Data Analysis Techniques for Testing General Relativity with GWs



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Motivation

- GW astronomy:
 - We want to get the "fingerprints" of GW sources from the shapes of their waveforms
 - Want to learn about **binary parameters** and **astrophysics and fundamental physics** of these systems
 - LIGO / Virgo have set the stage with about a dozen observations in O1 & O2, more coming as we speak in O3.
- Testing General Relativity with GWs:
 - GWs emitted from merging compact binaries allow us to test GR in the strong field regime
 - So far, no evidence for deviations from GR
- First, let's look at **basic GW data analysis**:
 - How can we **infer binary parameters** given a GW signal?
 - Ingredients: matched filtering, model of (GR) waveform and noise



Introduction to GW data analysis for compact binary mergers

See also: "Guide to LIGO data analysis", arXiv:1908.11170



Current GW detector network



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Matched filtering

GW data analysis

- Detection:
 - Build a **bank of template waveforms** that covers parameter space
 - Want to keep fraction of missed signals small
 - Cross-correlate templates against detector data
- Parameter estimation:
 - Given a likely detection follow it up with Bayesian methods to find the probability distribution of the binary parameters given the signal in the detectors and the PSD of the detector noise
 - Need high accuracy waveforms to make sure that we don't infer wrong / biased parameters
- Both need fast & accurate waveform models
 - Require O(10⁷) O(10⁸) waveform evaluations!

Interferometric GW detectors

Angular response of IFO $|F_A|$ (a) Plus (+)
(b) Cross (x)

$$h = h_A F^A \qquad F^A(\hat{n}, \psi) = D^{ab} e^A_{ab}$$
$$A = +, \times, \dots$$
$$D^{ab} = \frac{1}{2} \left(d^a_x d^b_x - d^a_y d^b_y \right)$$

What do we need to model?

• A waveform model is a parametrized function of the waveform polarizations $h_{+,\times}(t; \vec{\lambda})$ or complex modes $h_{lm}(t; \vec{\lambda})$:

$$(h_{+} - ih_{\times})(t; \overrightarrow{\lambda}) = \sum_{l,m} h_{lm}(t; \overrightarrow{\lambda})^{-2} Y_{lm}(\theta, \phi)$$

- Need to model the inspiral, merger and ringdown stages in binary black hole coalescence.
- GW detectors record GW strain:

$$h(t; \overrightarrow{\theta}) = h_+(t; \overrightarrow{\lambda})F_+(\hat{n}, \psi) + h_\times(t; \overrightarrow{\lambda})F_\times(\hat{n}, \psi)$$

[Baumgarte & Shapiro, Numerical Relativity]

numerical perturbation

methods

relativity

post-Newtonian

techniques

Model parameters

Intrinsic parameters:

masses, spins, eccentricity, tidal deformability

Extrinsic parameters:

time, sky position, distance, orientation, reference phase

Credit: LIGO/Virgo

To find the signal: Likelihood

- Subtract signal model from data $d h(\vec{\lambda})$, where $d = n + h_{true}$
- Assumptions: noise is Gaussian (zero mean) and stationary

$$\mathscr{L}(d \mid \overrightarrow{\lambda}) \propto \exp\left[-\frac{1}{2}\left(d - h(\overrightarrow{\lambda}), d - h(\overrightarrow{\lambda})\right)\right]$$

Parameter Estimation

• **Posterior probability** of model parameters $\vec{\lambda}$ given the data \vec{d} (Bayes' Theorem):

$$p(\overrightarrow{\lambda} | \overrightarrow{d}) \propto p(\overrightarrow{\lambda}) \mathscr{L}(\overrightarrow{d} | \overrightarrow{\lambda})$$

- Need Models for signal and noise & specify prior knowledge
- Numerically sample the posterior distribution

GWTC-1 Catalog

Tests of GR with GWs

See LVC, arXiv:1903.04467 Results on GWTC-1 Catalog

Residuals test

• **Residual:** detector data - maximum likelihood template

$$\mathscr{R}(t) := d(t) - h(\overrightarrow{\lambda}_{\max \mathscr{L}}; t)$$

- Compute 90% upper limit of SNR of residuals $\rho_{90}(\mathcal{R})$
 - describe signal as Gaussian noise + coherent signal
- Repeat analysis for noise-only detector data many times: $p(\rho_{90}(\mathcal{R}_{noise}))$
- At which percentile does $\rho_{90}(\mathcal{R})$ lie in $p(\rho_{90}(\mathcal{R}_{noise}))$?

• p-values > 0.05

• Meta p-value = 0.4

No statistically significant evidence for deviations from GR

Inspiral-merger-ringdown consistency test

- Check whether final mass M_f and final spin a_f inferred from the low and high frequency parts of the signal are consistent.
- Compute posterior distributions with different frequency bounds for likelihood integral:
 - Transition point: innermost stable circular orbit (ISCO)
 - Normalization:

$$\Delta M_f / \bar{M}_f := \frac{M_f^{\mathrm{I}} - M_f^{\mathrm{MR}}}{(M_f^{\mathrm{I}} + M_f^{\mathrm{MR}})/2}$$

- lines = 90% credible regions of $p(\Delta M_f/\bar{M}_f, \Delta a_f/\bar{a}_f | d)$
- Plus sign = GR value
- GR value recovered at < 80% credible level

All events consistent with GR

LSC

• Compute **posteriors** $p(\delta \hat{p}_n | d)$ of

Parametrized test of GW generation

- Compute posteriors $p(op_n | a)$ of deviations in coefficients p_n of the phase of waveform models
- Templates have to be phase coherent with the signal to a fractions of a radian to capture the SNR of an event
- post-Newtonian phase up to 3.5PN or $\mathcal{O}((v/c)^7)$

$$\Phi(\overrightarrow{\theta}; f) = 2\pi f t_c - \varphi_c - \pi/4$$
$$+ \frac{3}{128\eta} (\pi M f)^{-5/3} \sum_{i=0}^7 \varphi_i(\overrightarrow{\theta}) (\pi M f)^{i/3}$$

• Fractional deviations, $\delta \hat{\varphi}_n := (\varphi_n - \varphi_n^{\text{GR}})/\varphi_n^{\text{GR}}$ except for -1PN and 0.5PN which are *zero* in GR

• Combined posteriors for $\delta \hat{\varphi}_n$ for the most significant events

LSC

- Horizontal lines: 90% CR
- GR value = 0

All events consistent with GR

Parametrized test of GW propagation

 Assume a phenomenological dispersion relation

$$E^2 = p^2 c^2 + A_\alpha p^\alpha c^\alpha$$

where **GR** has $A_{\alpha} = 0$

- Massive graviton ($\alpha = 0, A_{\alpha} > 0$)
- Lorentz-violating theories
- Dispersive propagation of GWs except for $\alpha = 2$
- Use GR waveforms and only modify propagation

- 90% credible upper bounds on $|A_{\alpha}|$
- 1peV $\simeq h \times 250$ Hz

•
$$m_g \le 4.7 \times 10^{-23} eV/c^2$$

No indication for dispersive propagation of GWs.

Constraints on Alternative Polarizations

- For events observed with three or more GW detectors:
 - is possible to distinguish purely tensor from purely vector or purely scalar perturbations
 - detectors have different antenna patterns for different polarizations
- Compute **Bayes factors** between options:
 - GW170817: **BF > 10²⁰** in favor of **pure tensor polarizations**
 - Weaker constraints from binary black hole events (SNR and sky localizations play a role)

Conclusion & Outlook

- Will have many more events during O3 and future observing runs: statistical errors of the combined results will soon decrease significantly
- Potential sources of systematic errors:
 - imperfect modeling of GR waveforms
 - & missing physics problem for 2G design & 3G detectors (ET)
 - calibration uncertainties
 - noise artifacts
- Goal: inspiral-merger-ringdown waveforms for alternative theories of gravity
 - This is hard: need to solve initial value problem for 2-body problem: well-posedness?
 - Obtain waveform in the **inspiral** regime
 - Develop numerical relativity codes for alternative theories

Einstein Telescope and 3G Science

Science Targets for 3G GW Detectors

- determine the properties of the hottest and densest matter in the Universe
- *reveal* the merging BH population throughout the Universe; seeds
- investigate the particle physics of the primeval Universe and probe its dark sectors

- 3G Science Case Document, led subgroup on waveforms & data analysis
- 3G White paper on Extreme Gravity and Fundamental Physics, arXiv:1903.09221
- 3G network sensitivity ~ 20 times better than aLIGO/Virgo: precision measurements!

Extra slides

Phase contributions from dispersion relation

- Treat the GW as a stream of gravitons, which travel at the particle velocity $v_p/c = pc/E = 1 A_{\alpha}E^{\alpha-2}/2 + O(A_{\alpha}^2)$
- Length scale of modifications to the Newtonian potential $\lambda_A := hc |A_{\alpha}|^{1/(\alpha-2)}$
- Fourier domain dephasing due to modified dispersion relation:

$$\delta \Phi_{\alpha}(f) = \operatorname{sign}(A_{\alpha}) \begin{cases} \frac{\pi D_{\mathrm{L}}}{\alpha - 1} \lambda_{A, \operatorname{eff}}^{\alpha - 2} \left(\frac{f}{c}\right)^{\alpha - 1}, & \alpha \neq 1 \\ \frac{\pi D_{\mathrm{L}}}{\lambda_{A, \operatorname{eff}}} \ln \left(\frac{\pi G \mathcal{M}^{\operatorname{det}} f}{c^{3}}\right), & \alpha = 1 \end{cases}$$

$$\lambda_{A,\text{eff}} \coloneqq \left[\frac{(1+z)^{1-\alpha} D_{\text{L}}}{D_{\alpha}} \right]^{1/(\alpha-2)} \lambda_{A} \qquad D_{\alpha} = \frac{(1+z)^{1-\alpha}}{H_{0}} \int_{0}^{z} \frac{(1+\bar{z})^{\alpha-2}}{\sqrt{\Omega_{\text{m}}(1+\bar{z})^{3} + \Omega_{\Lambda}}} d\bar{z},$$

Mirshekari et al, PRD 85, 024041 (2012)