOR THEORETICAL PHYSICS

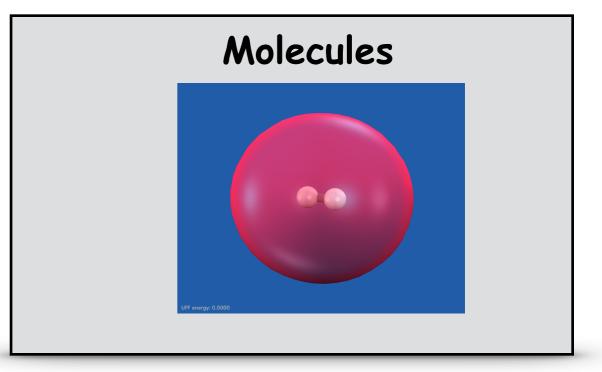


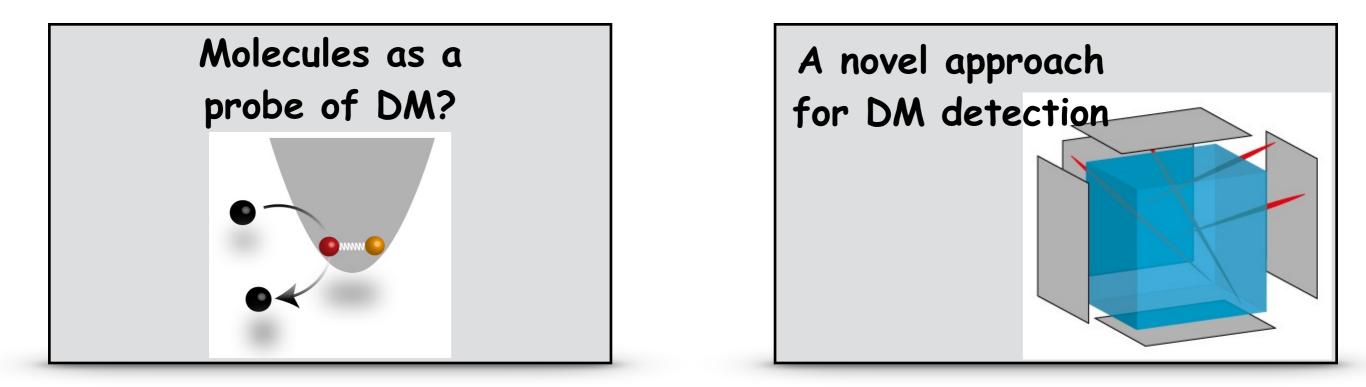


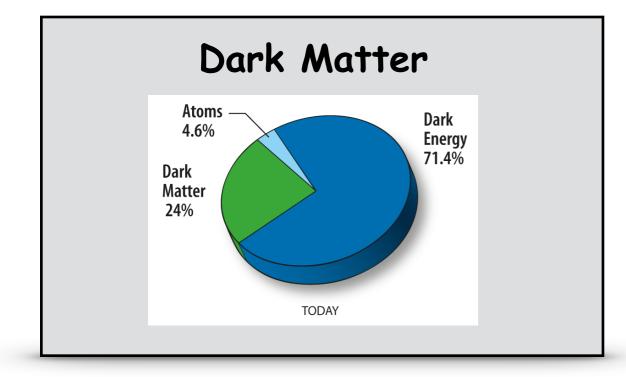


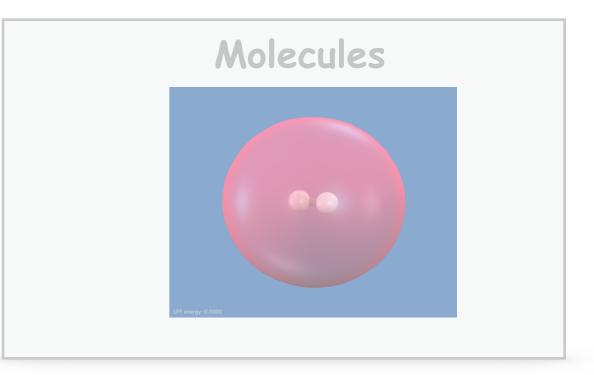


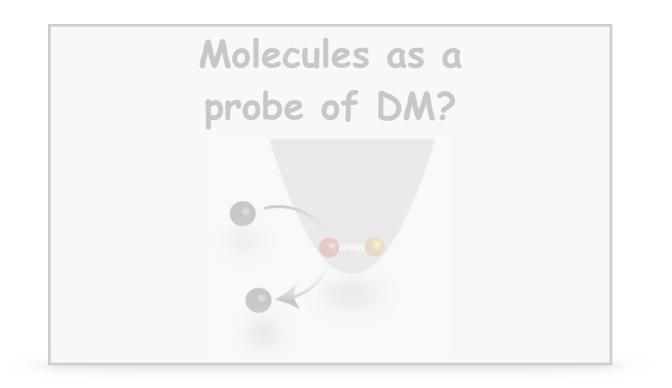






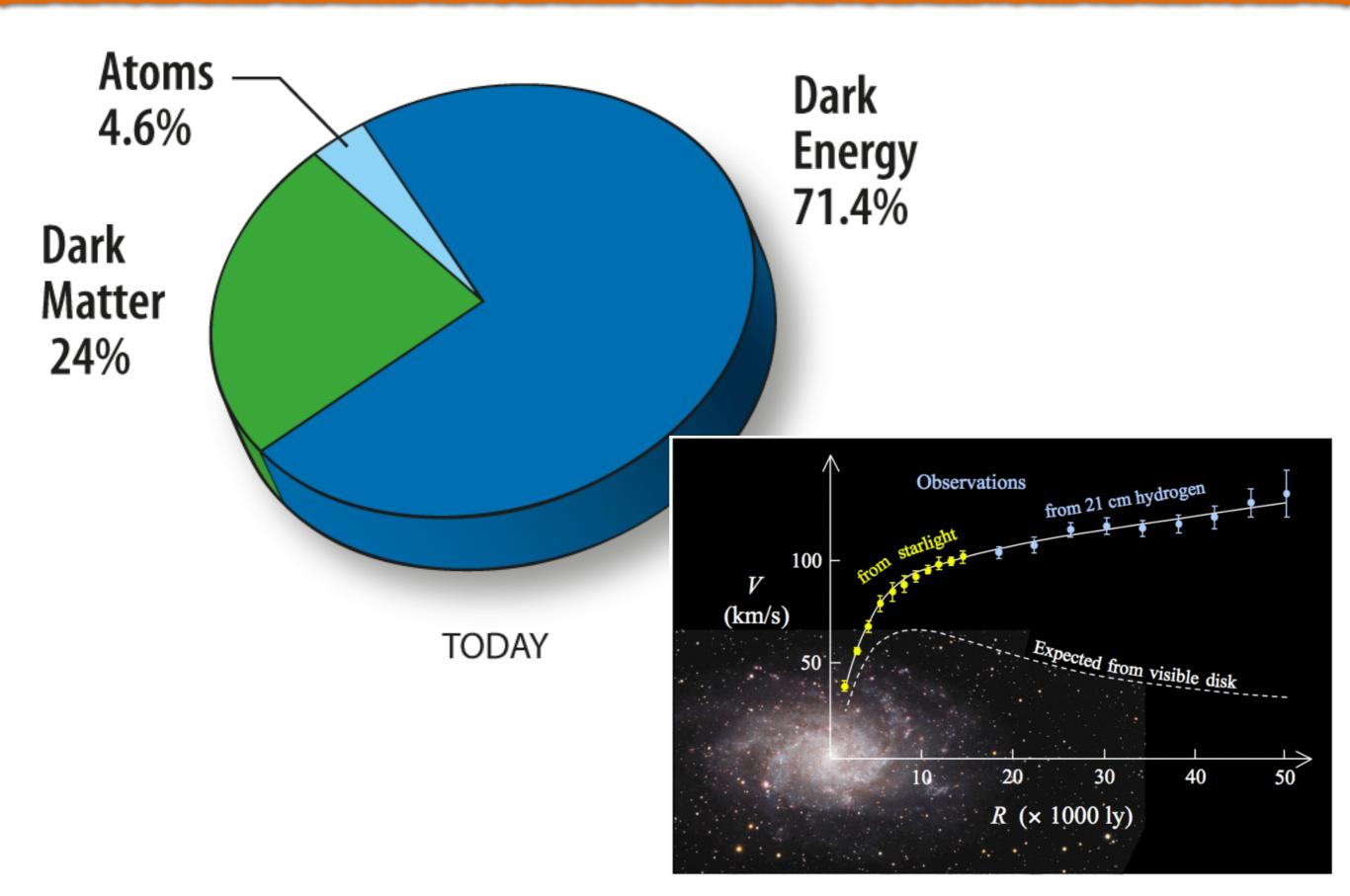








Dark Matter



Dark Matter

Dark Matter Candidates from Particle Physics and Methods of Detection

Jonathan L. Feng

Department of Physics and Astronomy, University of California, Irvine, California 92697; email: jlf[at]uci.edu

Annu. Rev. Astron. Astrophys. 2010. 48:495-545

	WIMPs	SuperWIMPs	Light \tilde{G}	Hidden DM	Sterile v	Axions
Motivation	GHP	GHP	GHP/NPFP	GHP/NPFP	v Mass	Strong CP
Naturally Correct Ω	Yes	Yes	No	Possible	No	No
Production Mechanism	Freeze Out	Decay	Thermal	Various	Various	Various
Mass Range	GeV-TeV	GeV-TeV	eV-keV	GeV-TeV	keV	µeV-meV
Temperature	Cold	Cold/Warm	Cold/Warm	Cold/Warm	Warm	Cold
Collisional				\checkmark		
Early Universe		$\sqrt{}$		\checkmark		
Direct Detection	$\sqrt{}$			\checkmark		$\sqrt{}$
Indirect Detection	$\sqrt{}$	\checkmark		\checkmark	$\sqrt{}$	
Particle Colliders	$\sqrt{}$	$\sqrt{}$	$\sqrt{}$	\checkmark		

Table 1 Summary of dark matter particle candidates, their properties, and their potential methods of detection

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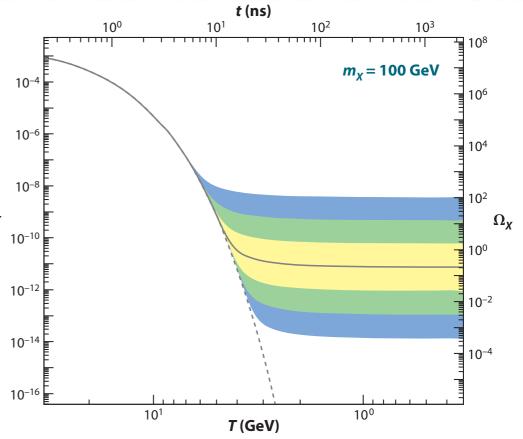


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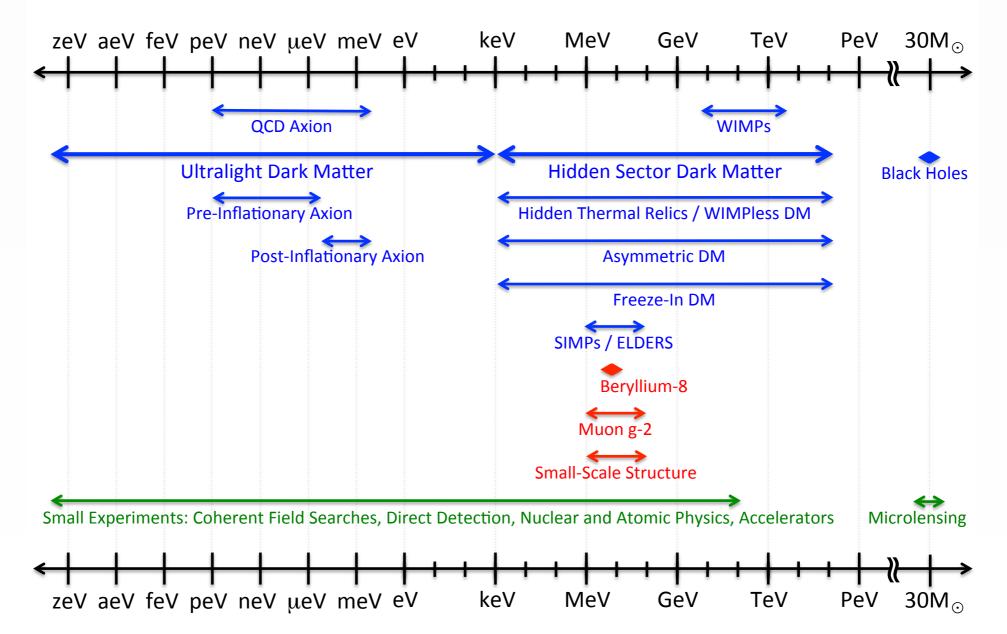
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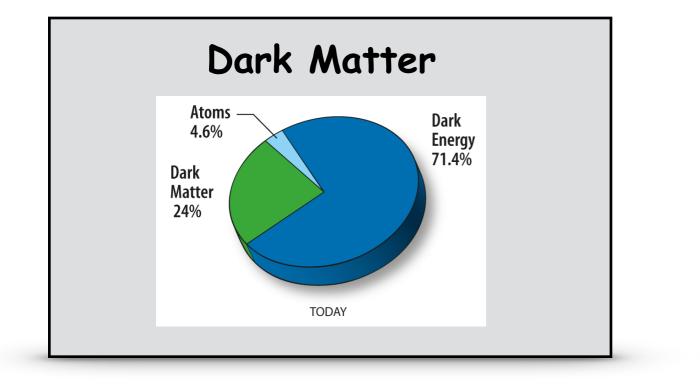


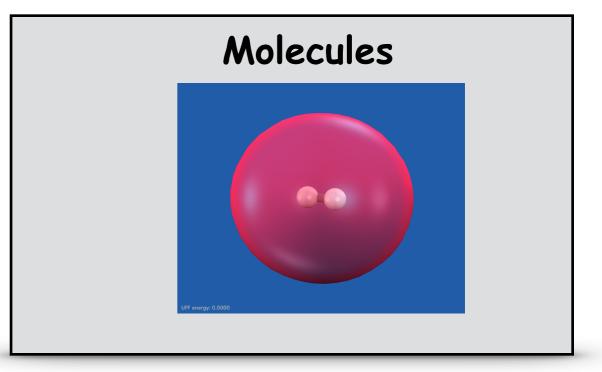
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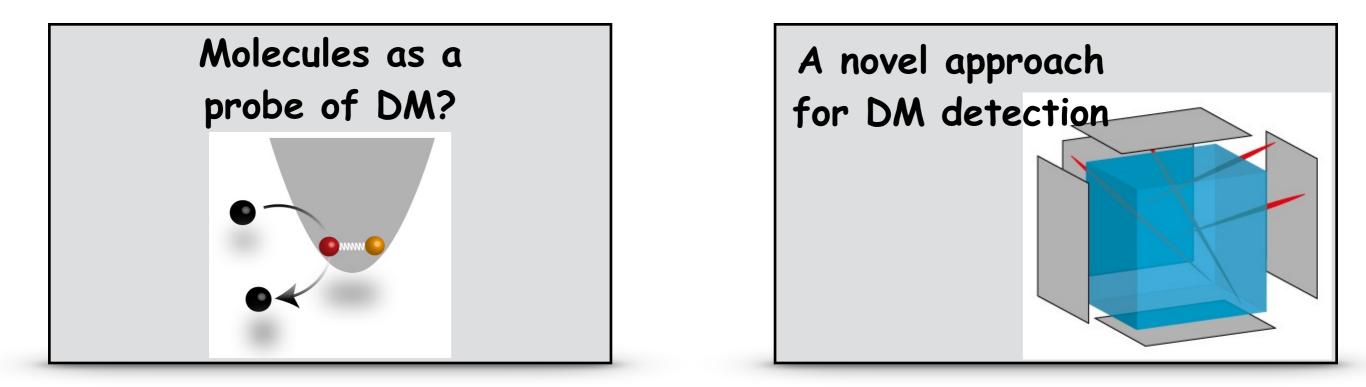
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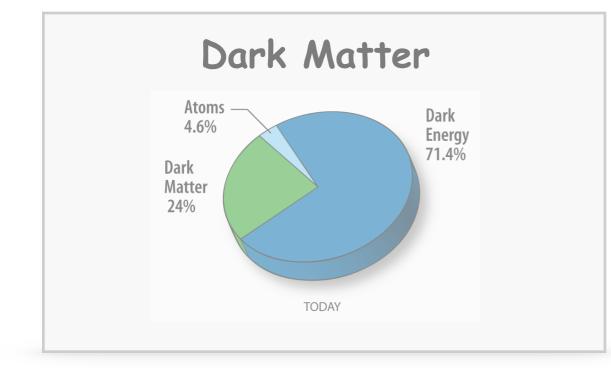


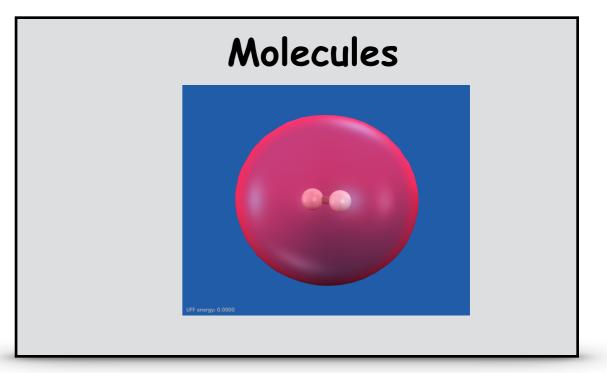


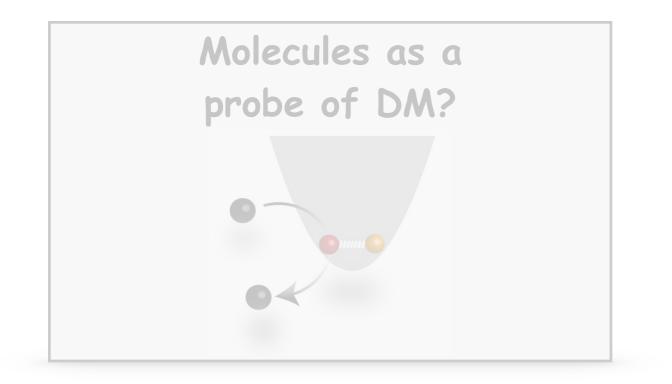
















$H = H_{\rm el} + T_N$

 $H_{\rm el} = T_e + V_{ee} + V_{NN} + V_{Ne}$



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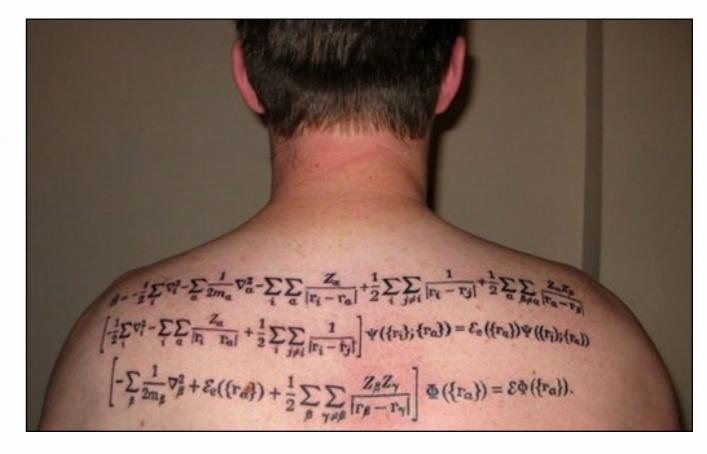
$$\Psi(\vec{r},\vec{R}) = \phi_{el}(\vec{r};\vec{R})\phi_N(\vec{R})$$



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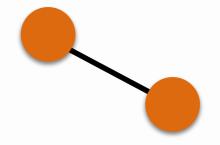


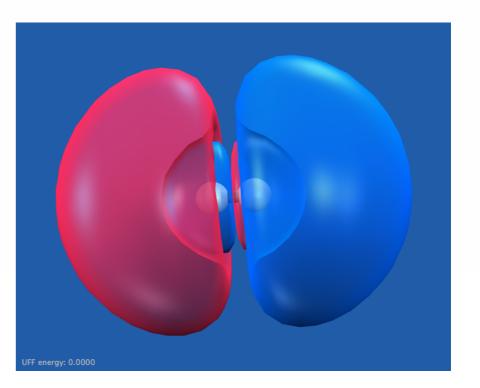


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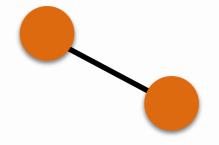


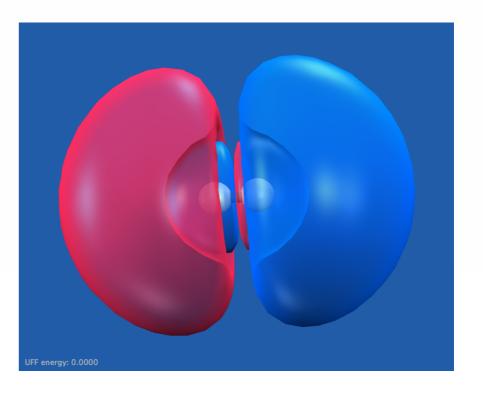
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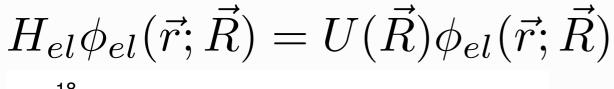
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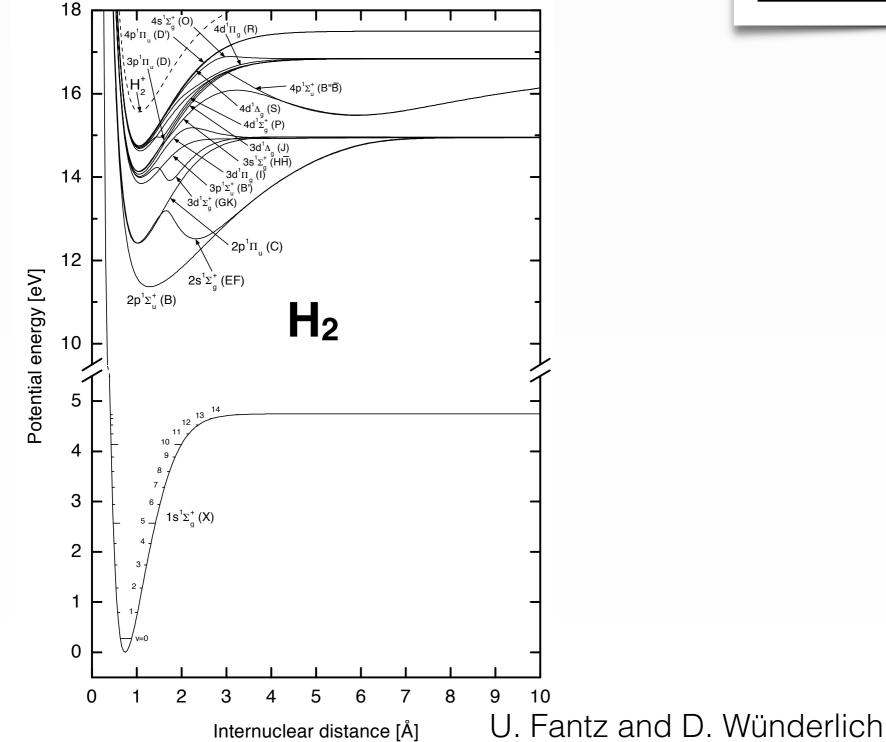
$$H_{el}\phi_{el}(\vec{r};\vec{R}) = U(\vec{R})\phi_{el}(\vec{r};\vec{R})$$

$$H = H_{\rm el} + T_N$$

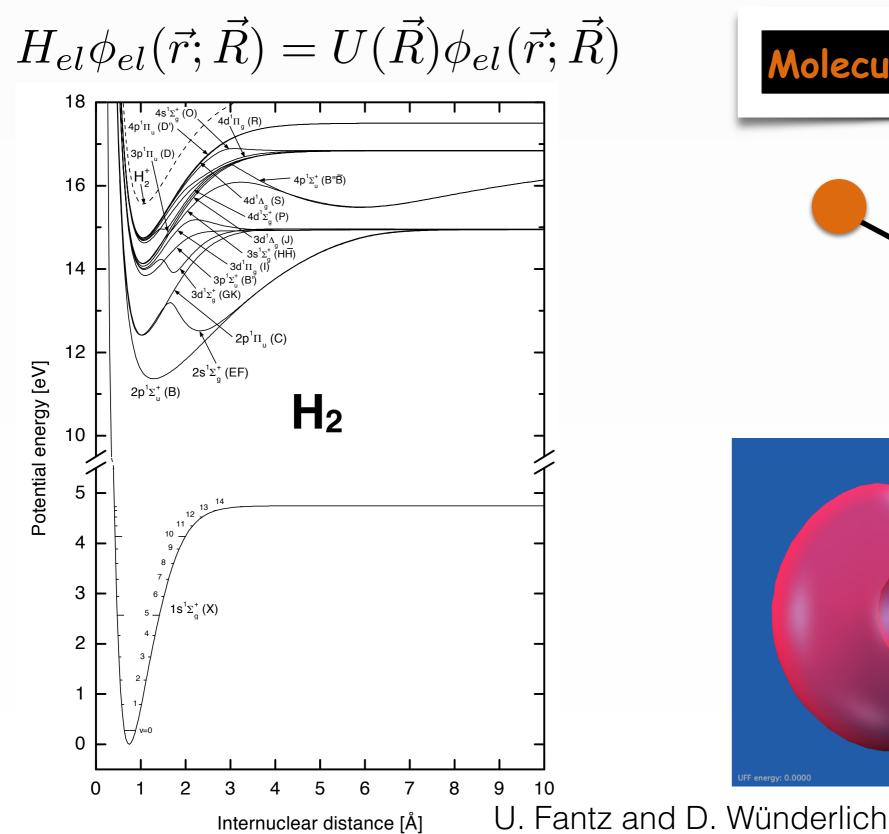


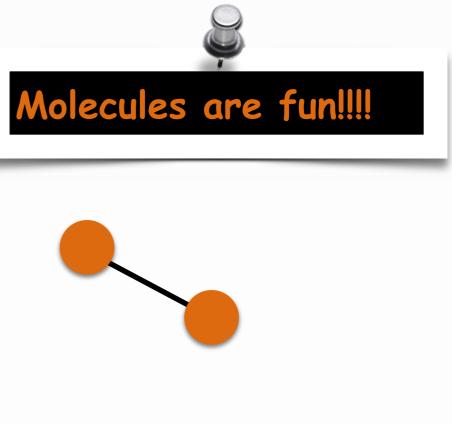


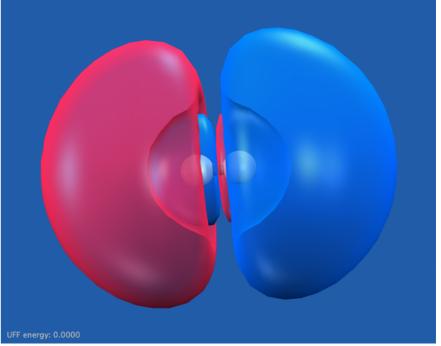








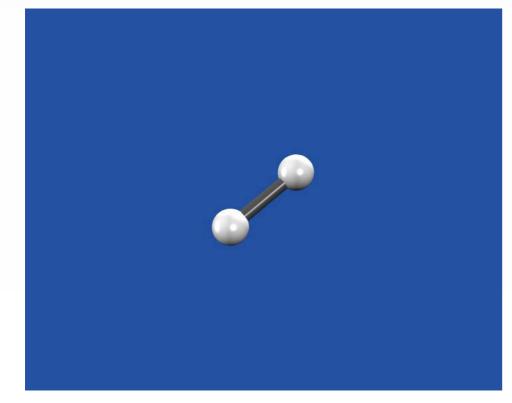


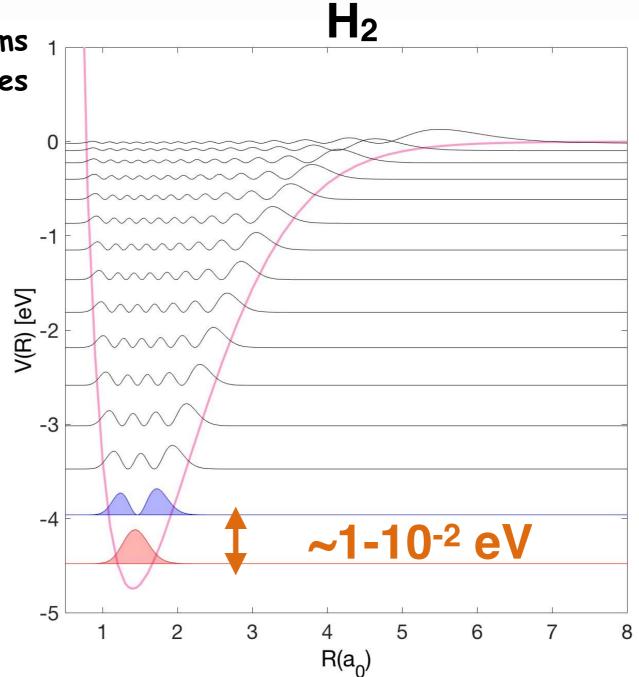




Molecules are more versatile than atoms thanks to the presence of internal degrees of freedom:

- \cdot Electronic
- \cdot Vibrational
- \cdot Rotational

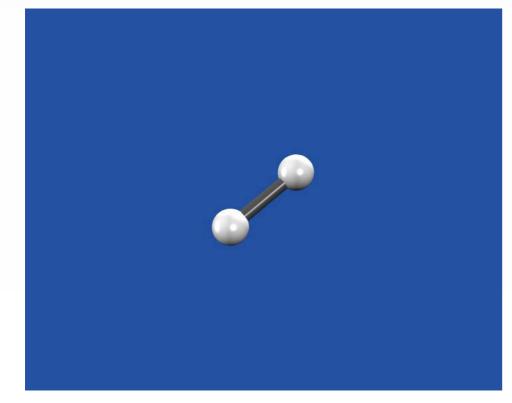


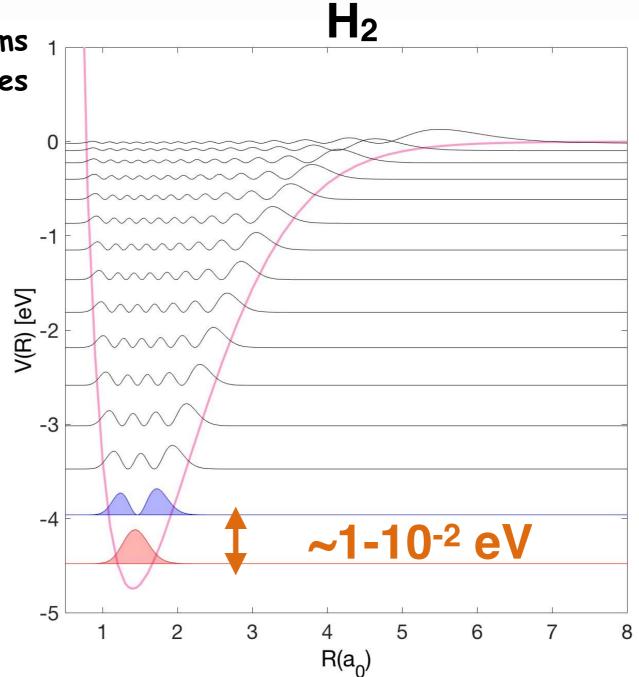


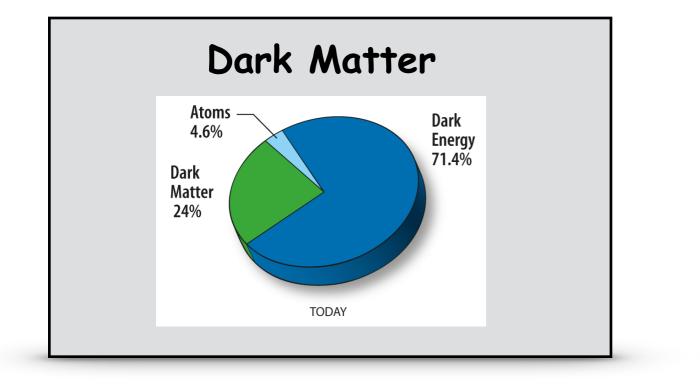


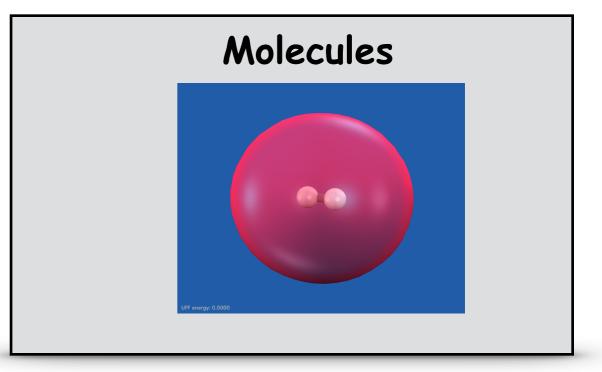
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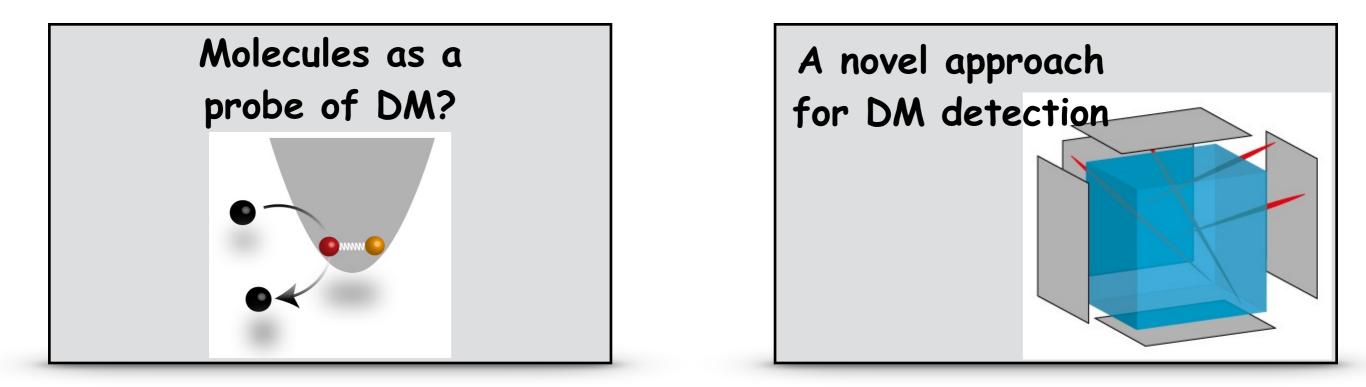
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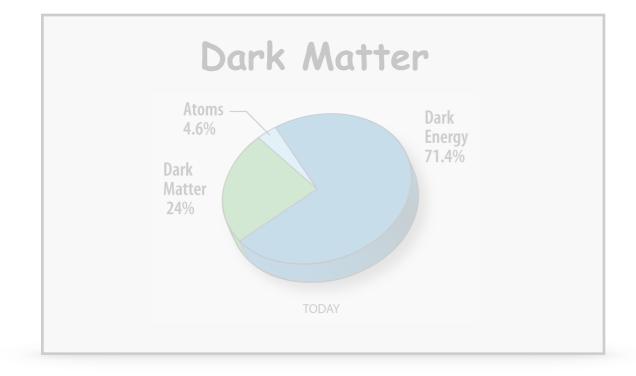




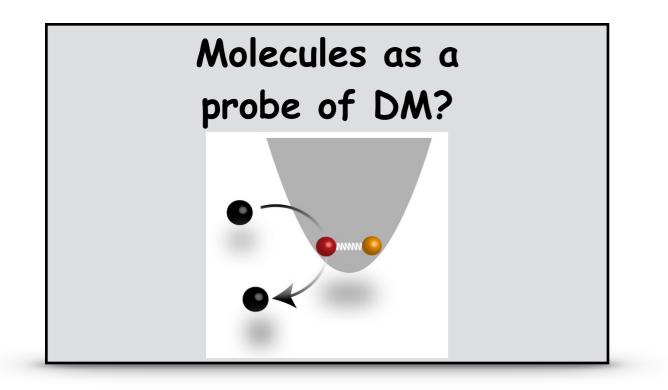














PHYSICAL REVIEW D 85, 076007 (2012)

Direct detection of sub-GeV dark matter

Rouven Essig,¹ Jeremy Mardon,^{2,3,4} and Tomer Volansky^{2,3}

¹SLAC National Accelerator Laboratory, Stanford University, Menlo Park, California 94025, USA ²Berkeley Center for Theoretical Physics, Department of Physics, University of California, Berkeley, California 94720, USA ³Lawrence Berkeley National Laboratory, Berkeley, California 94720, USA ⁴Stanford Institute for Theoretical Physics, Department of Physics, Stanford University, Stanford, California 94305, USA (Received 2 October 2011; published 9 April 2012)

$E_{tot} = m_{DM}v^2/2 = 50 \text{ eV x} (m_{DM}/100 \text{ MeV})$

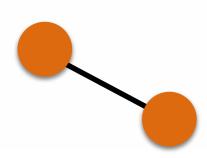
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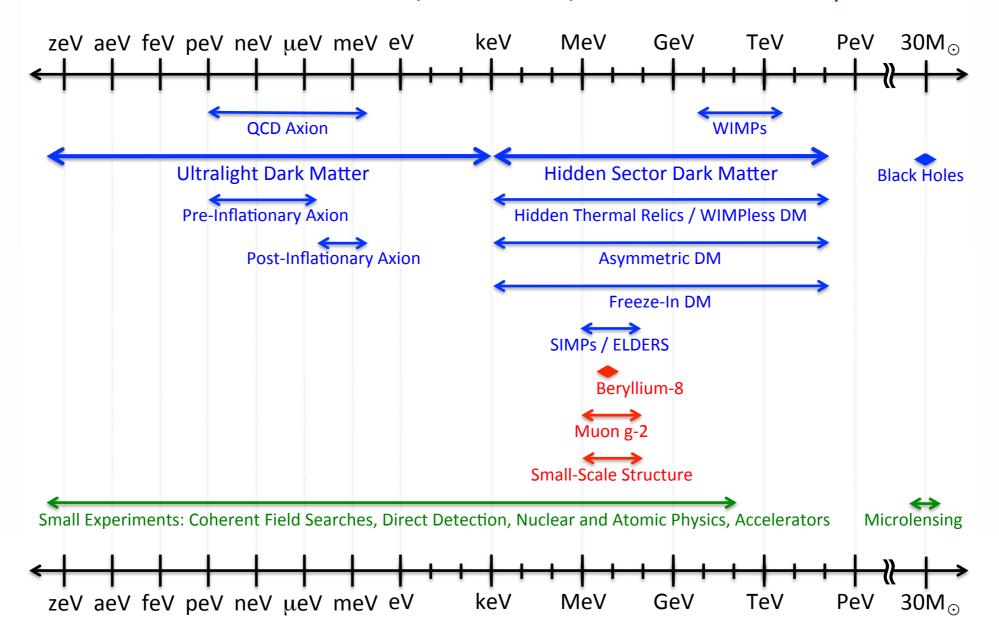


Vibration of molecules m_{DM} ~ MeV

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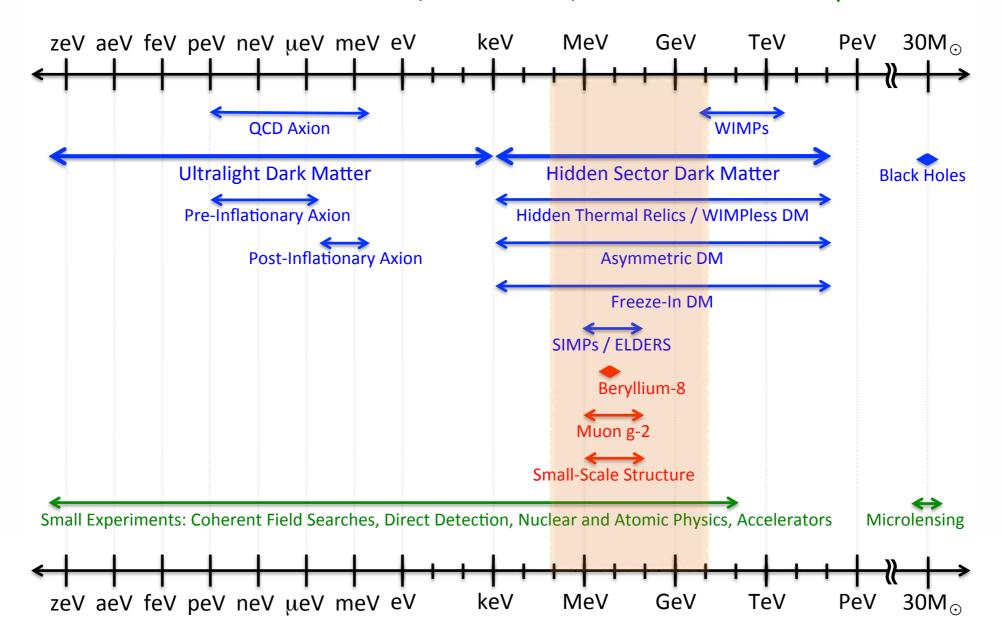
Dark Sector Candidates, Anomalies, and Search Techniques



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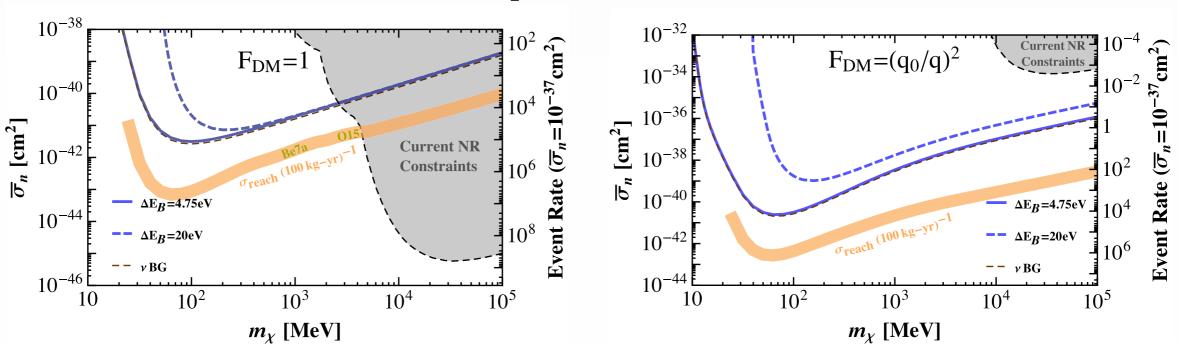
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PHYSICAL REVIEW D 95, 056011 (2017)

Detection of sub-GeV dark matter and solar neutrinos via chemical-bond breaking

Rouven Essig,^{1,*} Jeremy Mardon,^{2,†} Oren Slone,^{3,‡} and Tomer Volansky^{3,§} ¹C.N. Yang Institute for Theoretical Physics, Stony Brook University, Stony Brook, New York 11794, USA ²Stanford Institute for Theoretical Physics, Department of Physics, Stanford University, Stanford, California 94305, USA ³Raymond and Beverly Sackler School of Physics and Astronomy, Tel-Aviv University, Tel-Aviv 69978, Israel (Received 11 December 2016; published 8 March 2017)



H₂-like Molecule

Direct Detection of Light Dark Matter and Solar Neutrinos via Color Center Production in Crystals

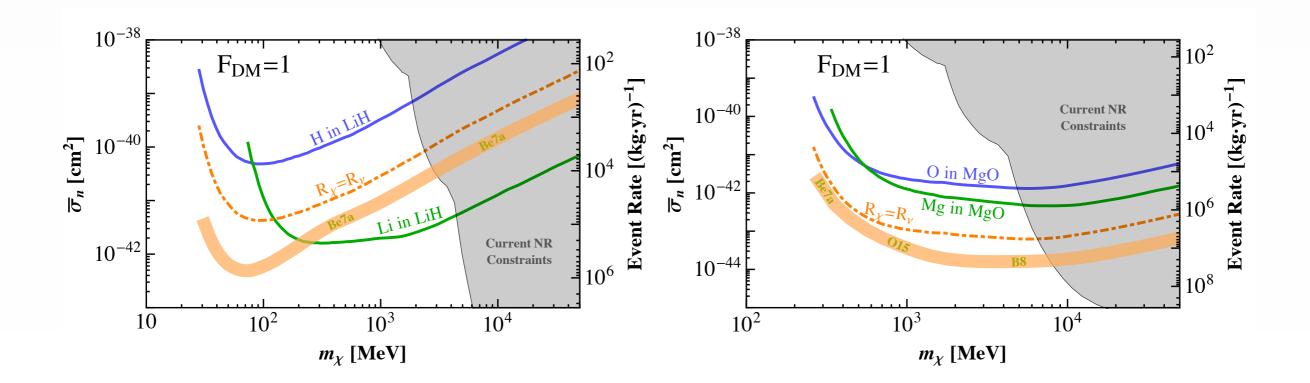
Ranny Budnik

Department of Particle Physics and Astrophysics, Weizmann Institute of Science, Rehovot, Israel*

Ori Chesnovsky Raymond and Beverly Sackler School of Chemistry, Tel-Aviv University, Tel-Aviv, Israel[†]

Oren Slone and Tomer Volansky

Raymond and Beverly Sackler School of Physics and Astronomy, Tel-Aviv University, Tel-Aviv, Israel[‡]

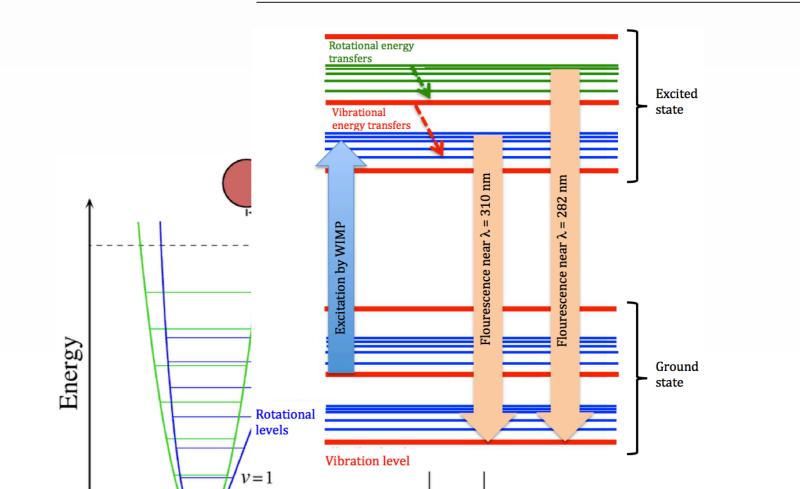


Molecular excitations: a new way to detect Dark matter

J.Va'vra

SLAC, Stanford University, CA94309, U.S.A. e-mail: jjv@slac.stanford.edu **Table 1.** A simple calculation of the transition wavelength for several frequency overtones of the OH-radicals. The last two modes correspond to visible wavelengths. Higher modes can reach the UV regime [6].

OH-band identity	Transition	Calculated wavelength [nm]
ν_1	$0 \rightarrow 1$	2803
$2v_1$	$0 \rightarrow 2$	1436
$3v_1$	$0 \rightarrow 3$	980
$4v_1$	$0 \rightarrow 4$	755
$5v_1$	$0 \rightarrow 5$	619.5



Resonant absorption of bosonic dark matter in molecules

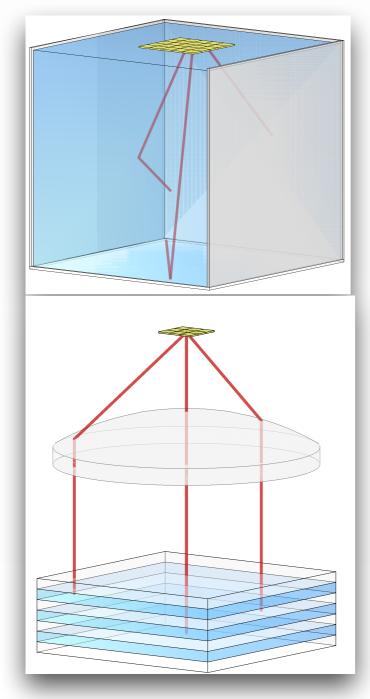
Asimina Arvanitaki,^{1, *} Savas Dimopoulos,^{2, †} and Ken Van Tilburg^{3, 4, ‡}

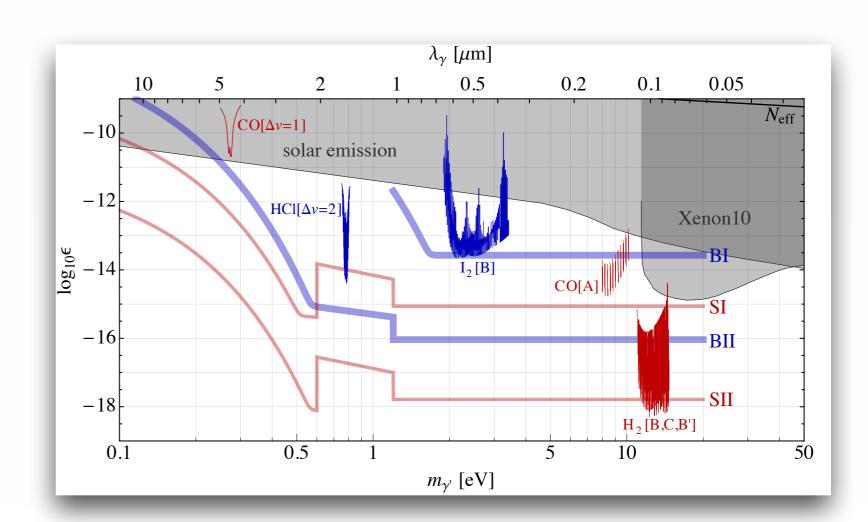
¹Perimeter Institute for Theoretical Physics, Waterloo, Ontario N2L 2Y5, Canada

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³School of Natural Sciences, Institute for Advanced Study, Princeton, NJ 08540, USA

⁴Center for Cosmology and Particle Physics, Department of Physics, New York University, New York, NY 10003



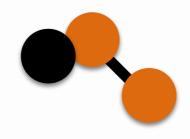






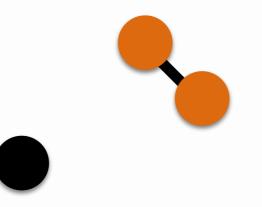
$$q_{\min/\max}^{2} = \mu_{\chi N}^{2} v_{max}^{2} [1 \mp \sqrt{1 - \frac{2\Delta E}{\mu_{\chi N} v_{max}^{2}}}]^{2}$$





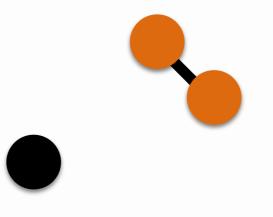
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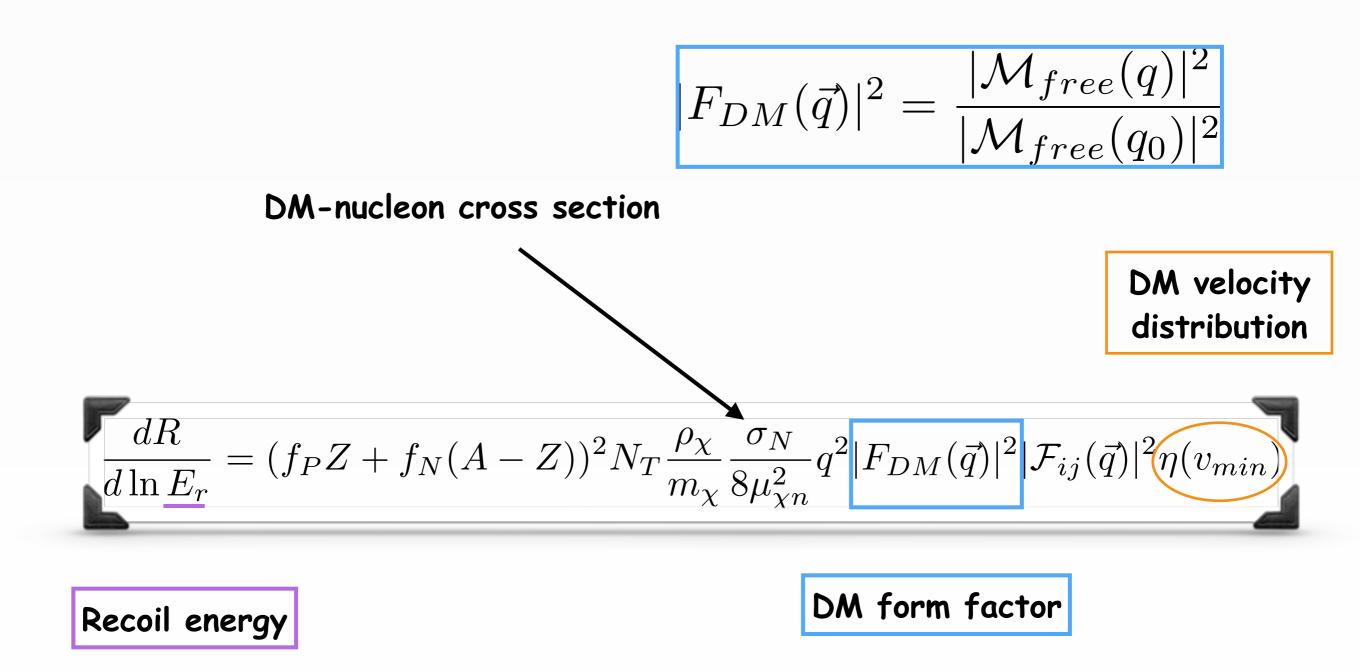
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$v_{\rm max} = v_{\rm SS} + v_{\rm esc}$

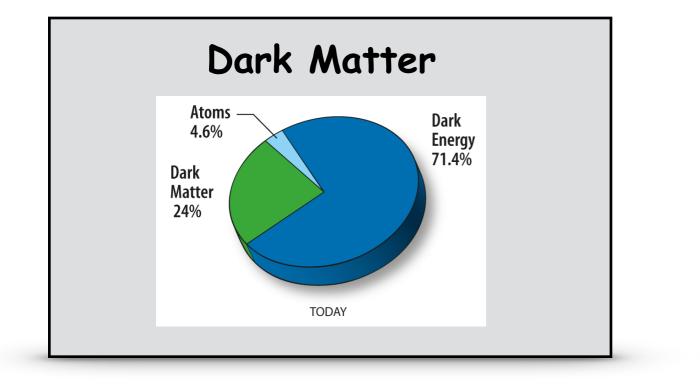
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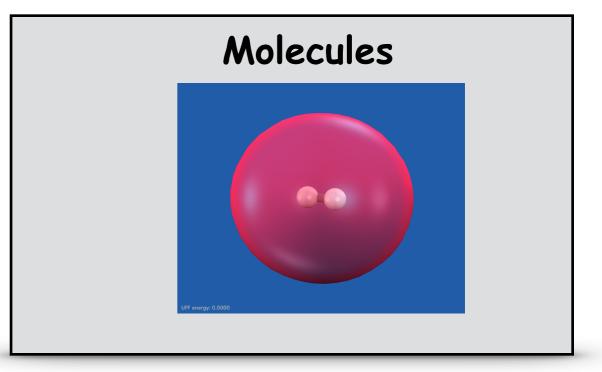


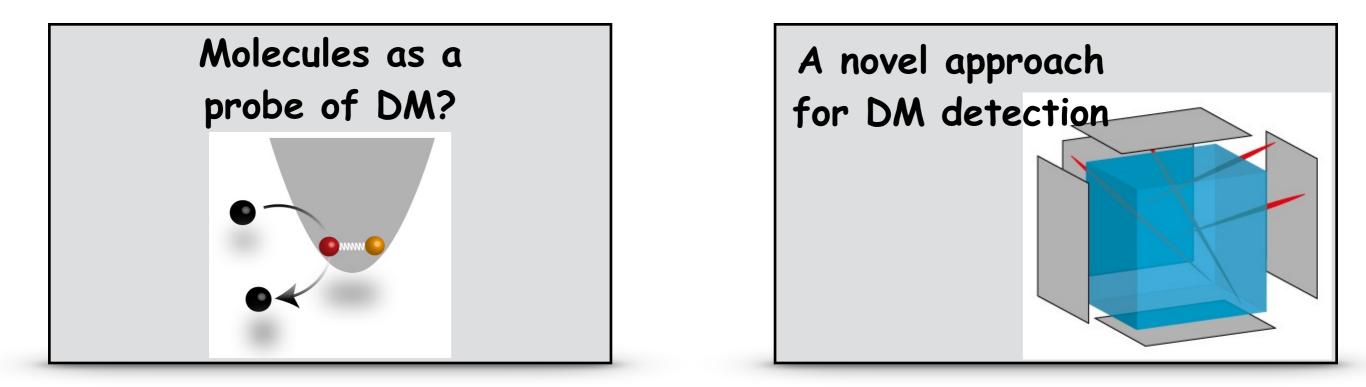
$$\frac{dR}{d\ln E_r} = (f_P Z + f_N (A - Z))^2 N_T \frac{\rho_\chi}{m_\chi} \frac{\sigma_N}{8\mu_{\chi n}^2} q^2 |F_{DM}(\vec{q})|^2 |\mathcal{F}_{ij}(\vec{q})|^2 \eta(v_{min})$$

Molecular form factor

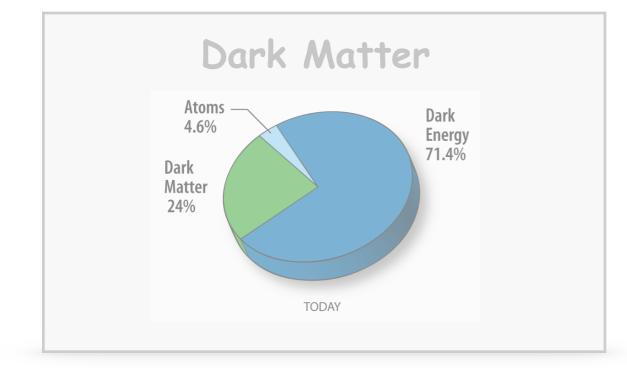
$$|\mathcal{F}_{ij}(q)|^2 = |\int d^3 \vec{r} e^{\frac{i\mu_{12}}{m_A}\vec{q}.\vec{r}} \Psi_i^J(\vec{r}) \Psi_j^{J'}(\vec{r})|^2$$



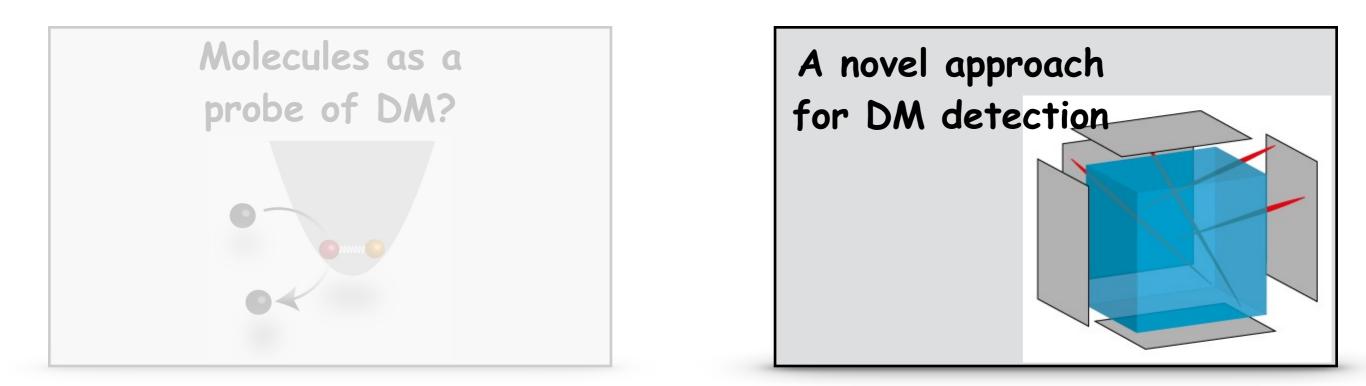


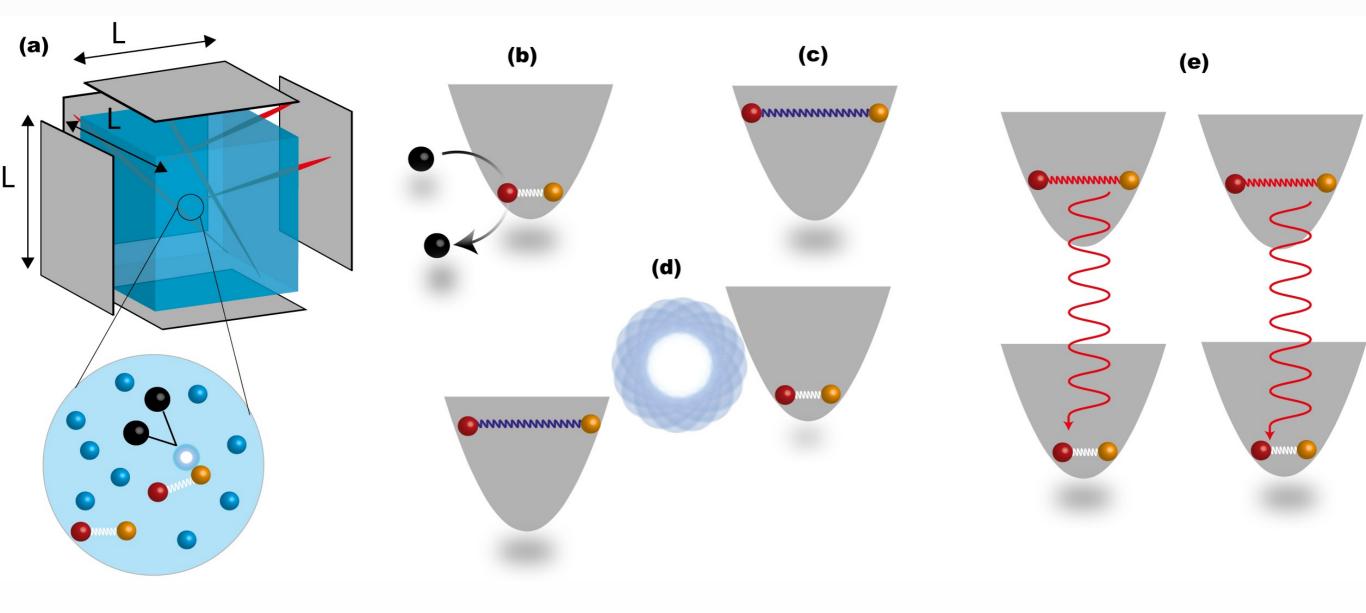


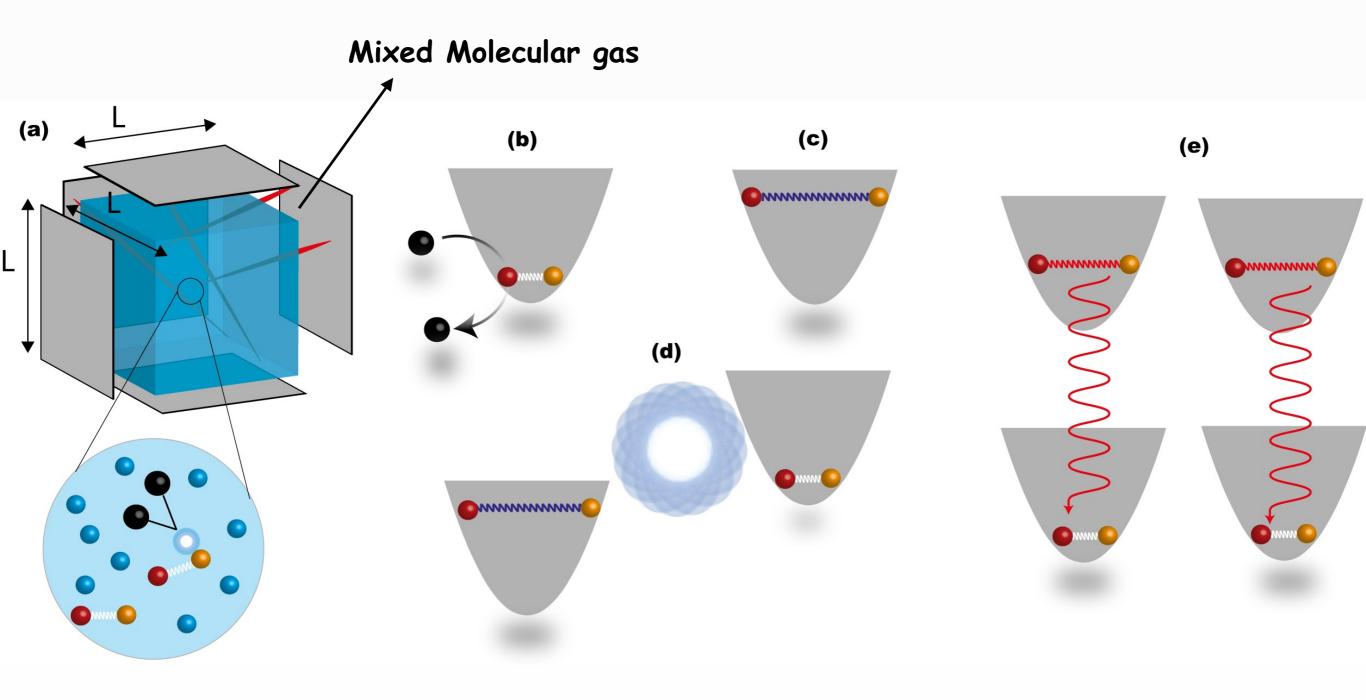
Outlook

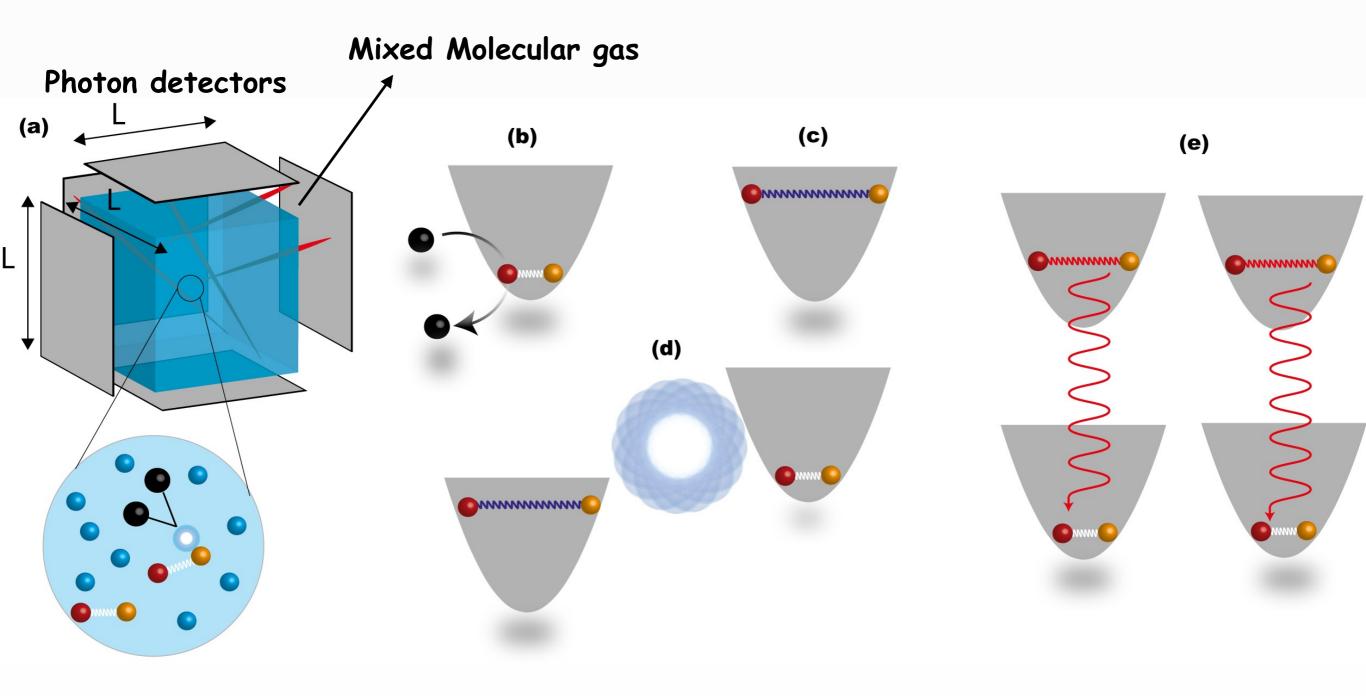


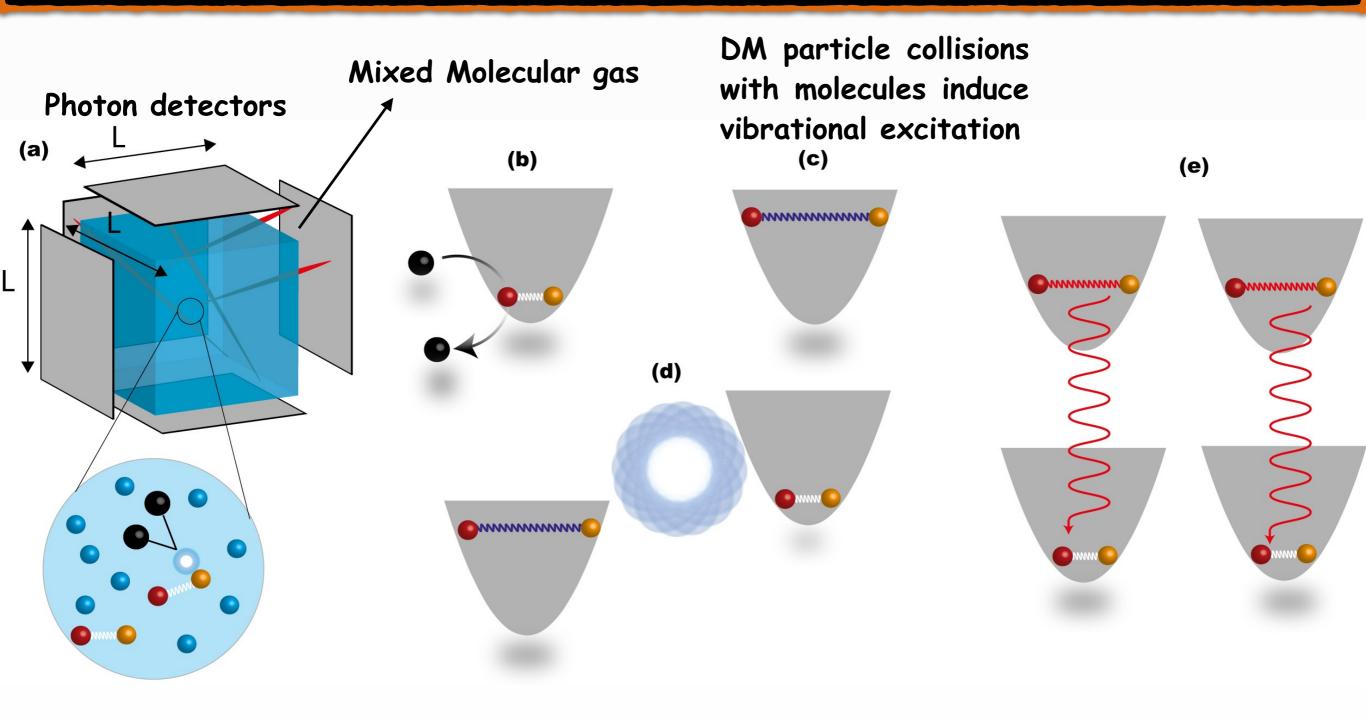


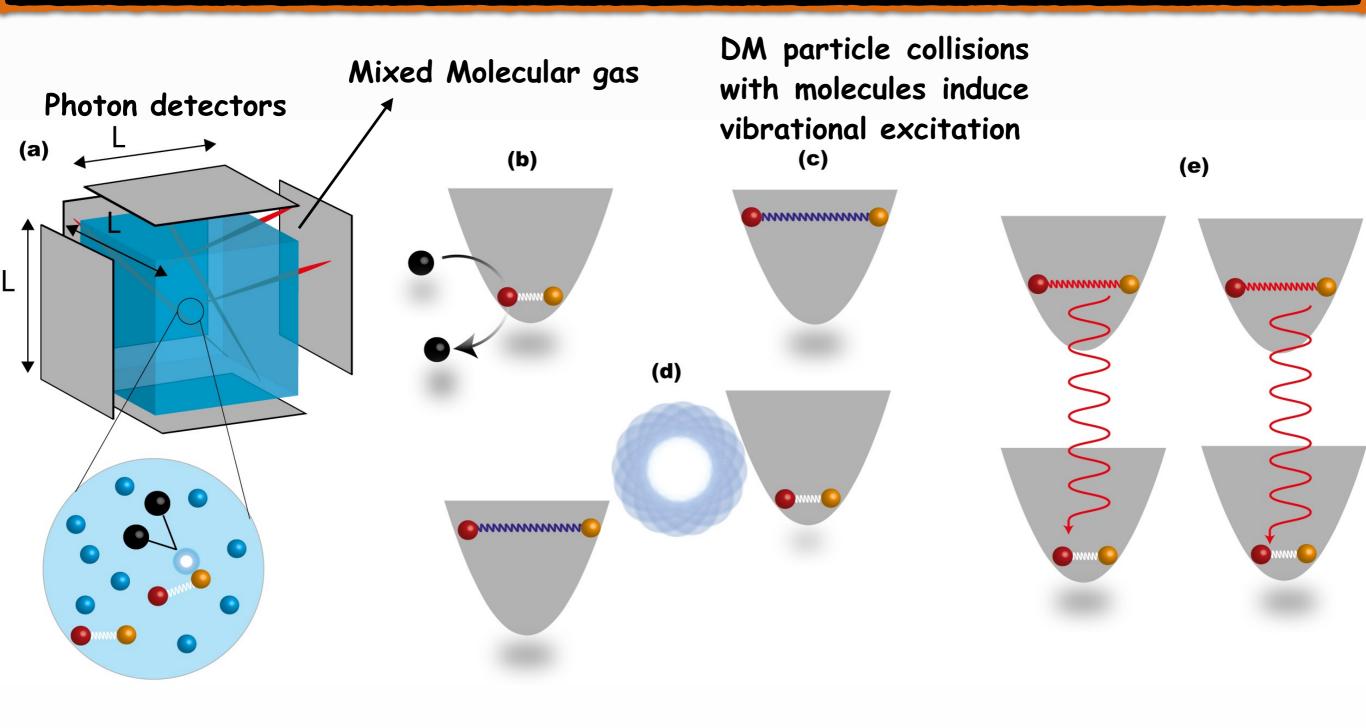




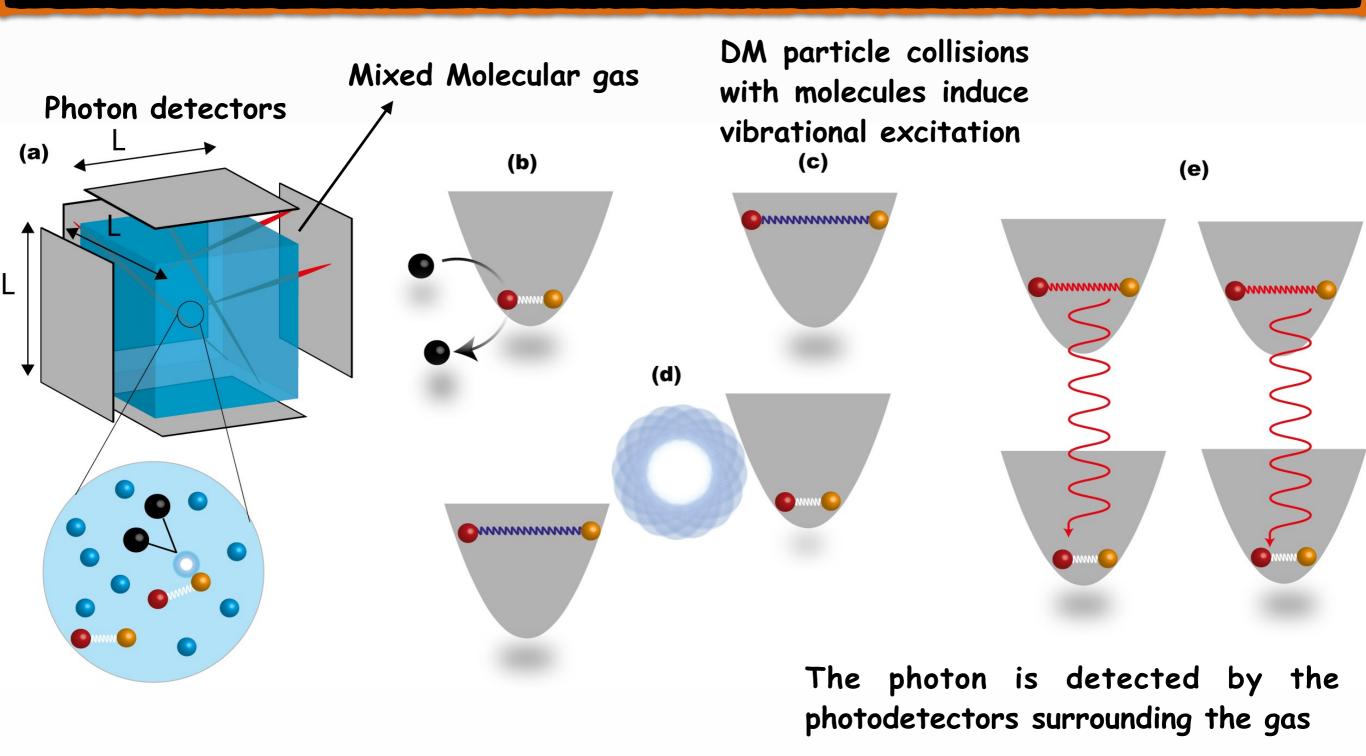




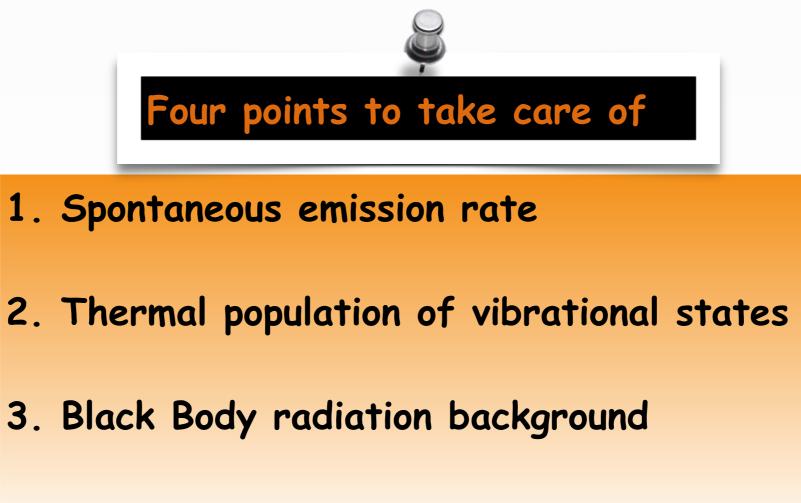




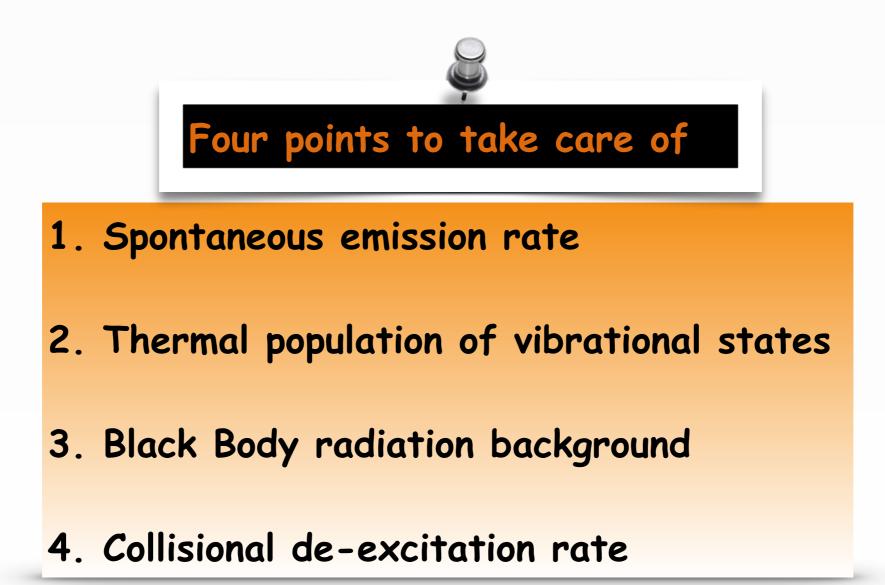
Excited molecules decay emitting a photon



Excited molecules decay emitting a photon

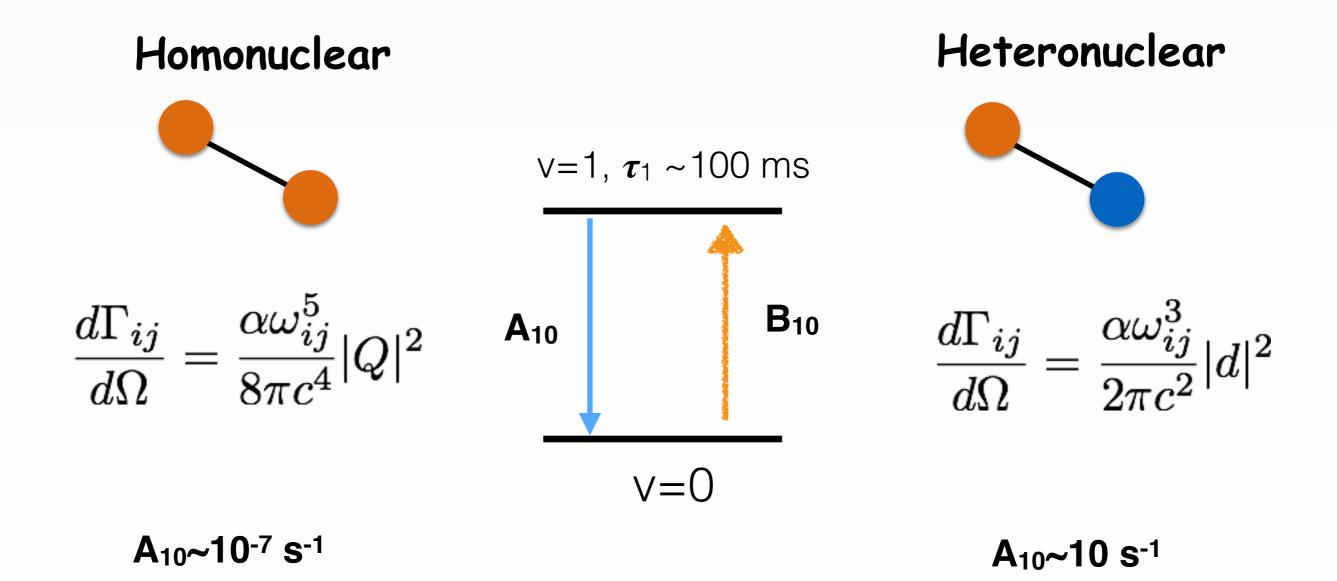


4. Collisional de-excitation rate

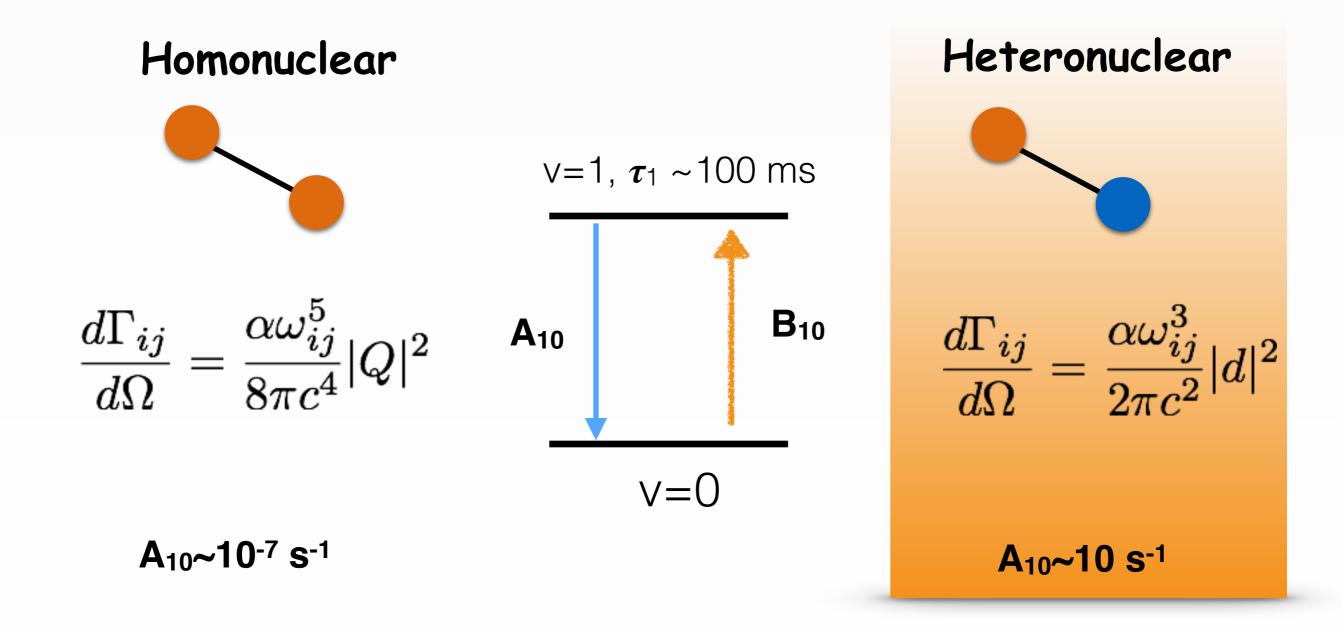


Bonus: Photon absorption

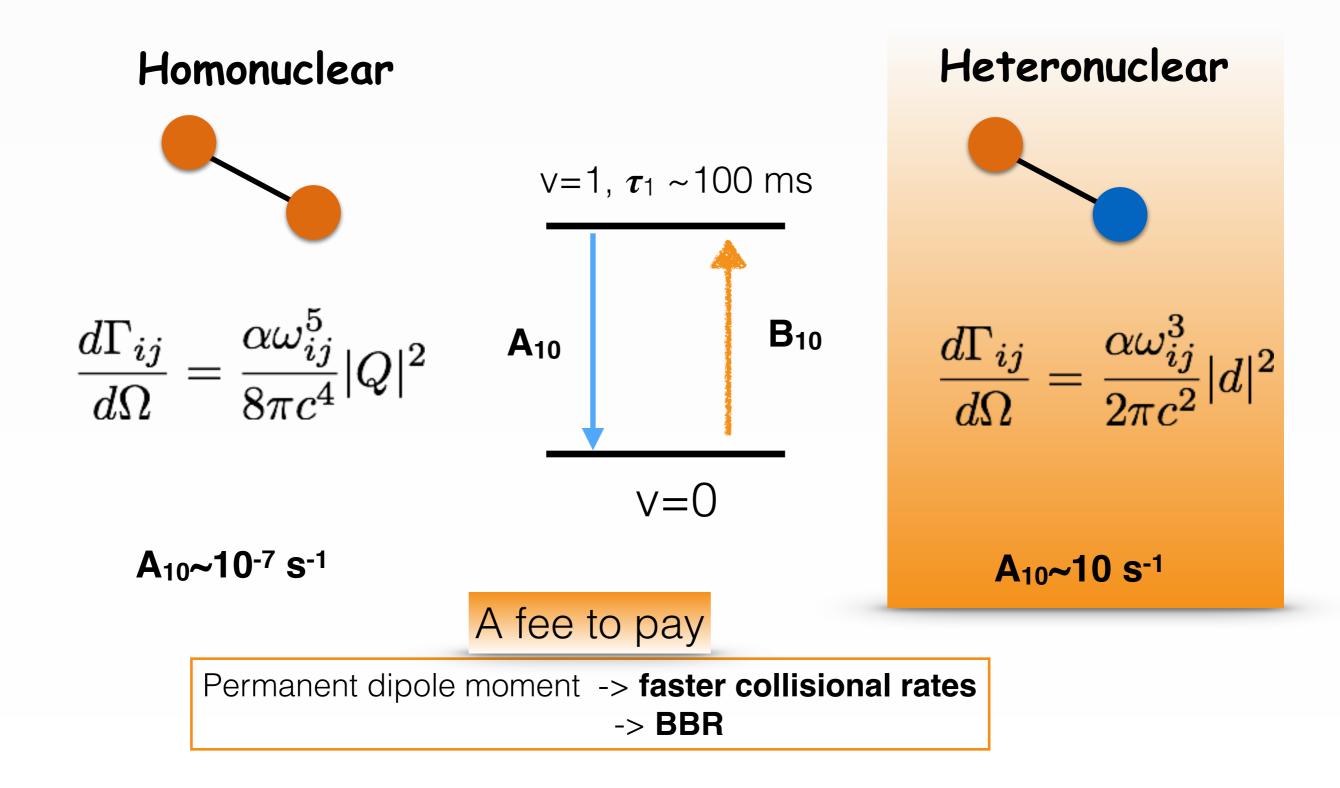
1. Spontaneous emission rate



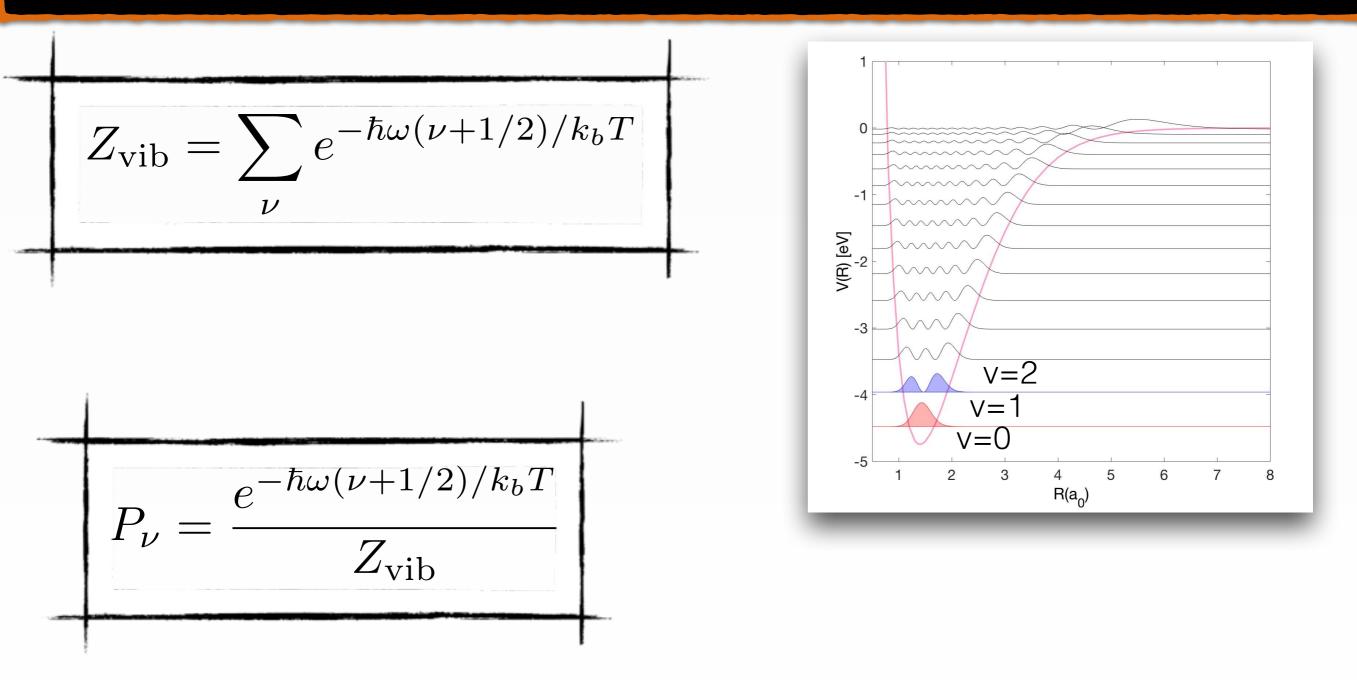
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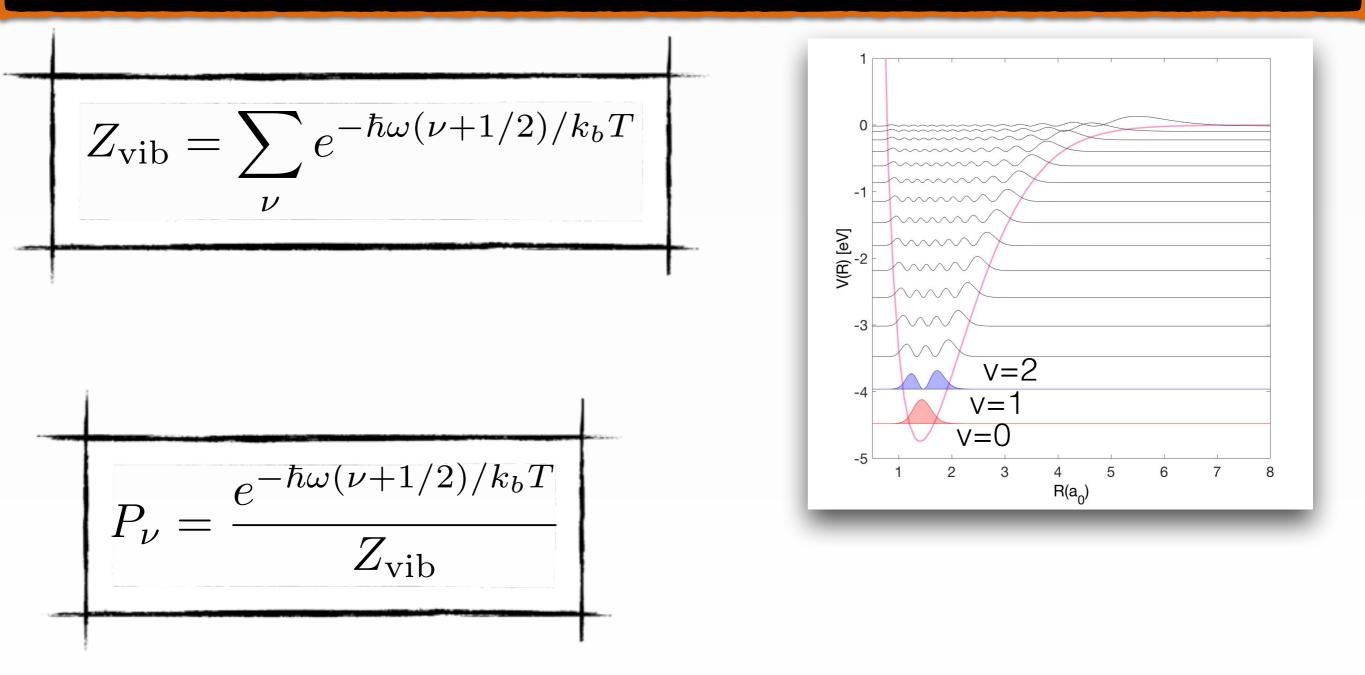
1. Spontaneous emission rate



2. Thermal population

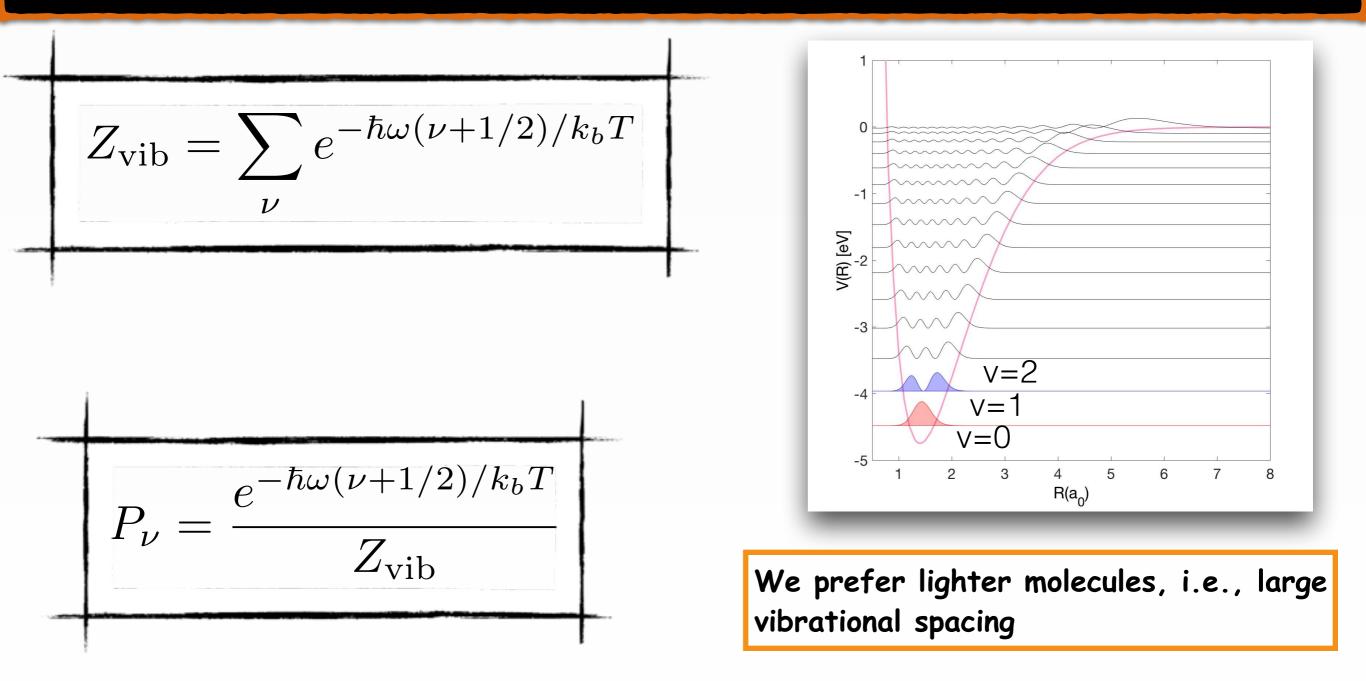


2. Thermal population

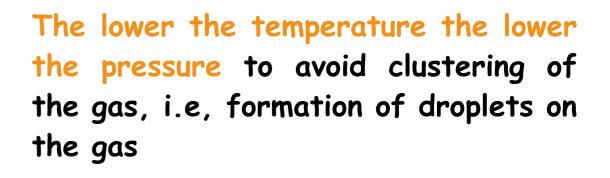


The temperature has to be low enough to ensure that almost every single molecule is in the vibrational ground state The lower the temperature the lower the pressure to avoid clustering of the gas, i.e, formation of droplets on the gas

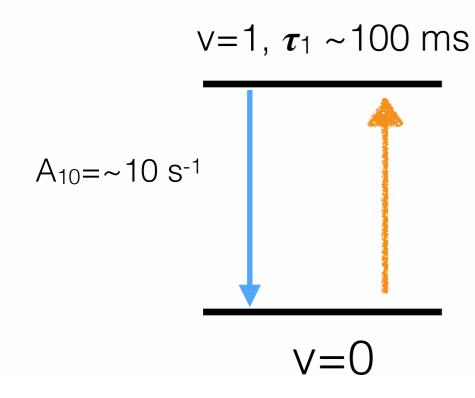
2. Thermal population



The temperature has to be low enough to ensure that almost every single molecule is in the vibrational ground state

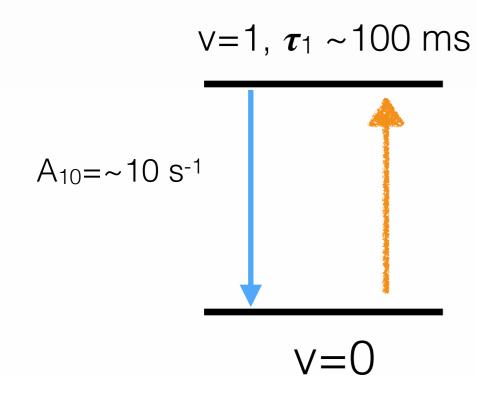


$$\Gamma_{ij}^{BBR} = \frac{8\pi^2 |d_{ij}|^2}{3\epsilon_o \hbar c^3} \frac{\nu^3}{\exp \frac{h\nu}{k_B T} - 1}$$



Dipole matrix element

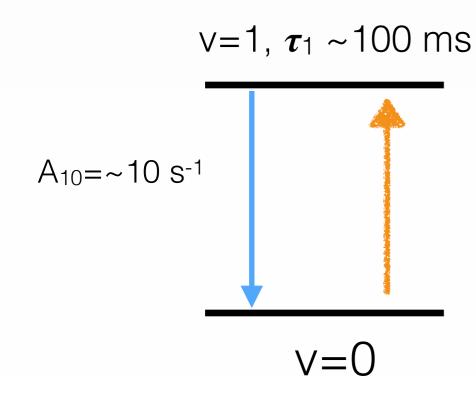
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Dipole matrix element

Vibrational frequency

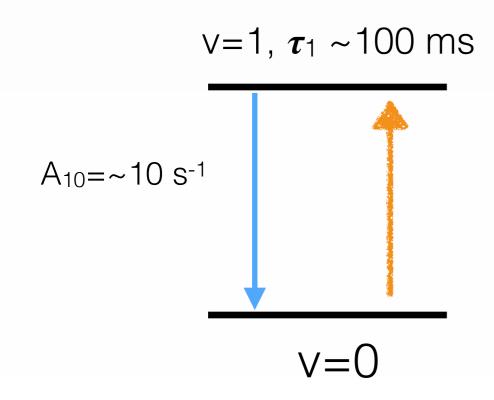
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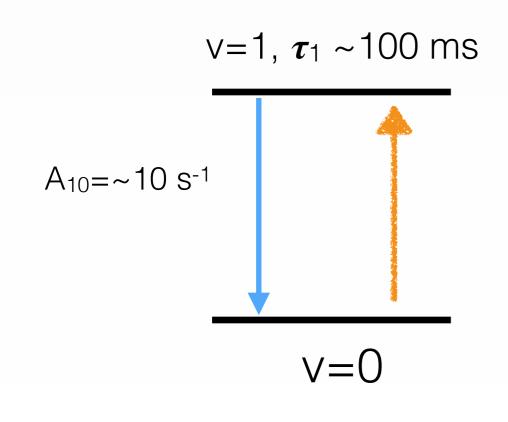


We prefer lighter molecules, i.e., large vibrational spacing

Dipole matrix element

Vibrational frequency

$$\Gamma_{ij}^{BBR} = \frac{8\pi^2 |d_{ij}|^2}{3\epsilon_o \hbar c^3} \frac{\nu^3}{\exp \frac{h\nu}{k_B T} - 1}$$



We prefer lighter molecules, i.e., large vibrational spacing

Molecules with smaller dipole moment will have smaller BBR rate

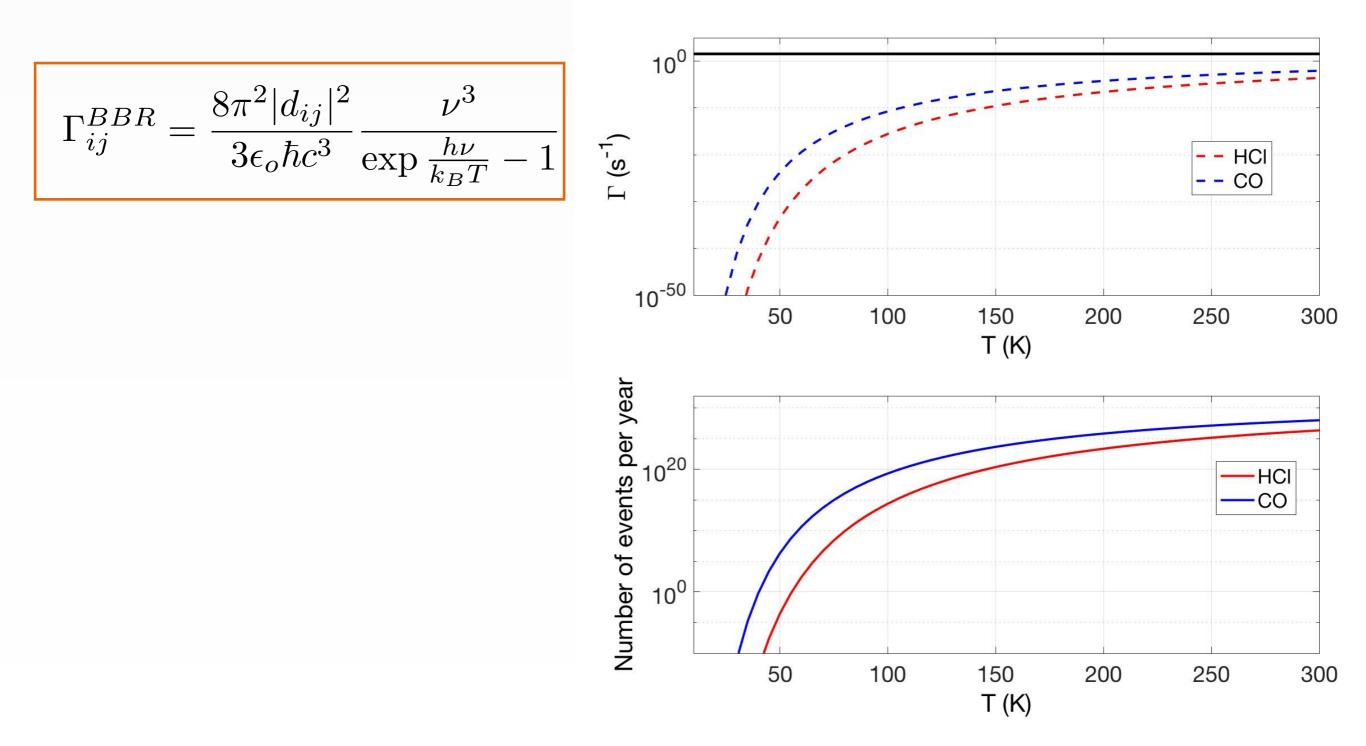
Low temperatures are better!!!

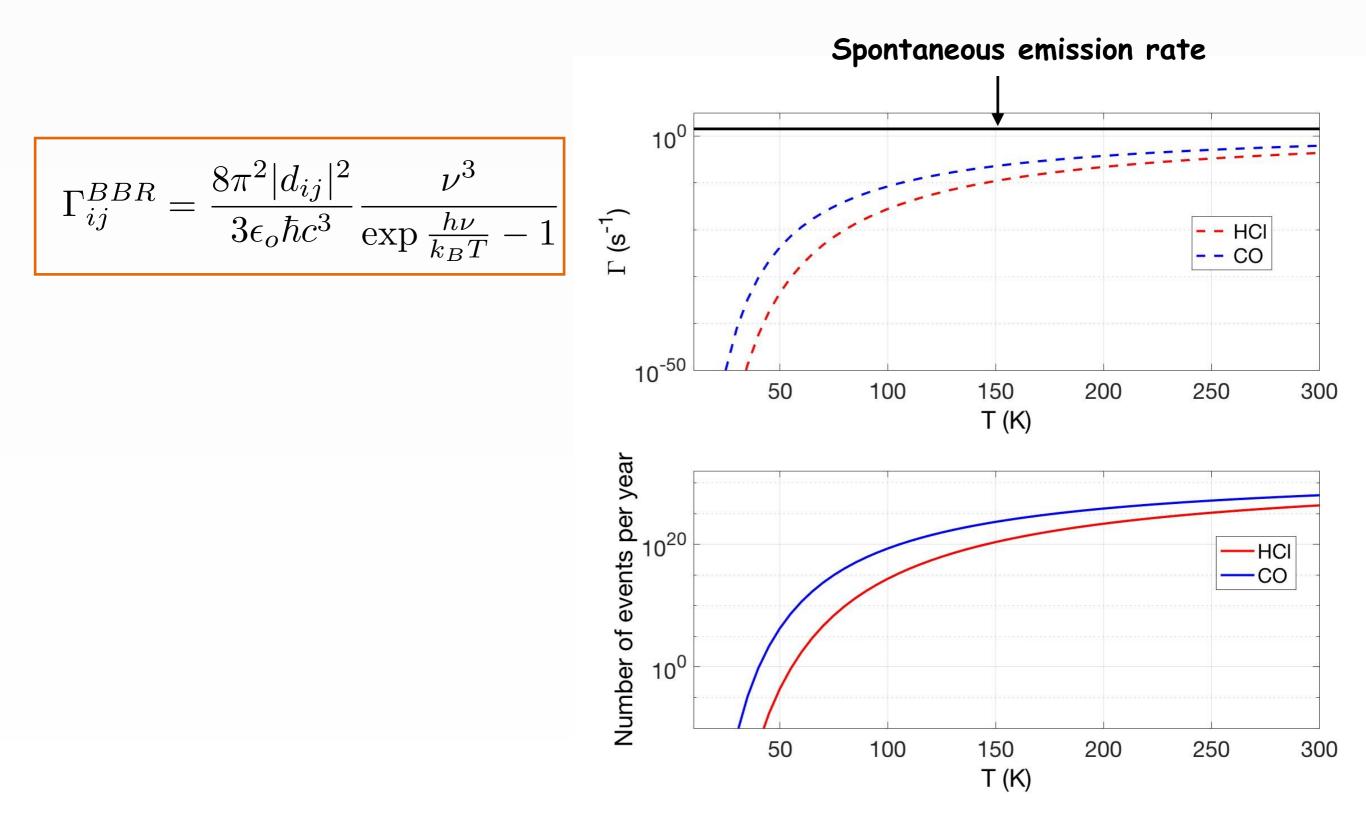
Dipole matrix element

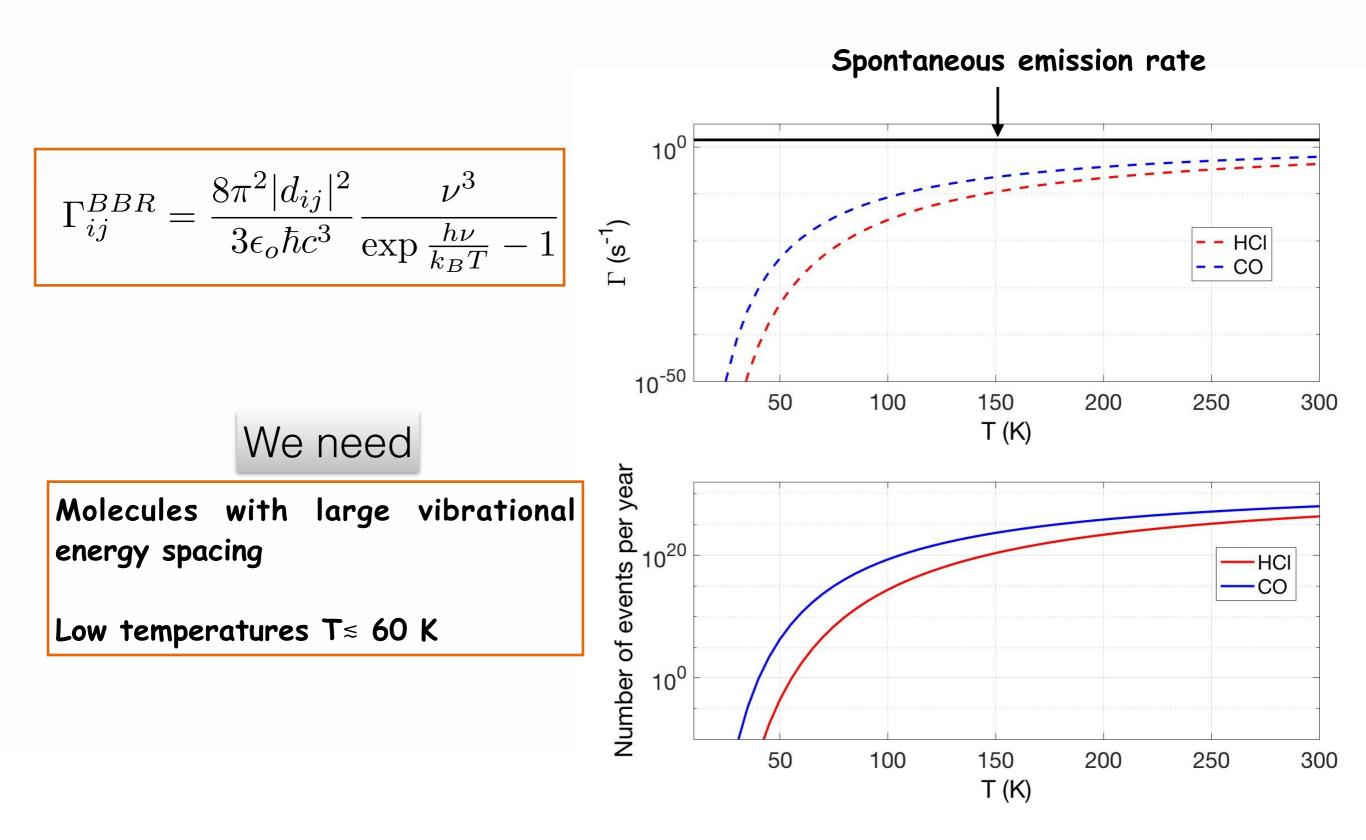
Vibrational frequency

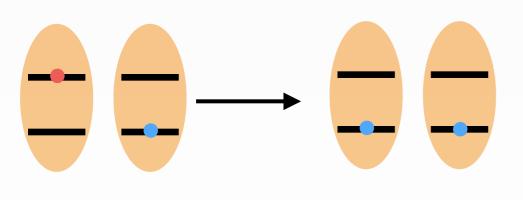
$$\Gamma_{ij}^{BBR} = \frac{8\pi^2 |d_{ij}|^2}{3\epsilon_o \hbar c^3} \frac{\nu^3}{\exp \frac{h\nu}{k_B T} - 1}$$

	Η	HF	CO	ΝΟ	HCI
r		1.7	2.1	2.2	2.4
ω	453 198 466	513	269	236	371
d(D)	1.80	1.98	0.12	0.16	1.03





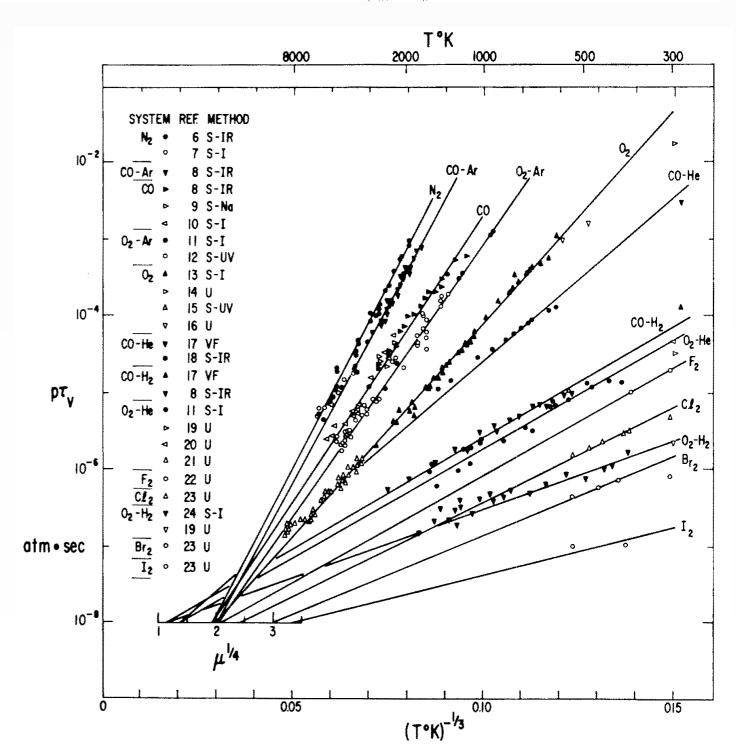


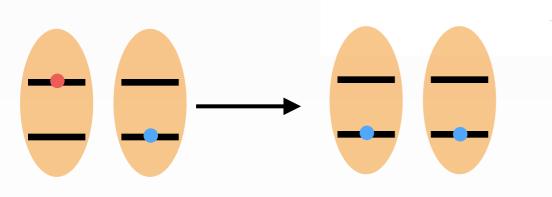


THE JOURNAL OF CHEMICAL PHYSICS VOLUME 39, NUMBER 12 15 DECEMBER 1963

Systematics of Vibrational Relaxation*

ROGER C. MILLIKAN AND DONALD R. WHITE General Electric Research Laboratory, Schenectady, New York

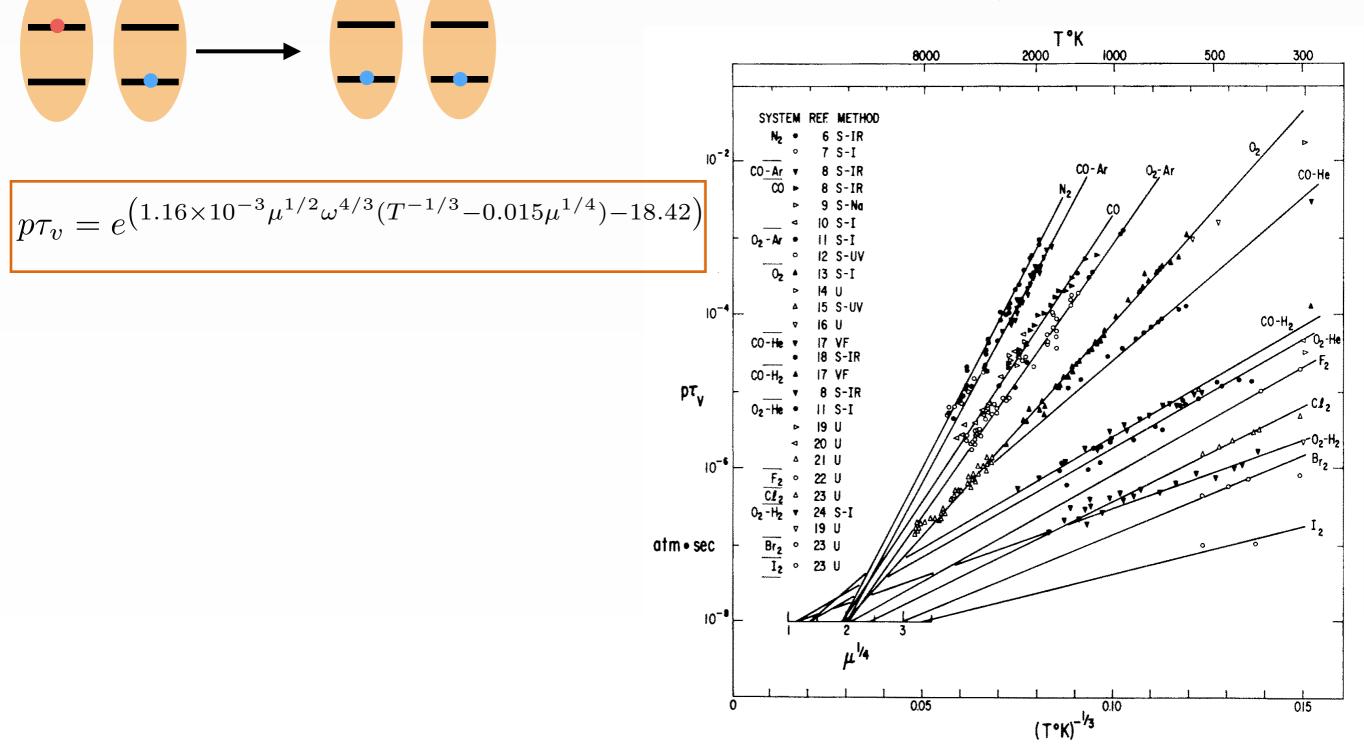




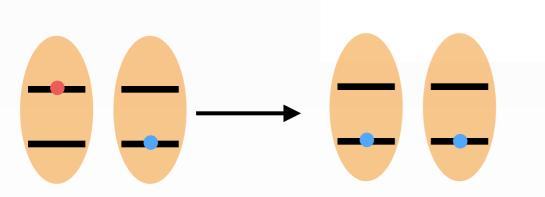
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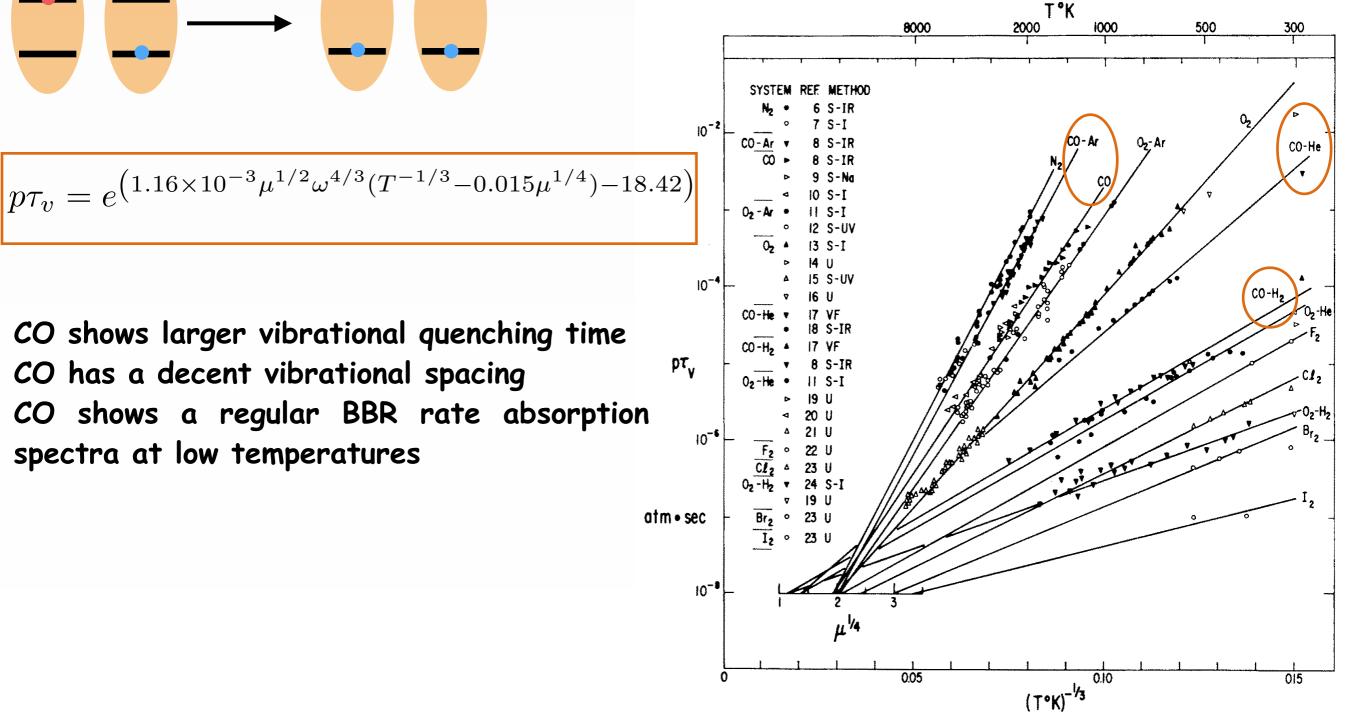


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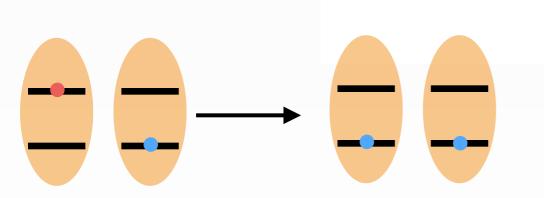
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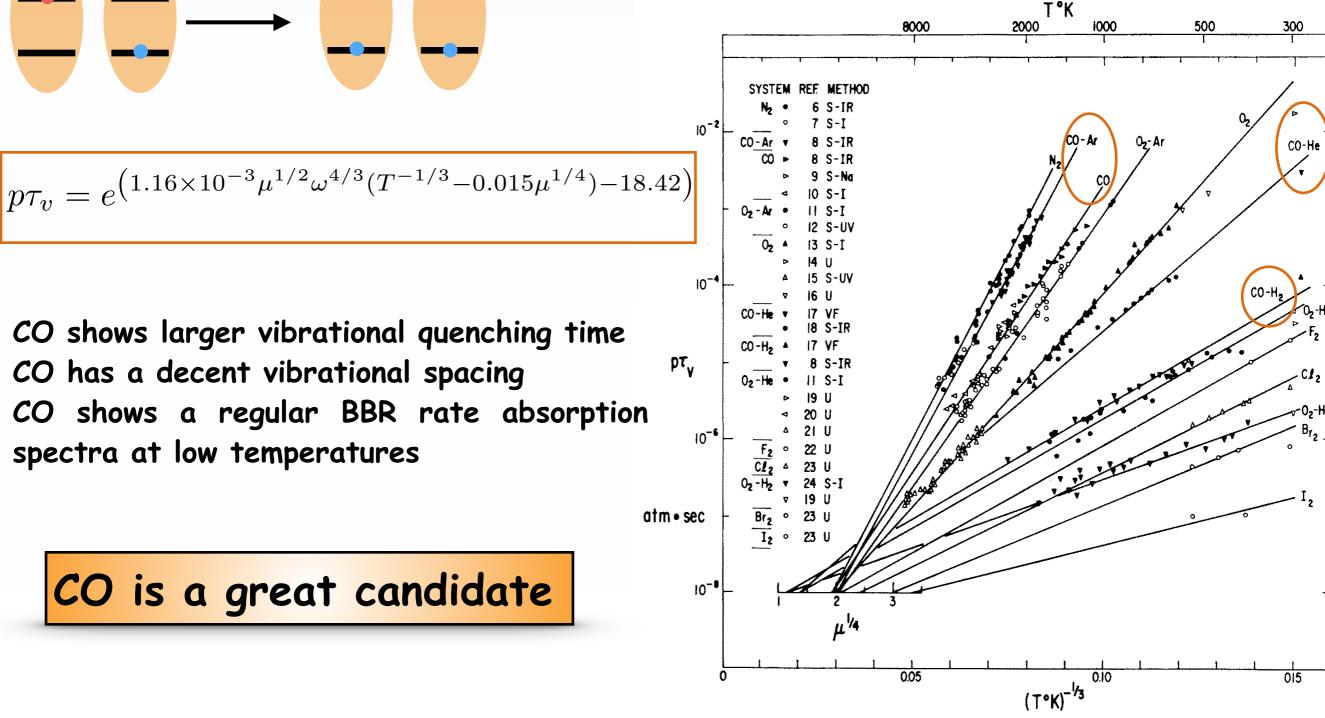
CO shows larger vibrational quenching time CO has a decent vibrational spacing CO shows a regular BBR rate absorption spectra at low temperatures

THE JOURNAL OF CHEMICAL PHYSICS VOLUME 39, NUMBER 12 15 DECEMBER 1963

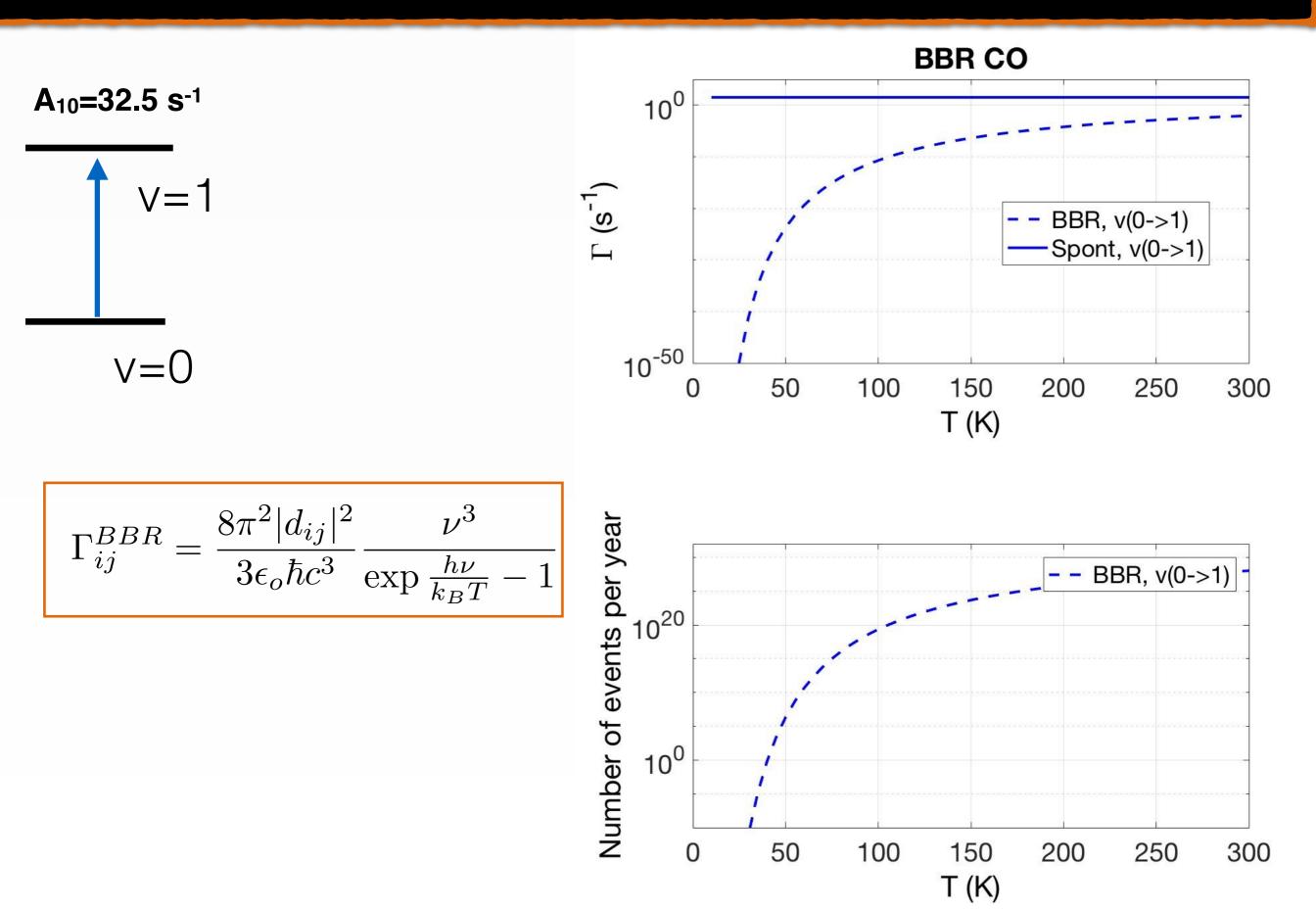


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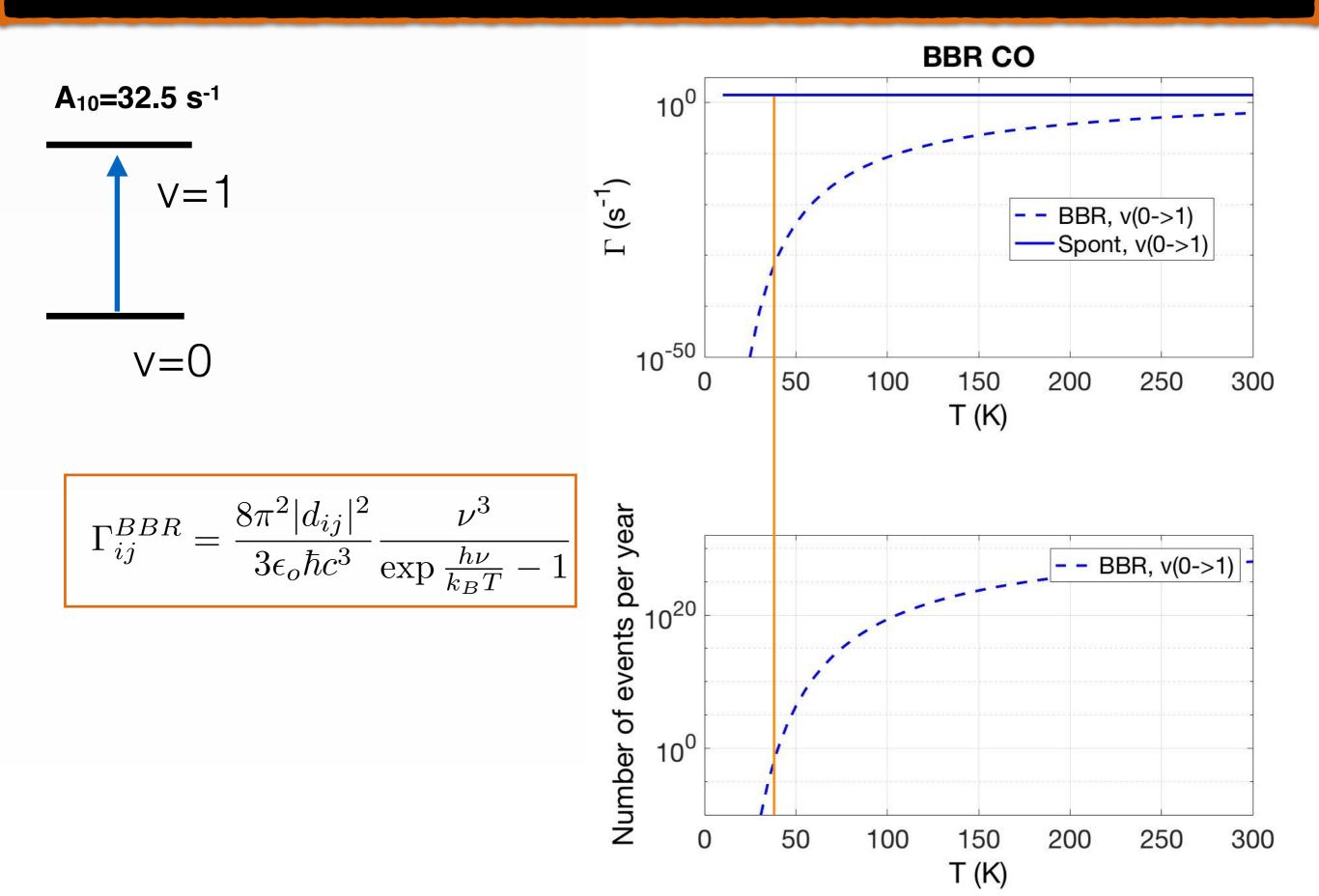
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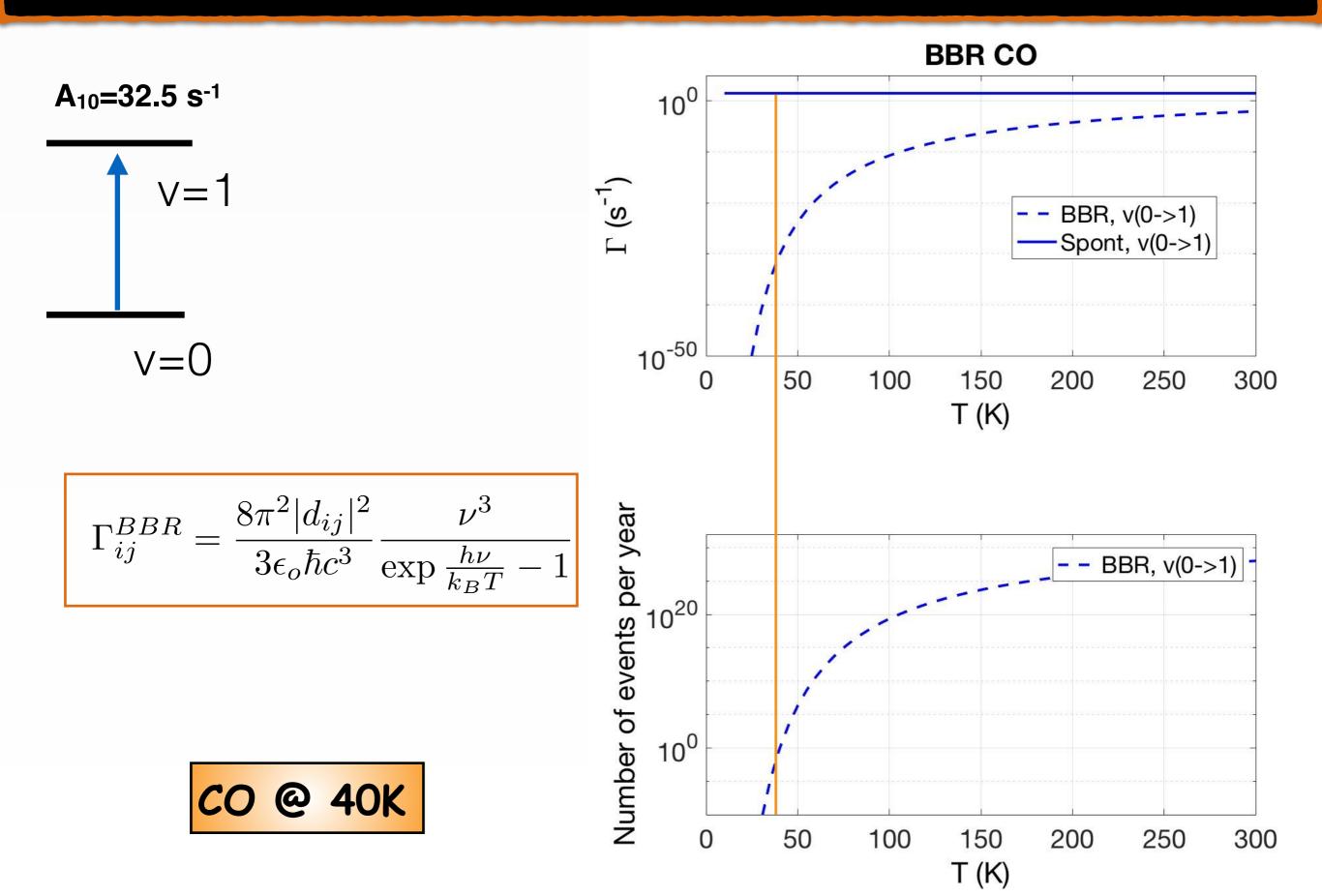
CO in detail (BBR)



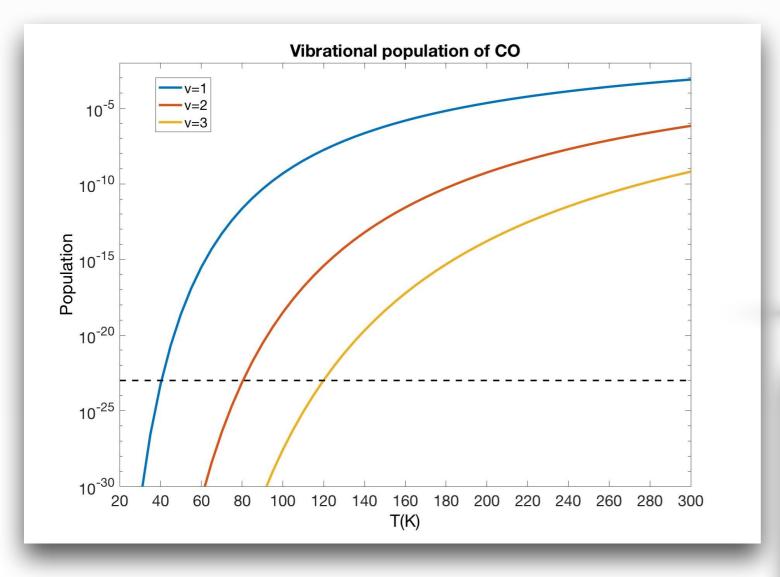
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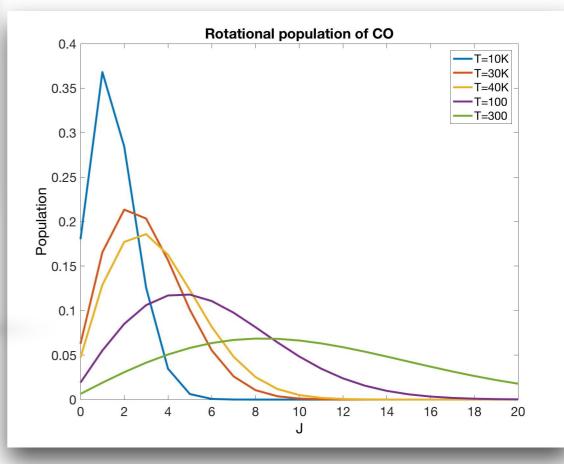
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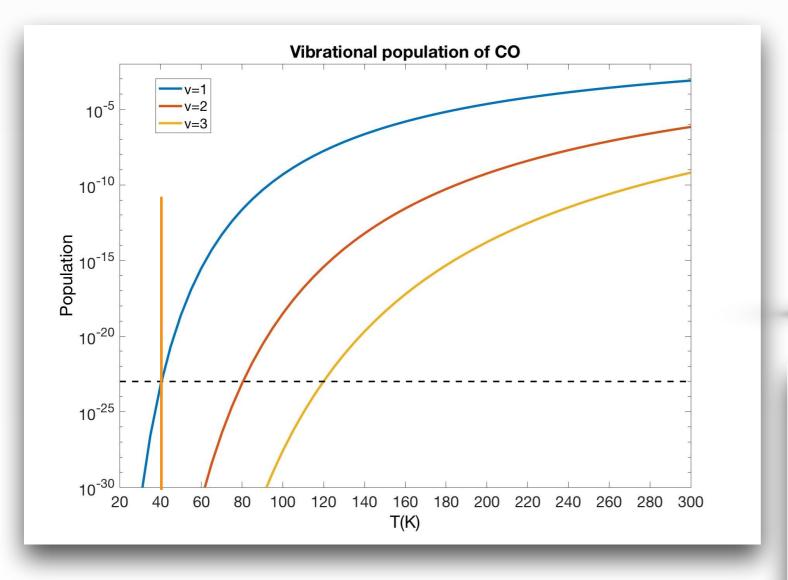
CO in detail (Rovibrational population)



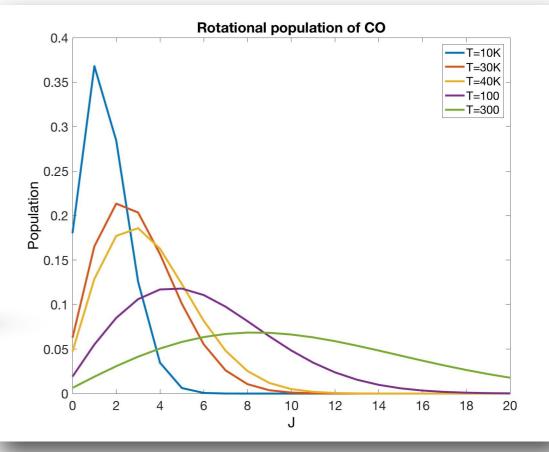
The rotational level distribution is accounted for in the rate calculation For CO @ 40 K the pressure must be $\leq 0.1 \text{ mbar}$ to have all the molecules in gas phase



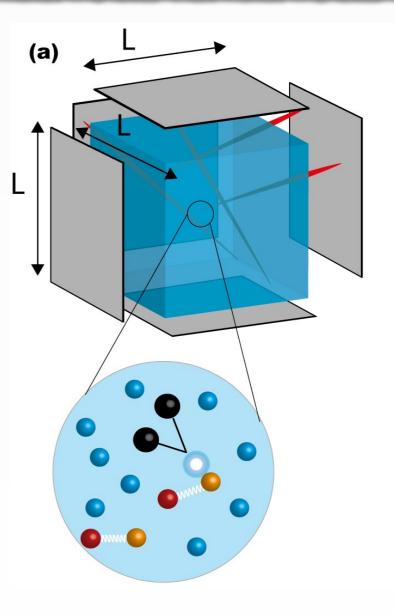
CO in detail (Rovibrational population)



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CO in detail (bonus as I promised)



$$\sigma_{\rm abs}(\omega) = \hbar \omega B_{ij} g(\omega)$$

$$g(\omega) = \frac{1}{2\pi} \frac{\gamma_{\rm el}}{\gamma_{\rm el}^2 + (\omega - \omega_0)^2}$$

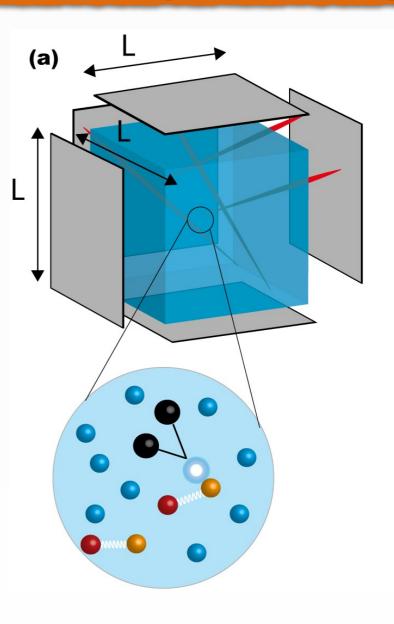
$$\gamma_{\rm el} = \rho_{\rm He} \langle \sigma_{\rm el}(v) v \rangle$$

CO in detail (bonus as I promised)

The emitted photon can be absorbed by another molecule before it reaches the detector

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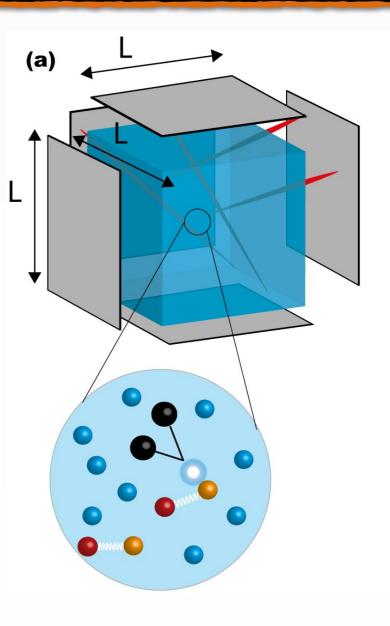
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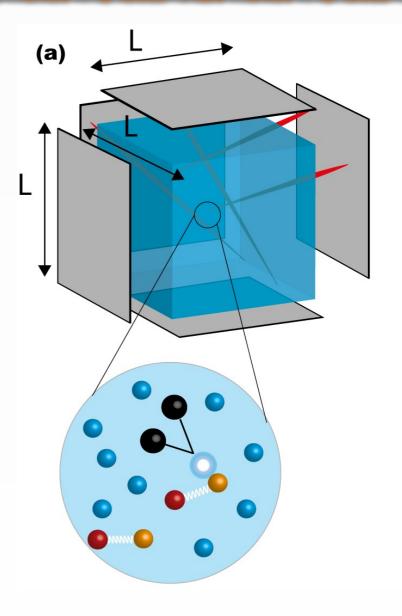
Absorption cross section

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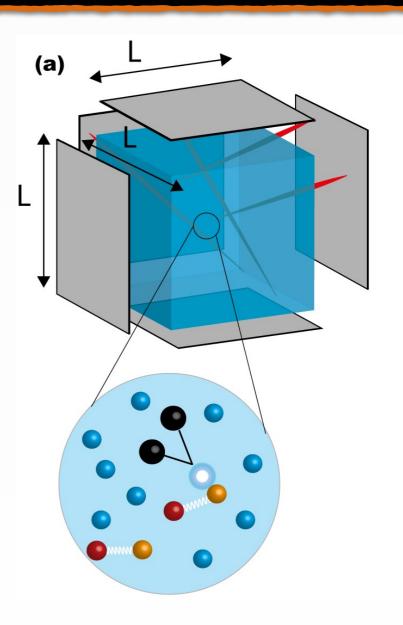
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$$l(\omega) = \frac{1}{\rho_{\rm CO}\sigma_{\rm abs}(\omega)}$$

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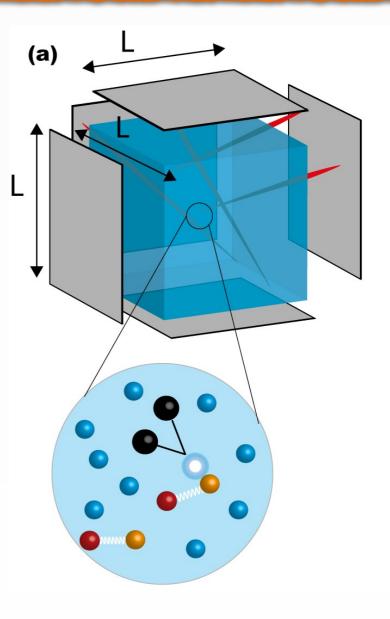
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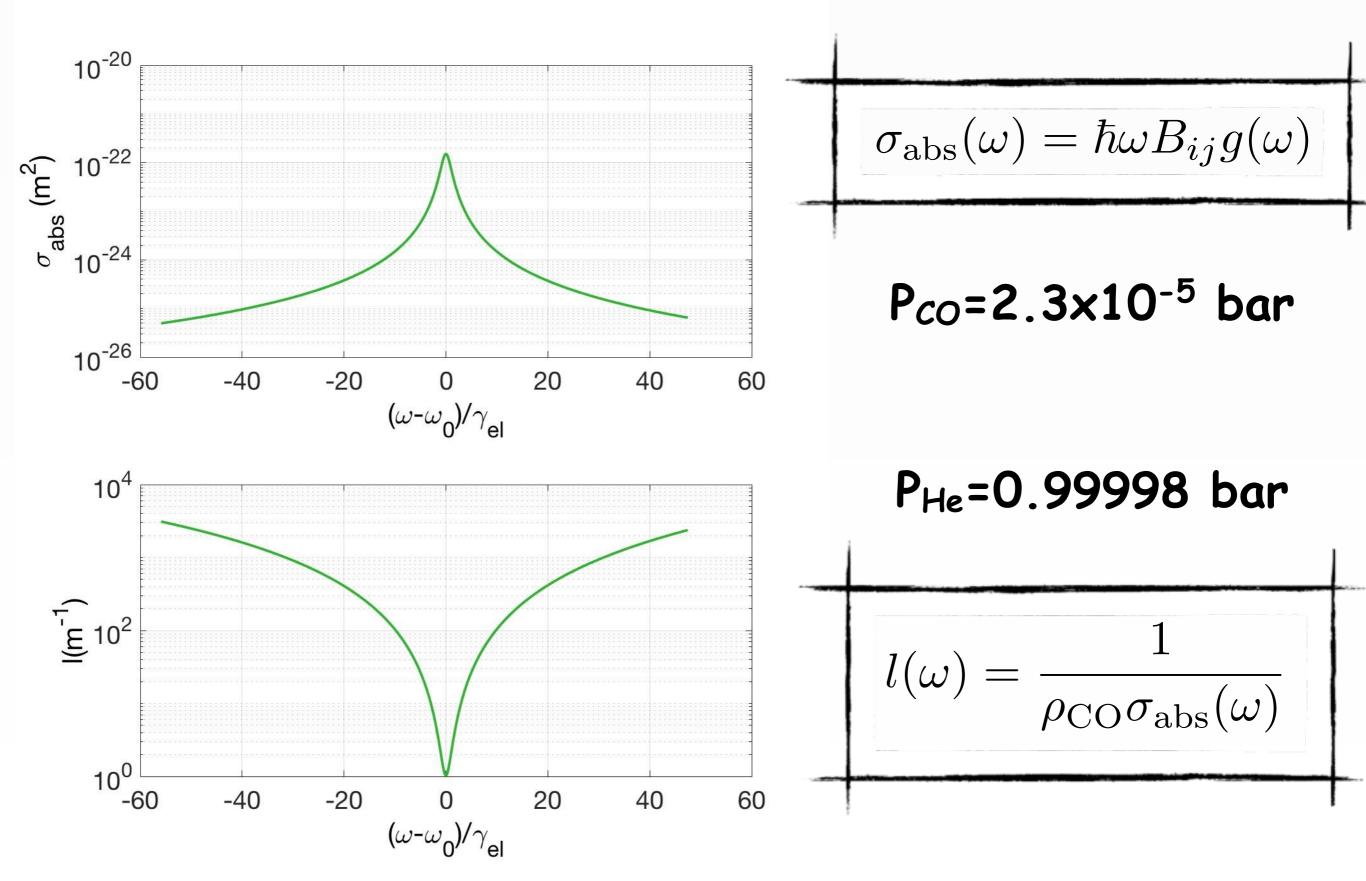
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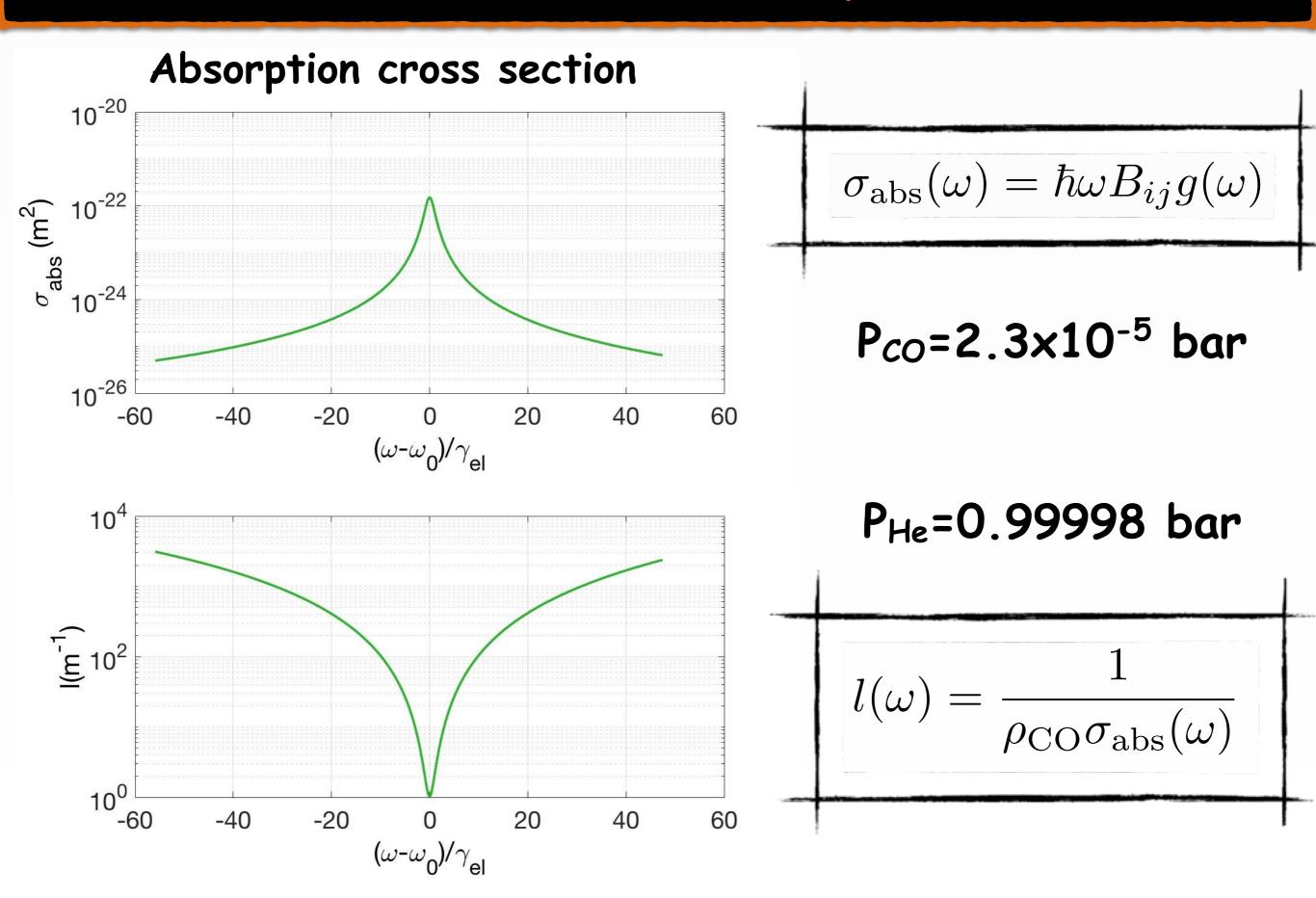
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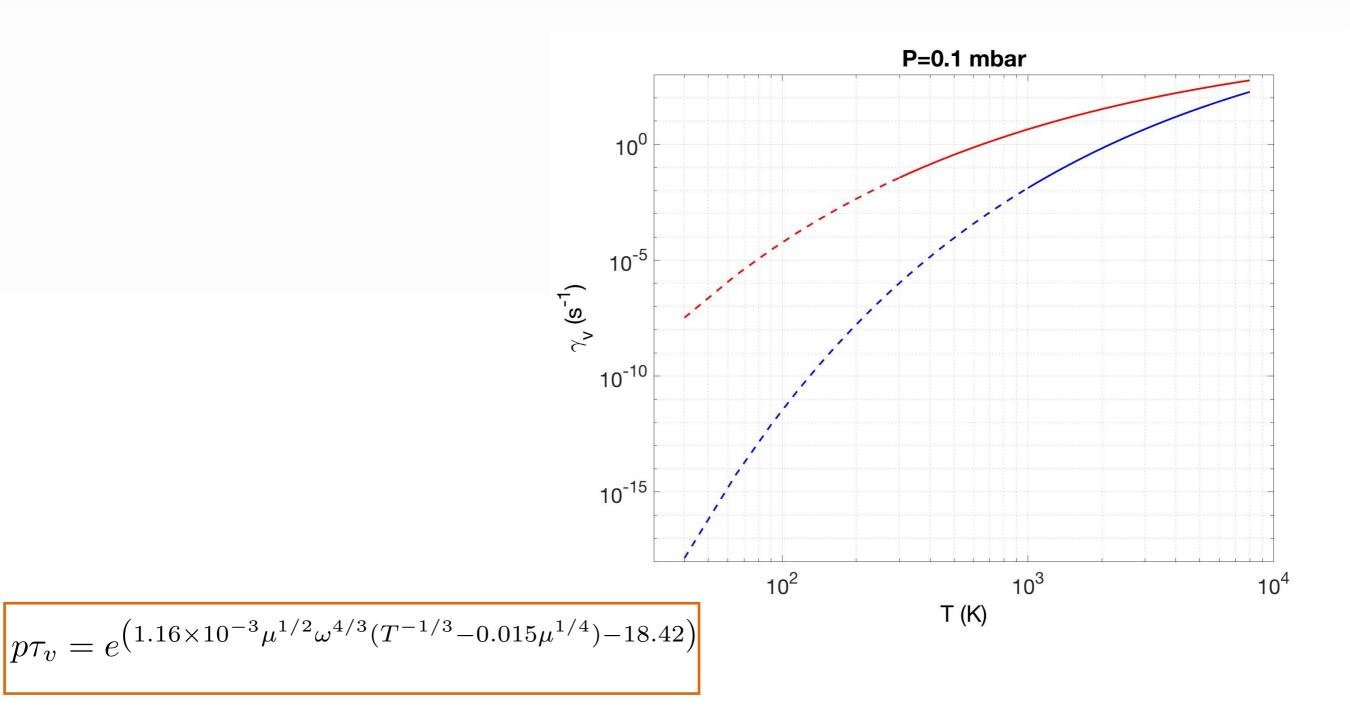
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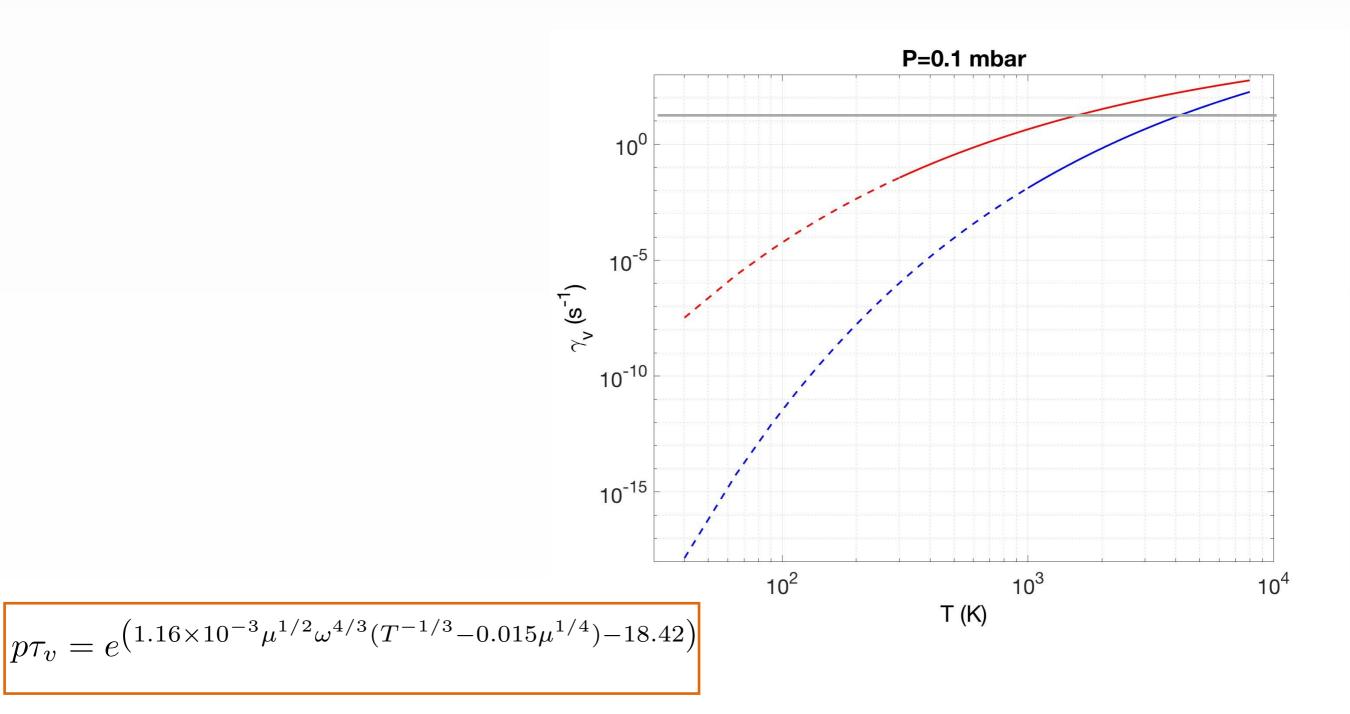
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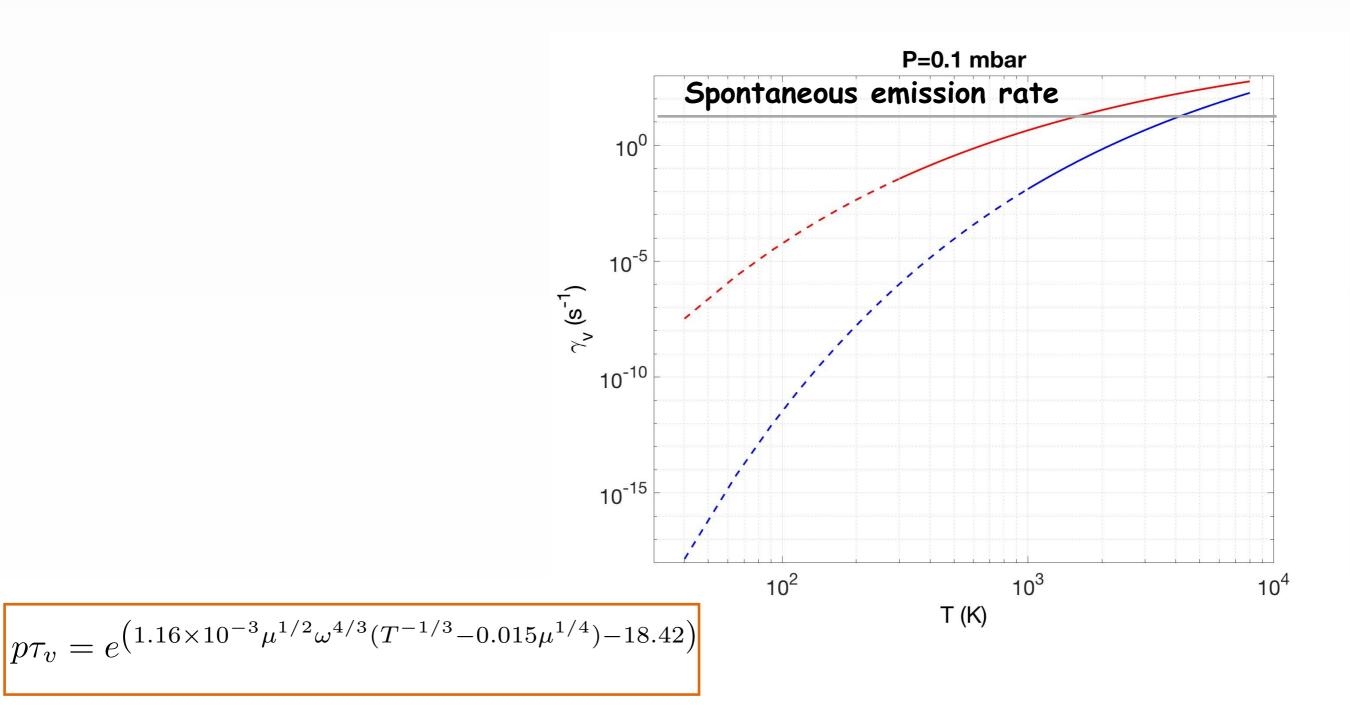












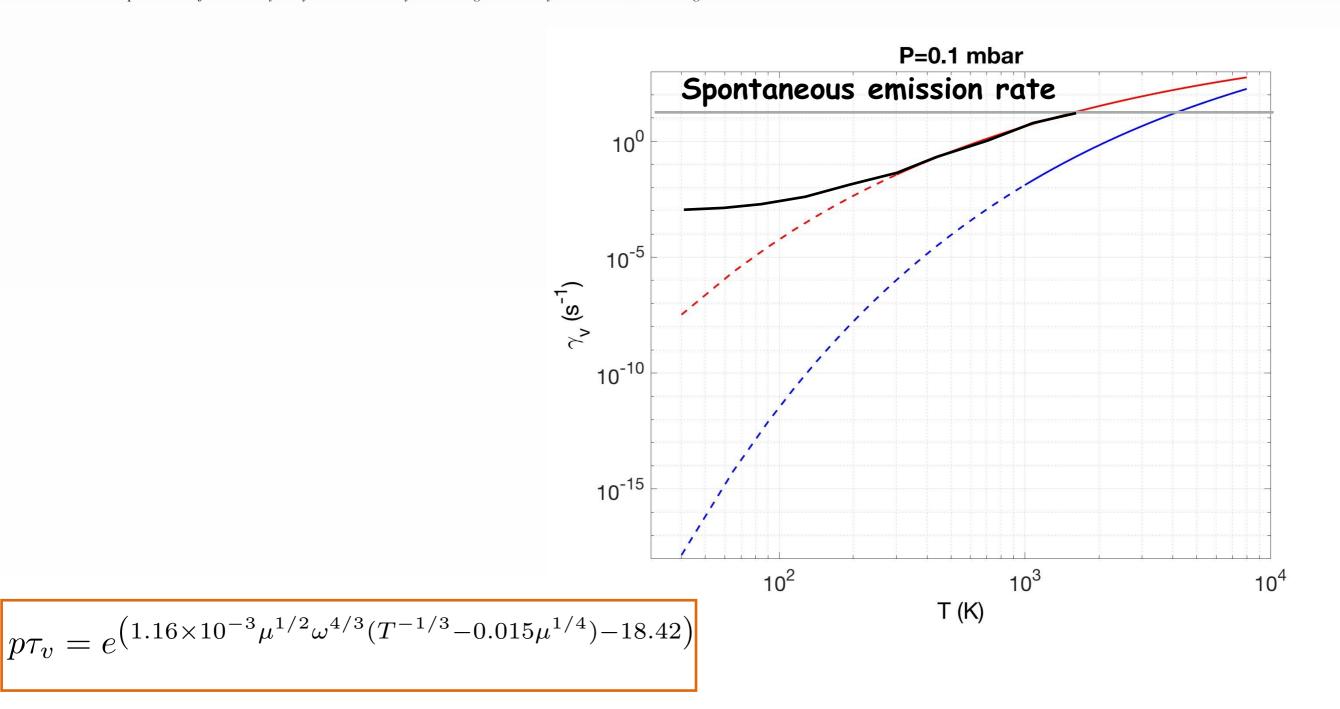
JOURNAL OF CHEMICAL PHYSICS

VOLUME 116, NUMBER 11

15 MARCH 2002

Vibrational relaxation of vibrationally and rotationally excited CO molecules by He atoms

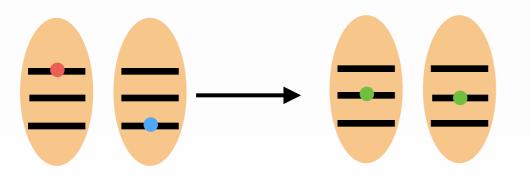
Roman V. Krems^{a)} Department of Chemistry, Physical Chemistry, Göteborg University, SE-412 96, Göteborg, Sweden



Vibrational energy exchange in CO–CO collisions at low temperature

M. C. Gower, G. Srinivasan, and K. W. Billman

Citation: The Journal of Chemical Physics **63**, 4206 (1975); View online: https://doi.org/10.1063/1.431190 View Table of Contents: http://aip.scitation.org/toc/jcp/63/10 Published by the American Institute of Physics

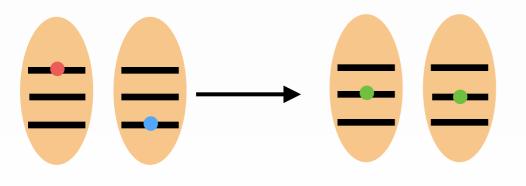


$k_{v,v-1}=1.8 \times 10^{-12} \text{ cm}^3/\text{s}$

Vibrational energy exchange in CO–CO collisions at low temperature

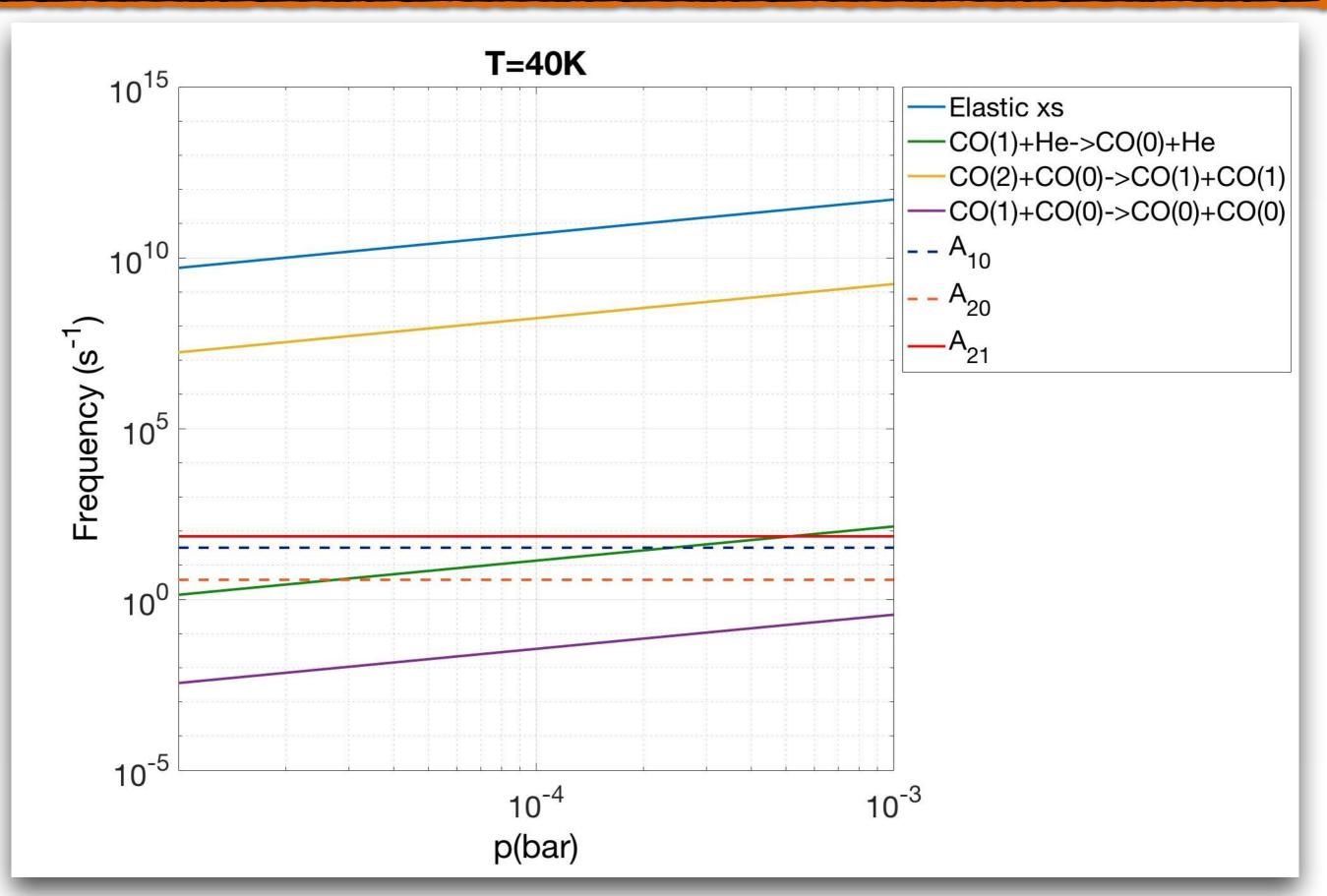
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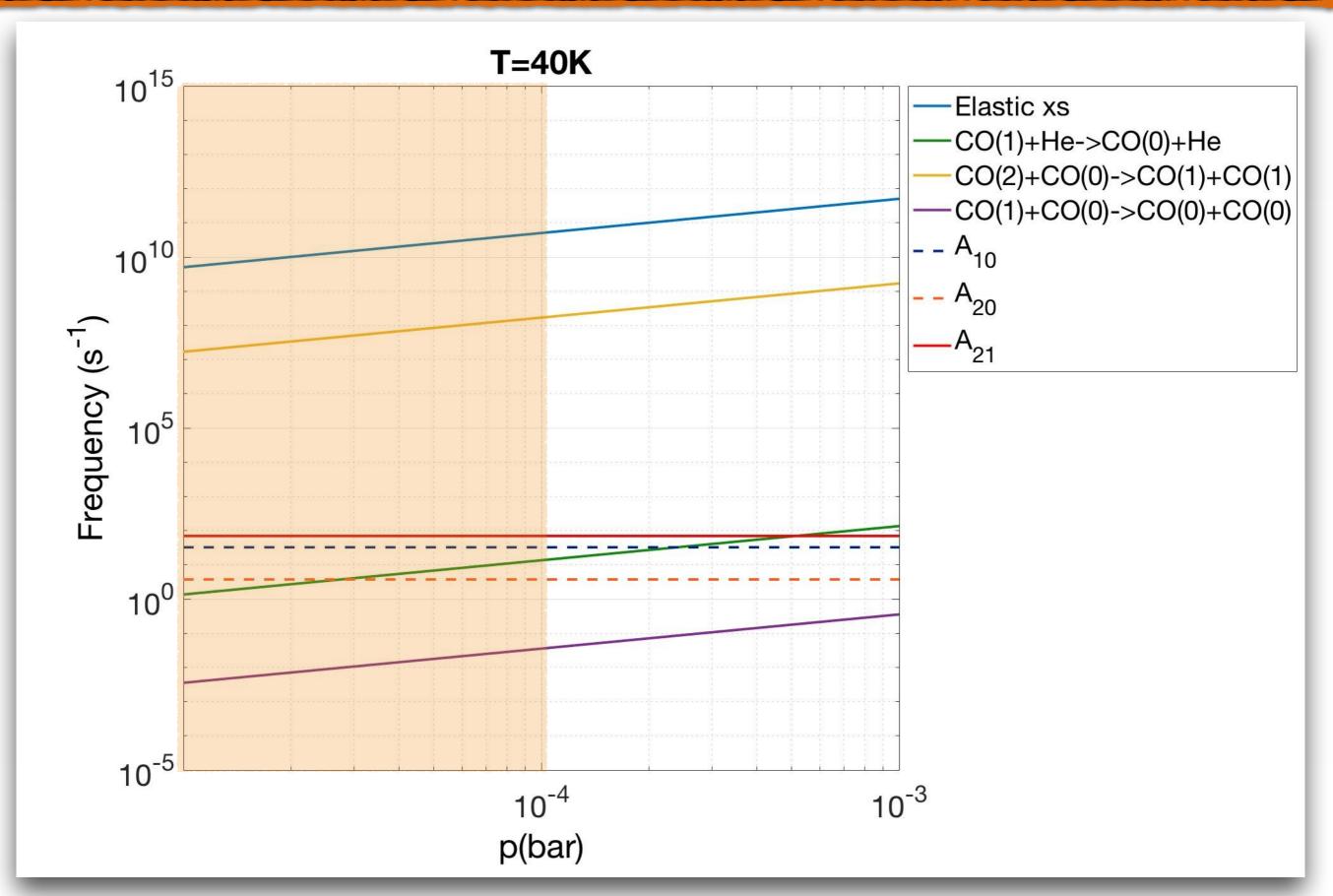
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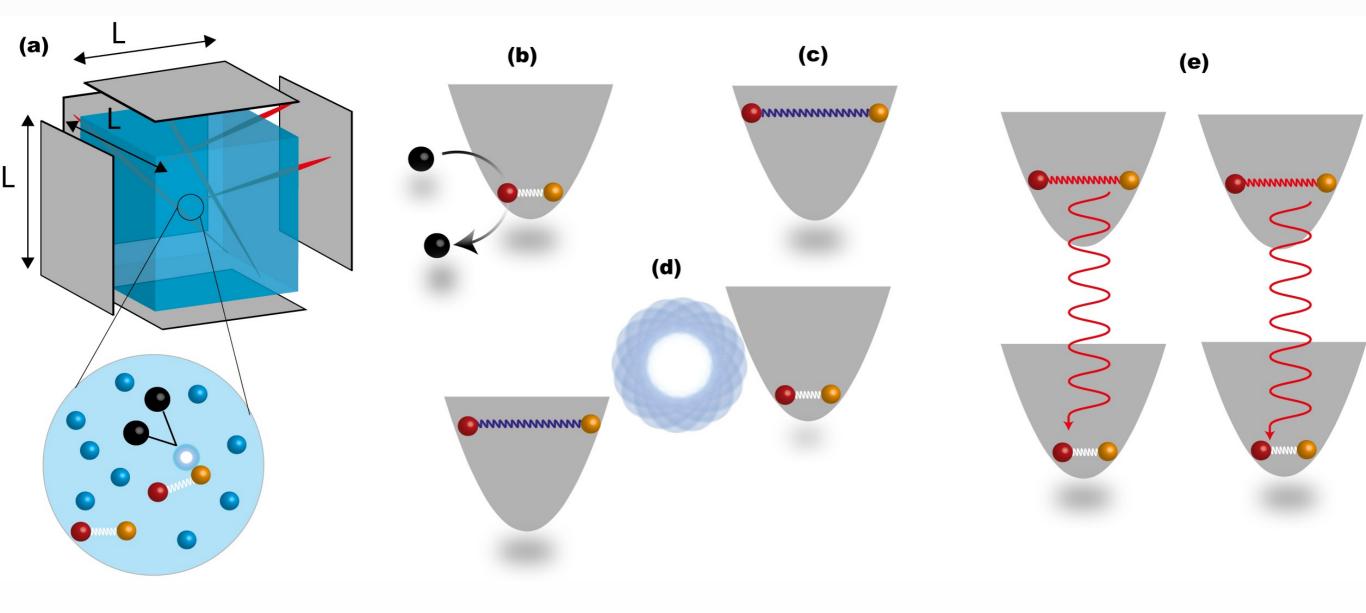


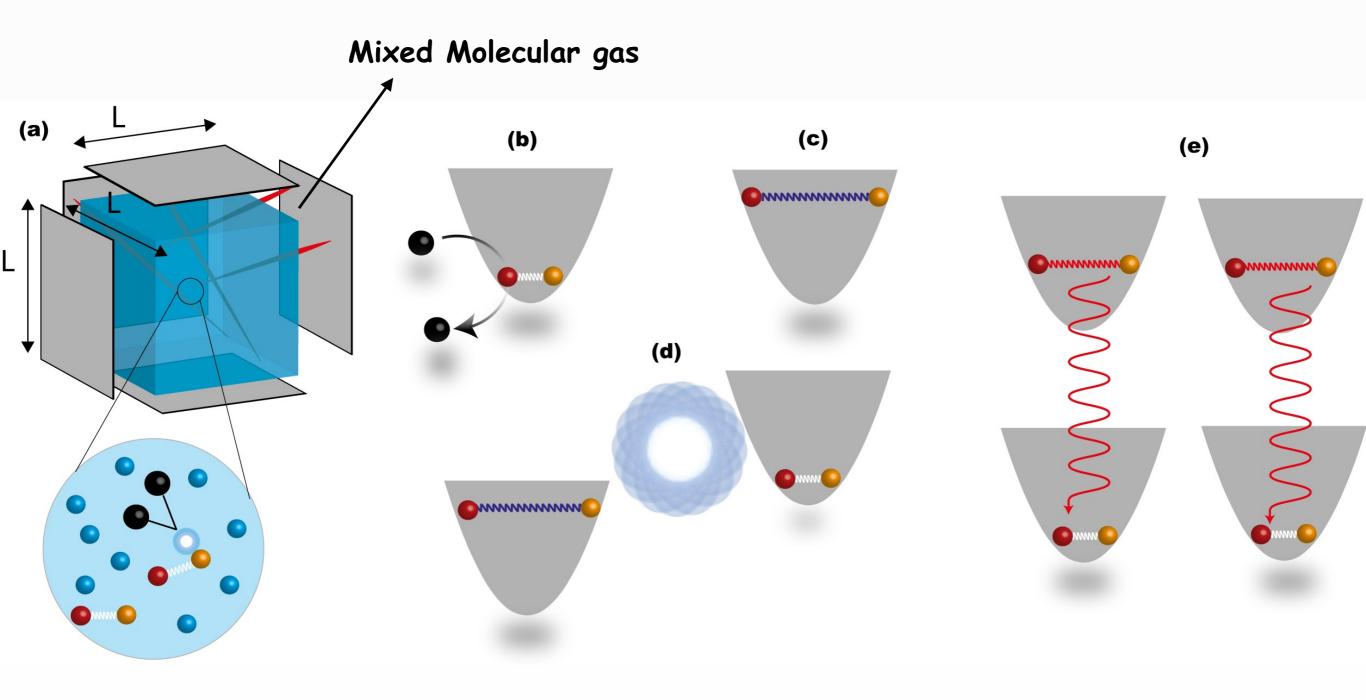
CO @ 40 K

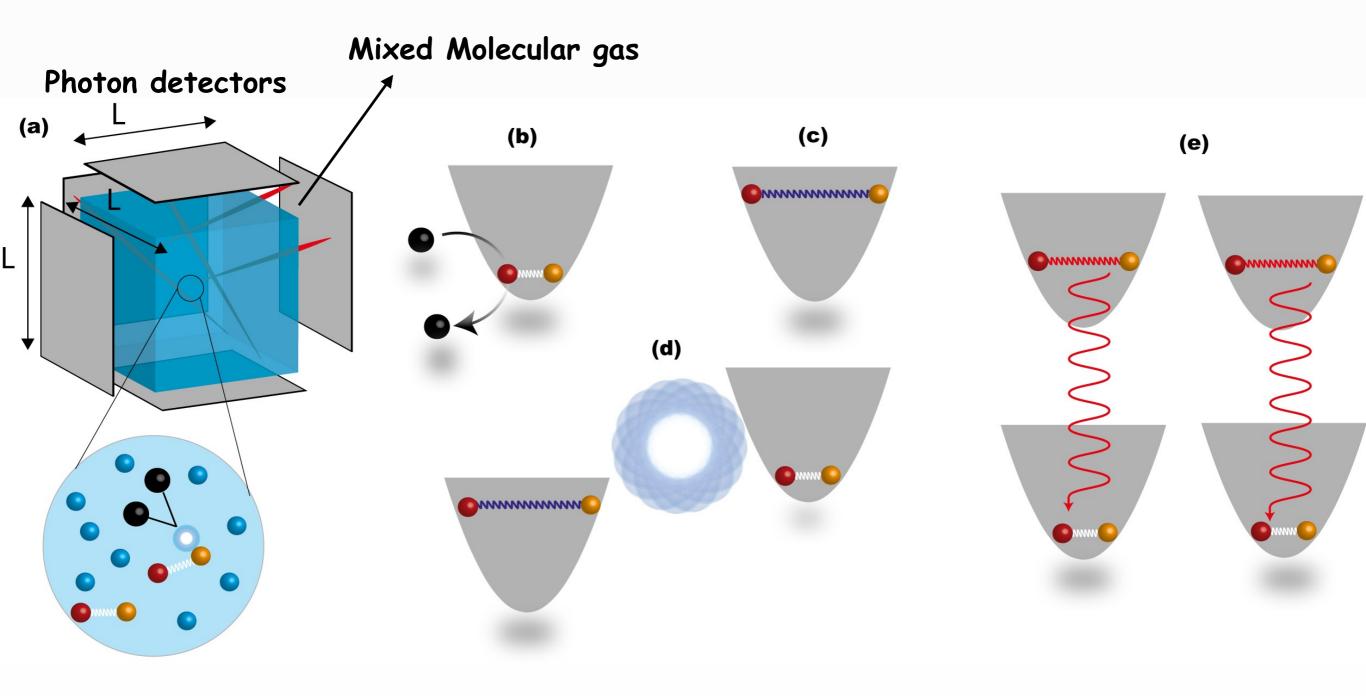
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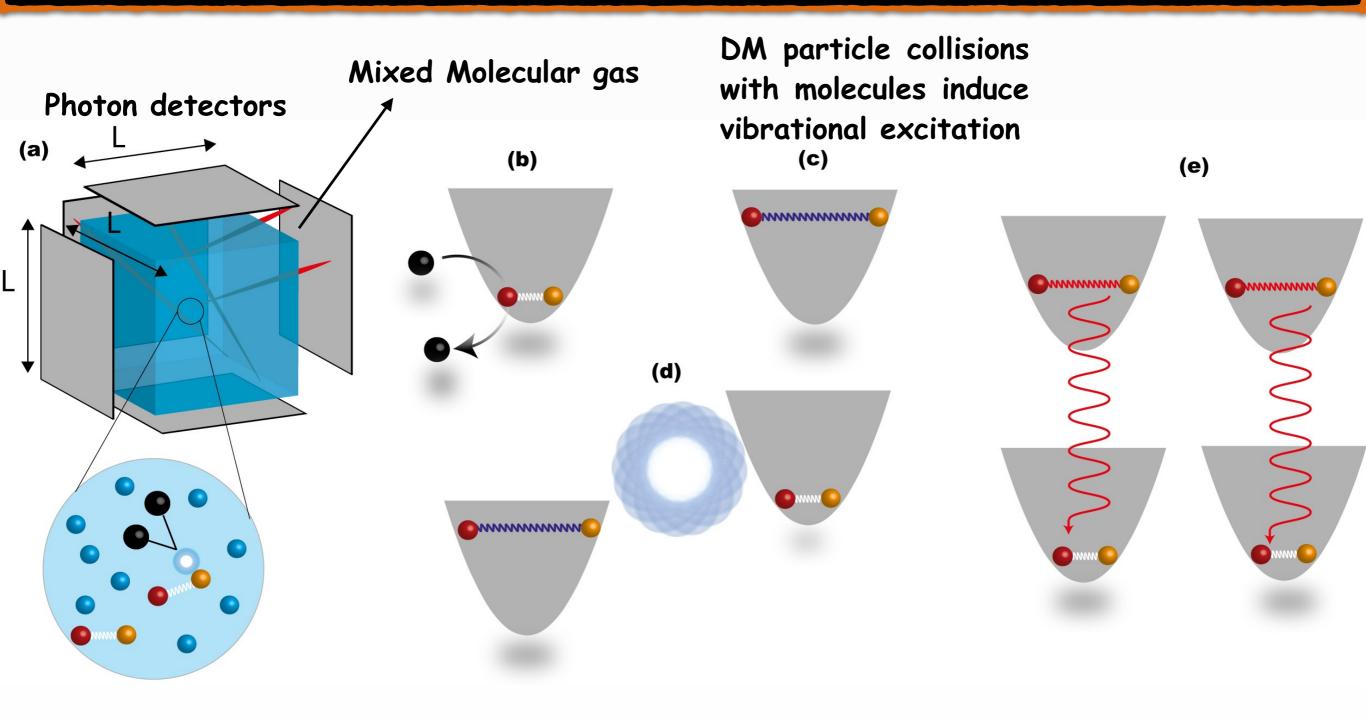


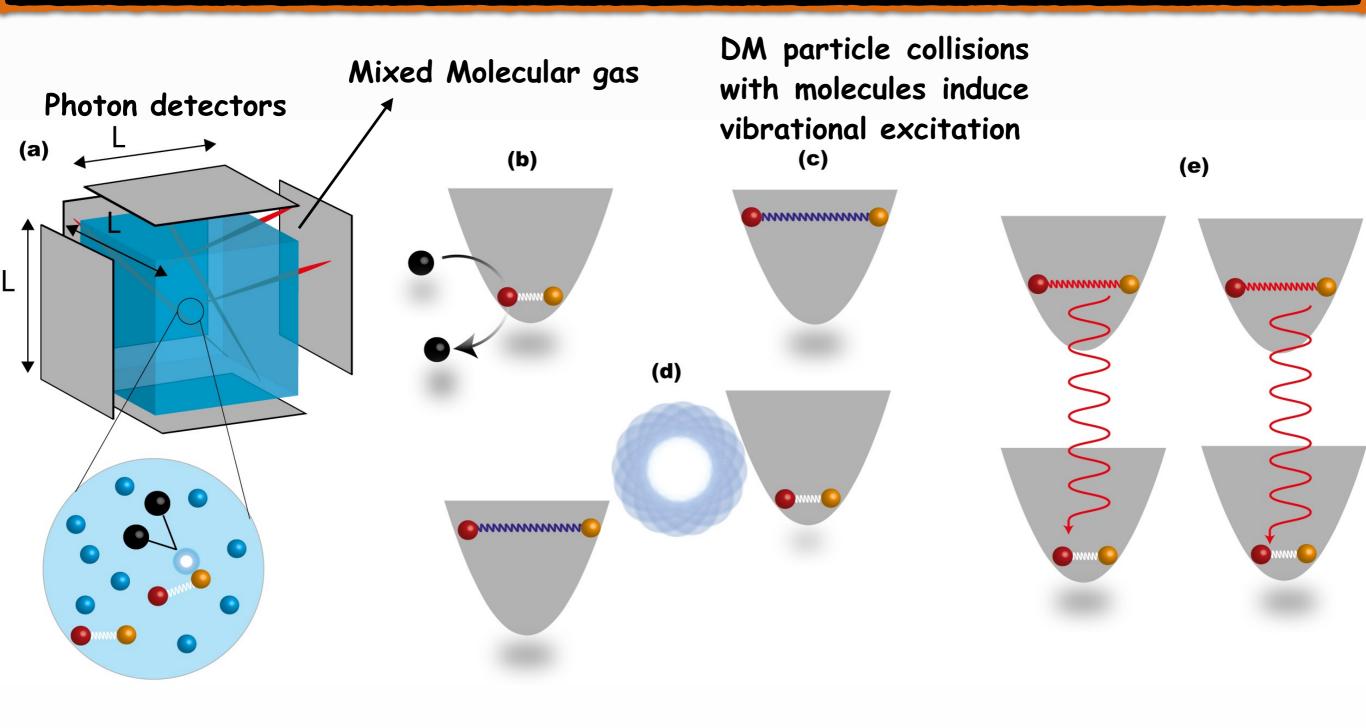




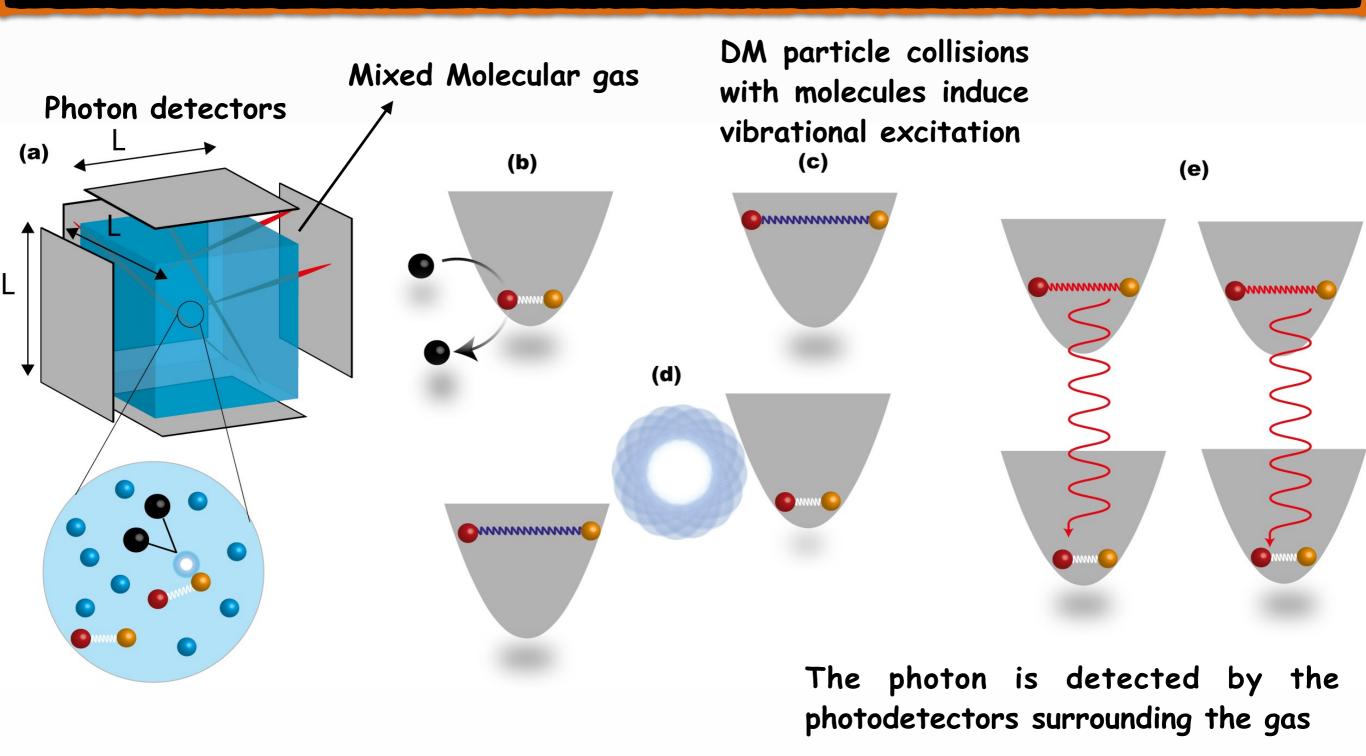






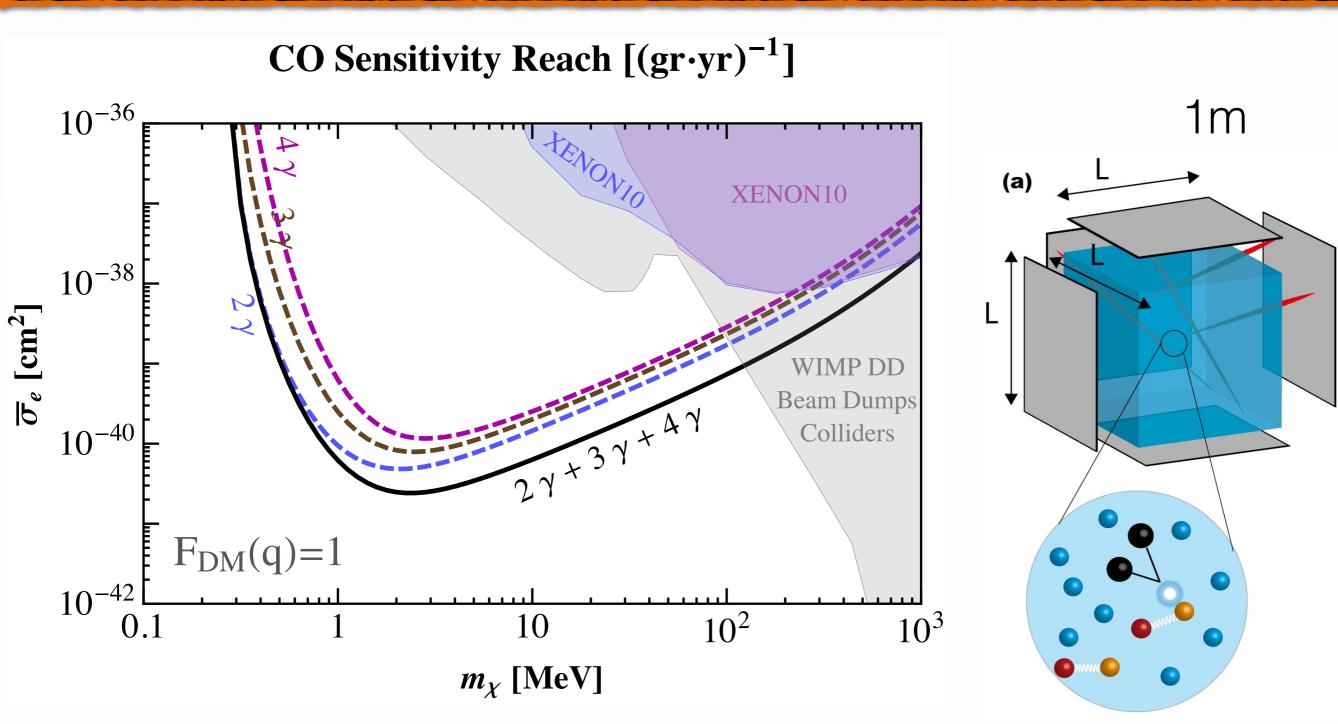


Excited molecules decay emitting a photon



Excited molecules decay emitting a photon

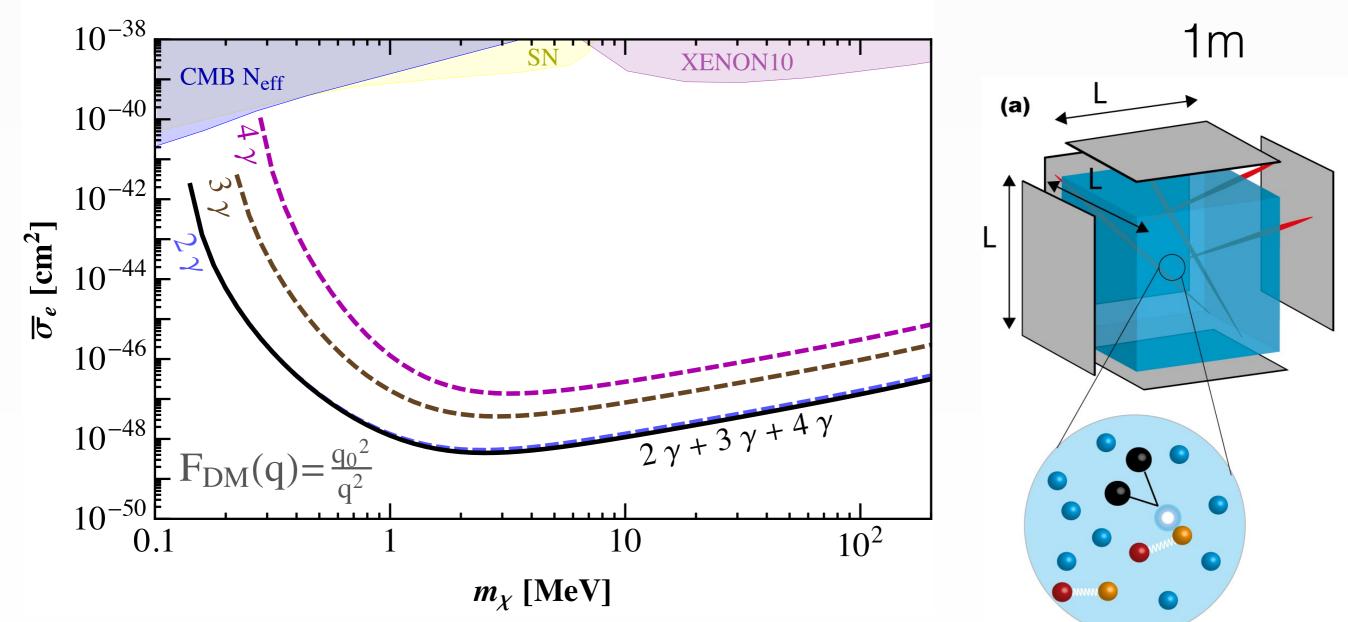
CO in detail



We will be able to explore DM particles in a mass range of almost 3 orders of magnitude, from 100 keV to 100 MeV.

CO in detail



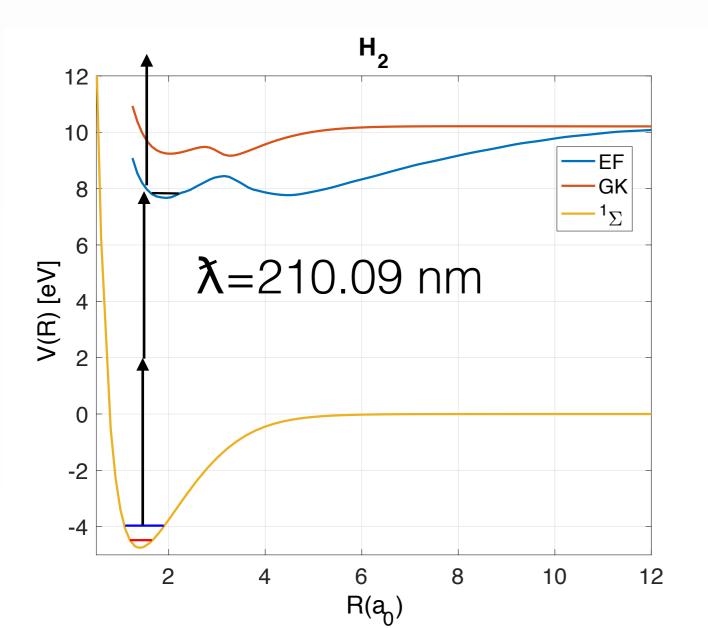


We will be able to explore DM particles in a mass range of almost 3 orders of magnitude, from 100 keV to 100 MeV.

Some future work

Electronic excitations?

REMPI through a dark state



That's all folks



Thank you so much for your attention!!!!!!!!