

THEORY STATUS AFTER ICHEP 2014

Searching for New Physics
at the Next LHC Run

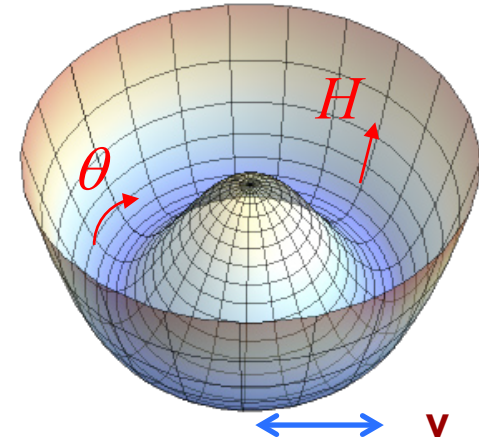
A. Pich

IFIC, Univ. València - CSIC

IFAE, Bellaterra, Spain, 24 November 2014

Great success of the Standard Model

BEGHHK (\equiv Higgs) Mechanism

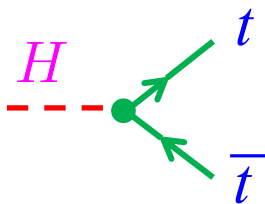


$$SU(2)_L \otimes U(1)_Y \quad v = 246 \text{ GeV}$$

$$M_Z \cos \theta_W = M_W = \frac{1}{2} v g$$



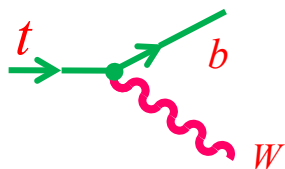
The Heaviest Mass Scale



$$y_t = \frac{\sqrt{2}}{v} m_t = 2^{3/4} G_F^{1/2} m_t = 1 \quad (0.995)$$

The top quark:

- Sensitive probe of Electroweak Symmetry Breaking
- Non-perturbative (strong) dynamics ?
- Very different from other quarks: $y_b = 0.025$, $y_c = 0.007 \dots$
- Is it really a SM quark?



So far, we only know
the decay $t \rightarrow b W^+$

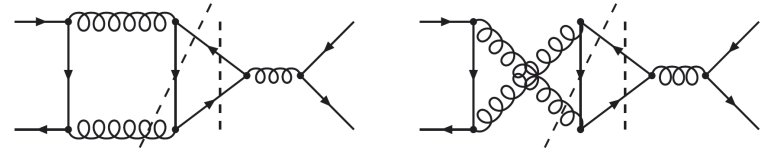
| Single-top | $ V_{tb} $ |
|------------|-----------------|
| ATLAS '14 | > 0.88 (95% CL) |
| CMS'14 | > 0.92 (95% CL) |
| CDF'14 | > 0.84 (95% CL) |
| D0 '13 | > 0.92 (95% CL) |

$t\bar{t}$ Production Asymmetries

Tevatron: $A_{FB} \equiv A_{t\bar{t}} = \frac{N(\Delta y > 0) - N(\Delta y < 0)}{N(\Delta y > 0) + N(\Delta y < 0)}$

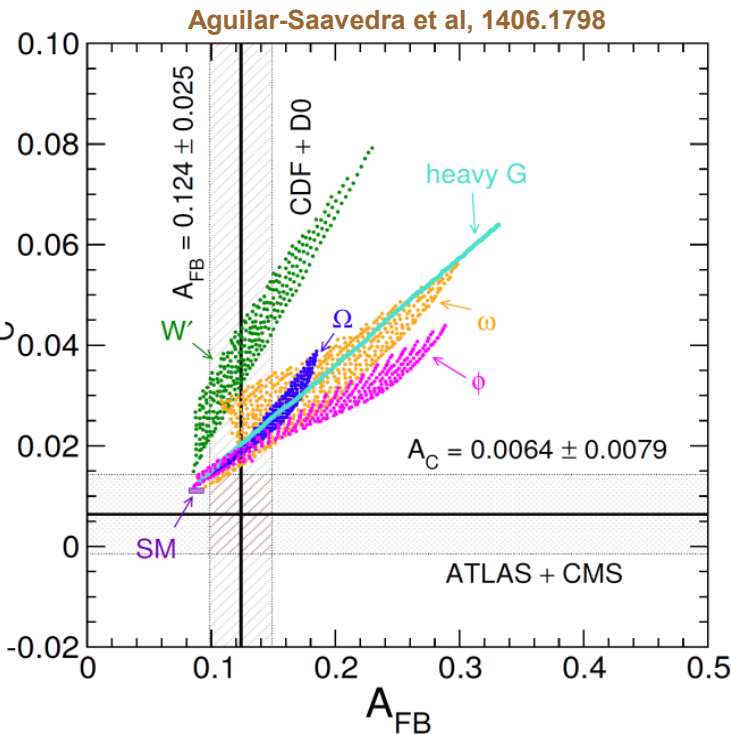
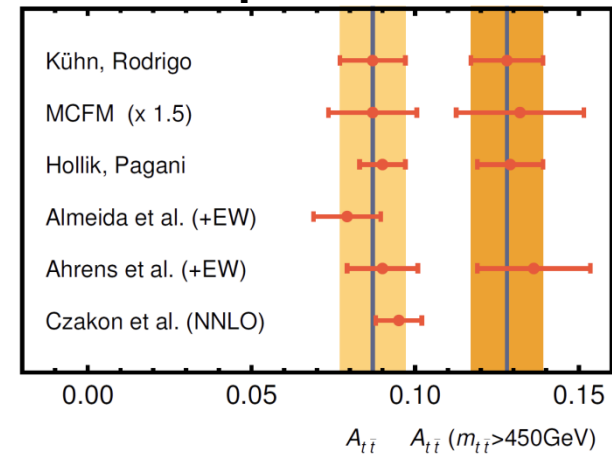
LHC: $A_C = \frac{N(\Delta|y| > 0) - N(\Delta|y| < 0)}{N(\Delta|y| > 0) + N(\Delta|y| < 0)}$

$\Delta y = y_t - y_{\bar{t}}$, $\Delta|y| = |y_t| - |y_{\bar{t}}|$



Kühn-Rodrigo, 1411.4675

SM predictions



Data is now consistent with the SM
(still 1.7 excess at CDF)

Models predicting larger asymmetries don't pass other phenomenological tests or are rather ad-hoc

TOP MASS

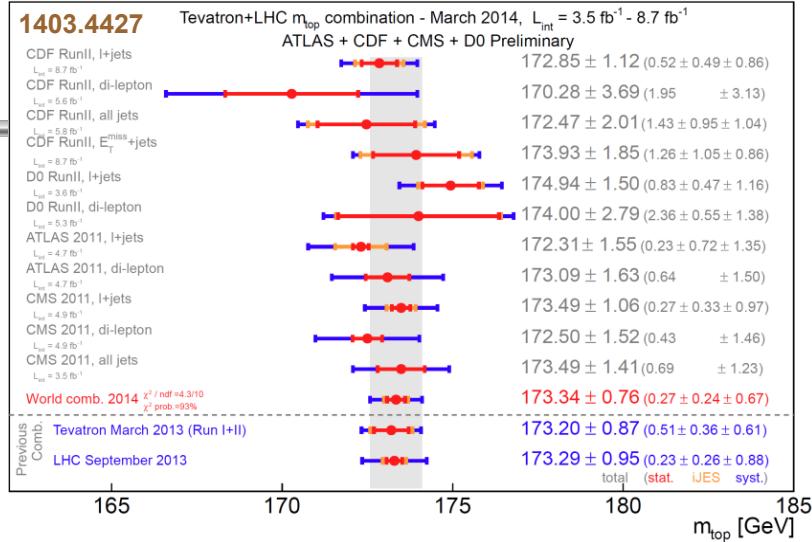
Monte Carlo mass:

$$M_t^{\text{MC}} = (173.34 \pm 0.76) \text{ GeV}$$

Lacks a proper QCD definition: $M_t^{\text{pole}} = M_t^{\text{MC}} + \Delta M_t^{\text{th}}$

$$|\Delta M_t^{\text{th}}| \approx \mathcal{O}(1 \text{ GeV})$$

Hoang-Stewart, 0808.0222

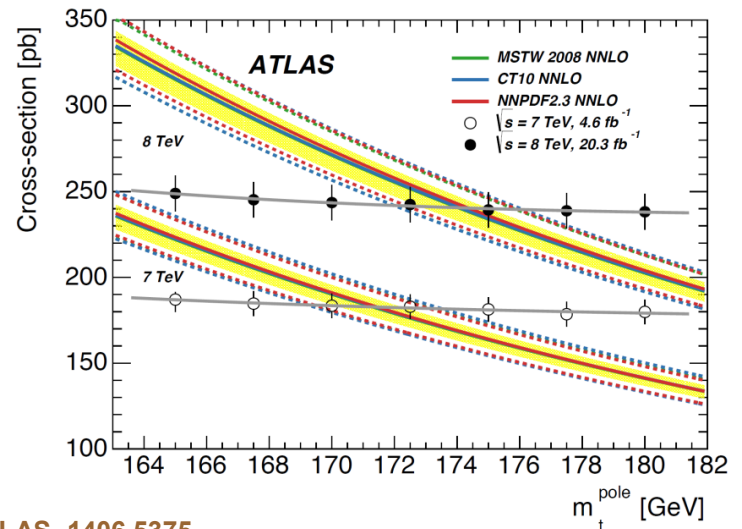
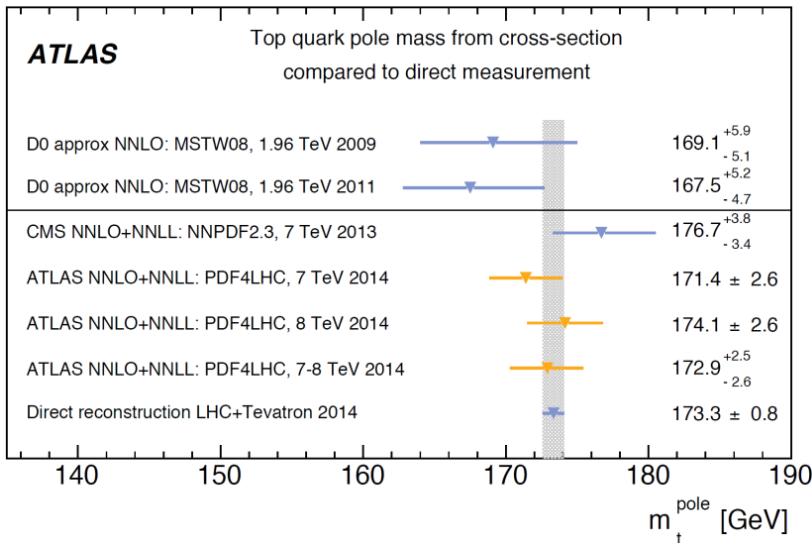


Cross section:

$\sigma_{t\bar{t}}$ NNLO+NNLL

Well-defined mass

Czakon et al, Bärnreuther et al, Cacciari et al



ATLAS, 1406.5375



Possible Improvements:

- Differential distribution in** $\rho_s = m_0 / \sqrt{s_{t\bar{t}+j}}$

Alioli et al, 1303.6415

$$R(m_t^{\text{pole}}, \rho_s) \equiv \frac{1}{\sigma_{t\bar{t}+1\text{jet}}} \frac{d\sigma_{t\bar{t}+1\text{jet}}}{d\rho_s}$$

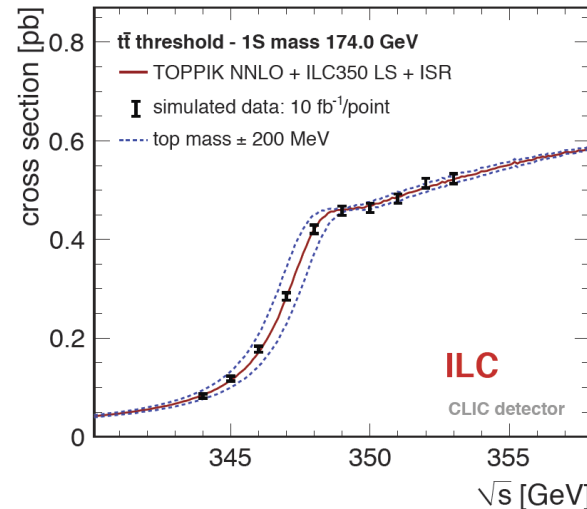
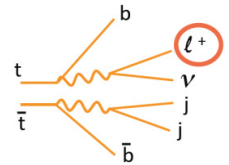
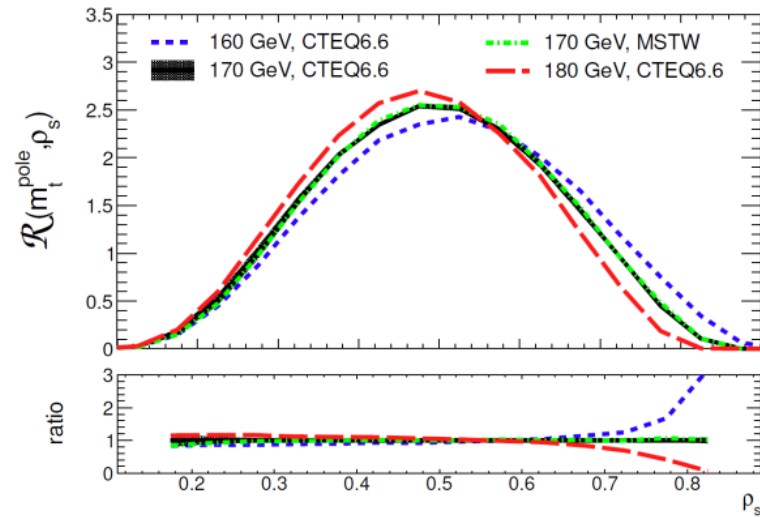
- Weight function method:** Lepton energy distribution

Kawabata et al, 1405.2395

- $\sigma(e^+e^- \rightarrow t\bar{t})_{\text{threshold}}$

Hoang et al, Beneke et al, Ruiz-Femenía, Martínez-Miquel

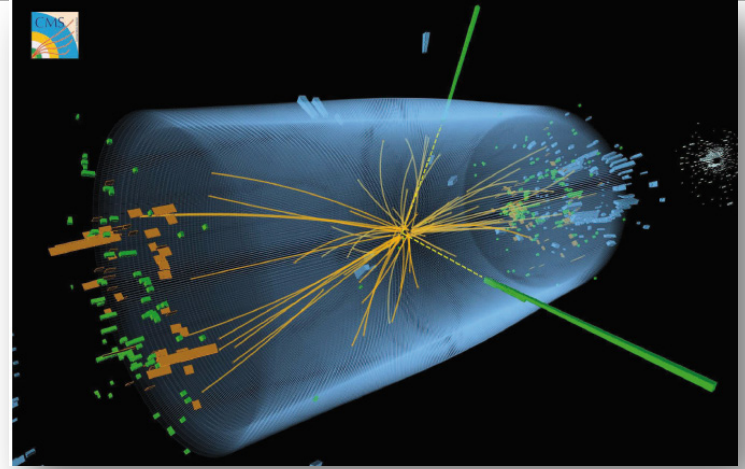
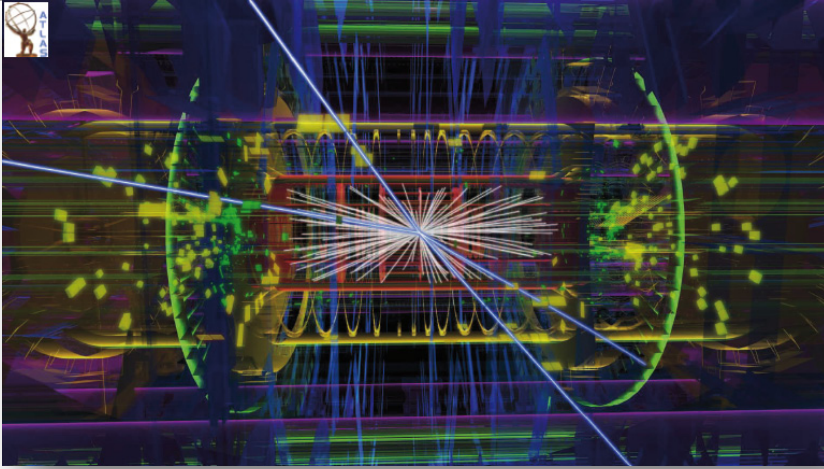
$$\delta m_t < 100 \text{ MeV}$$



García-García

Precision measurement needed to test the EW theory

A New Higgs-like Boson

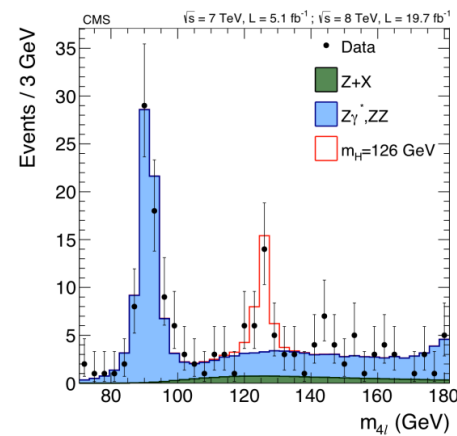
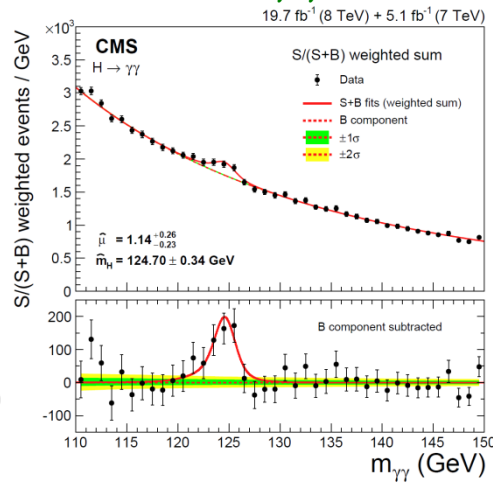
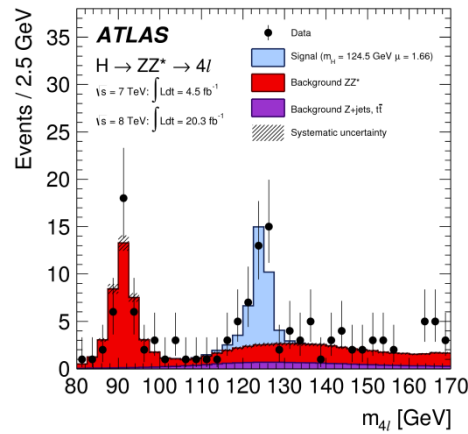
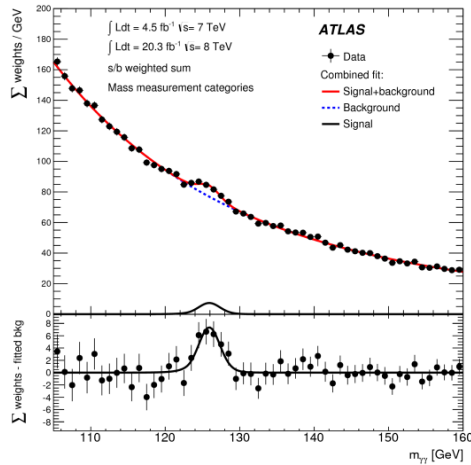


$H \rightarrow \gamma\gamma$

$H \rightarrow ZZ^* \rightarrow 4l$

$H \rightarrow \gamma\gamma$

$H \rightarrow ZZ^* \rightarrow 4l$



$$M_H^{\text{ATLAS}} = (125.36 \pm 0.37 \pm 0.18) \text{ GeV}$$

$$M_H^{\text{CMS}} = \left(\begin{matrix} 125.03 & +0.26 & +0.13 \\ & -0.27 & -0.15 \end{matrix} \right) \text{ GeV}$$



Beautiful Discovery

Boson (J = 0)

Fermions = Matter ; Bosons = Forces

- **Fundamental Boson:** New interaction which is not gauge
- **Composite Boson:** New underlying dynamics

If New Physics exists at Λ_{NP}

$$\delta M_H^2 \sim \frac{g^2}{(4\pi)^2} \Lambda_{\text{NP}}^2 \log\left(\frac{\Lambda_{\text{NP}}^2}{M_H^2}\right)$$

Which symmetry keeps M_H away from Λ_{NP} ?

- **Fermions:** Chiral Symmetry
- **Gauge Bosons:** Gauge Symmetry
- **Scalar Bosons:** Supersymmetry, Scale/Conformal Symmetry ... ?

Symmetries & Mass Scales

Fermions: $\psi_{L,R} \longrightarrow e^{i\alpha_{L,R}} \psi_{L,R}$ **Chiral symmetry**

$$\mathcal{L}_\psi = \bar{\psi} (i\partial - m_\psi) \psi = \bar{\psi}_L i\partial \psi_L + \bar{\psi}_R i\partial \psi_R - m_\psi (\bar{\psi}_L \psi_R + \bar{\psi}_R \psi_L)$$

Symmetry recovered at $m_\psi = 0$ \longrightarrow $\delta m_\psi \propto m_\psi$

Vectors: $A_\mu \longrightarrow A_\mu + \partial_\mu \theta$ **Gauge symmetry**

$$\mathcal{L}_A = -\frac{1}{4} F_{\mu\nu} F^{\mu\nu} + \frac{1}{2} m_A^2 A_\mu A^\mu$$

Symmetry recovered at $m_A = 0$ \longrightarrow $\delta m_A^2 \propto m_A^2$

Scalars: $\mathcal{L}_\phi = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m_\phi^2 \phi^2$ **Any symmetry?**

No additional symmetry at $m_\phi = 0$ \longrightarrow $\delta m_\phi^2 \propto M^2$ ($M = \text{any scale}$)

Symmetries & Mass Scales

Scalars: $\mathcal{L}_\phi = \frac{1}{2} \partial_\mu \phi \partial^\mu \phi - \frac{1}{2} m_\phi^2 \phi^2$ **Any symmetry?**

No additional symmetry at $m_\phi = 0$ \rightarrow $\delta m_\phi^2 \propto M^2$ ($M = \text{any scale}$)

- **Shift symmetry:** $\phi \rightarrow \phi + c$

Pseudo-Goldstone Boson

- **Scale symmetry:** $x \rightarrow x/\lambda$, $\phi(x) \rightarrow \lambda \phi(x/\lambda)$

$$M = 0 \quad , \quad \forall M$$

Conformal Invariance. Dilaton



Possible Scenarios of EWSB

1. SM scalar:

Favoured by EW precision tests

2. Alternative perturbative EWSB:

Scalar Doublets and Singlets

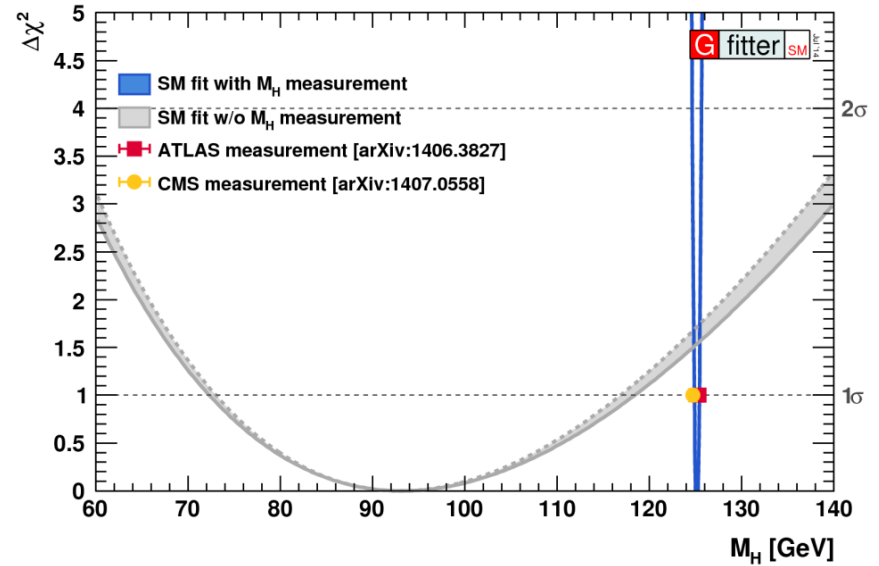
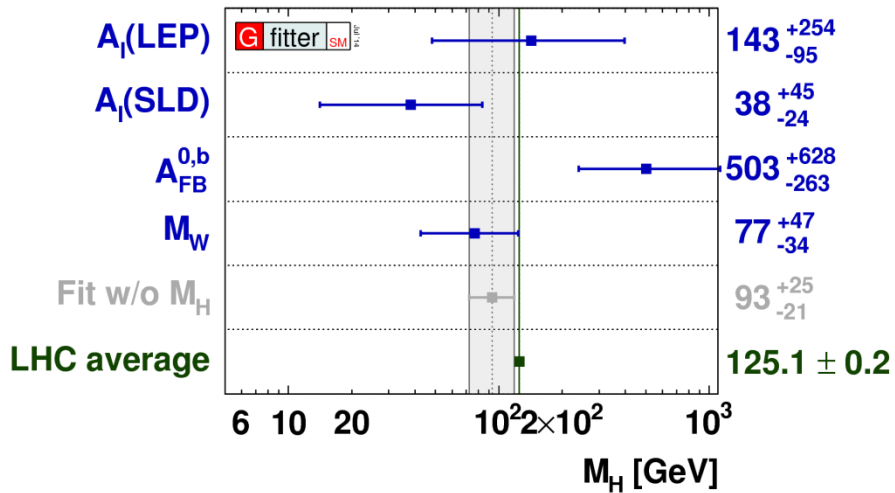
$$\rho_{\text{tree}} = \frac{M_W^2}{M_Z^2 c_W^2} = \frac{\sum_k v_k^2 [T_k(T_k + 1) - Y_k^2]}{2 \sum_k v_k^2 Y_k^2}$$

3. Dynamical (non-perturbative) EWSB:

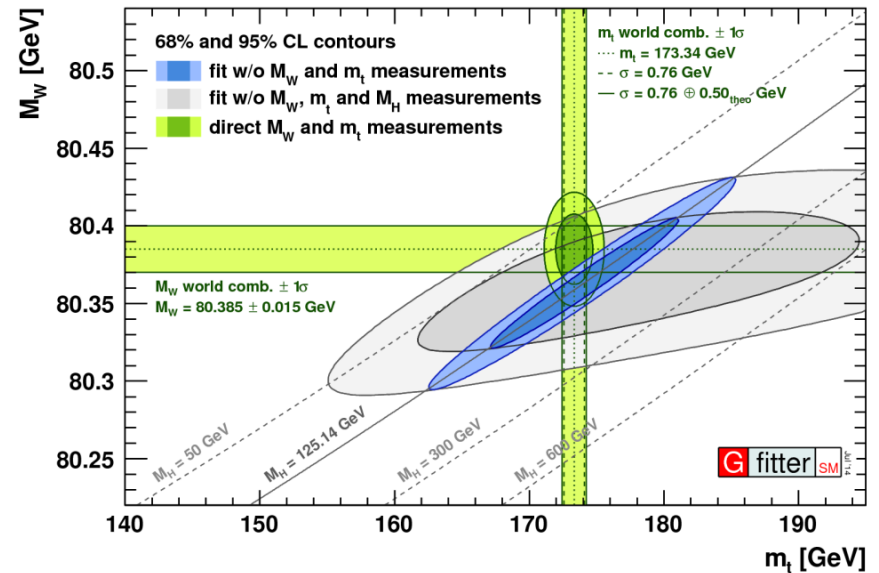
Pseudo-Goldstone Boson

Scalar Resonance

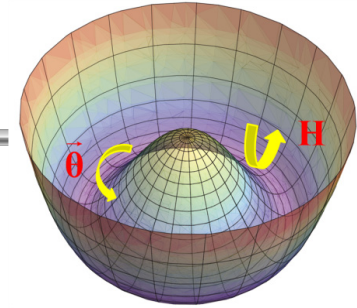
SM Higgs



**Favoured by
EW precision tests**



SM Scalar Potential



$$V(\Phi) - V_0 = \lambda \left(|\Phi|^2 - \frac{v^2}{2} \right)^2 = \frac{1}{2} M_H^2 H^2 \left(1 + \frac{H}{v} + \frac{H^2}{4v^2} \right)$$

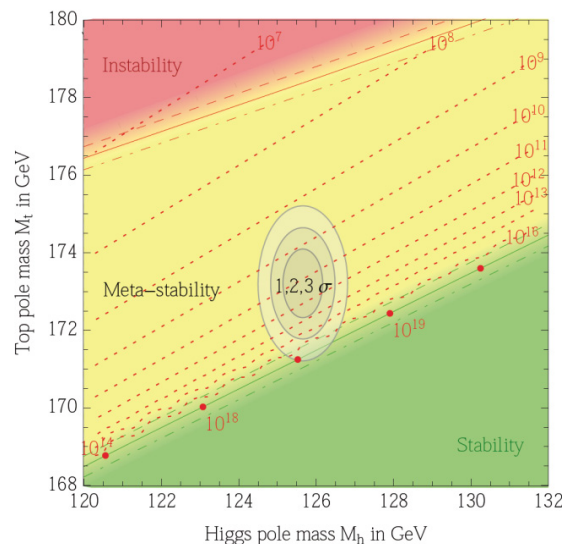
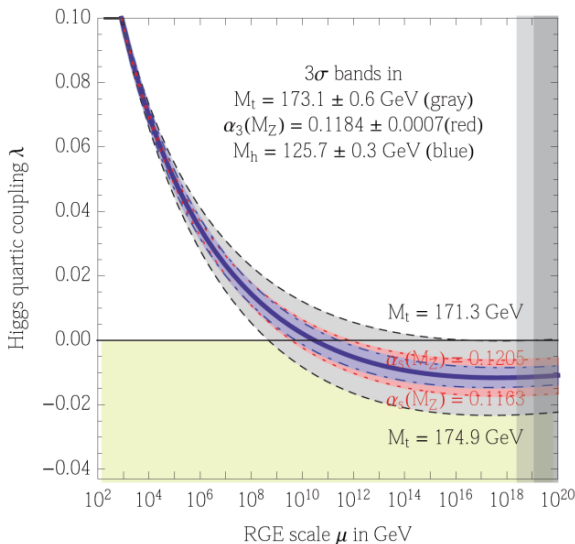
$$M_H = 125.14 \pm 0.24 \quad \rightarrow \quad \lambda = \frac{M_H^2}{2v^2} = 0.13$$

Loop corrections:

$$M_H^2 = 2\lambda(\mu) v^2 + \frac{2y_t^2 v^2}{(4\pi)^2} \left[2\lambda + 3(\lambda - y_t^2) \log(m_t^2 / \mu^2) \right] + \dots$$

Vacuum stability: $\lambda(\Lambda) \geq 0$

Meta-stable vacuum



$$\Lambda = M_{\text{Planck}}$$

↓

$$M_H > 129.1 \pm 1.5 \text{ GeV}$$

$$M_t < 171.53 \pm 0.42 \text{ GeV}$$

Buttazzo et al, 1307.3536

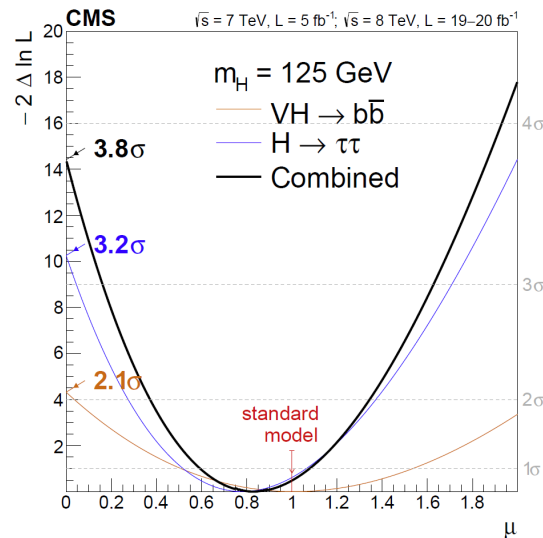
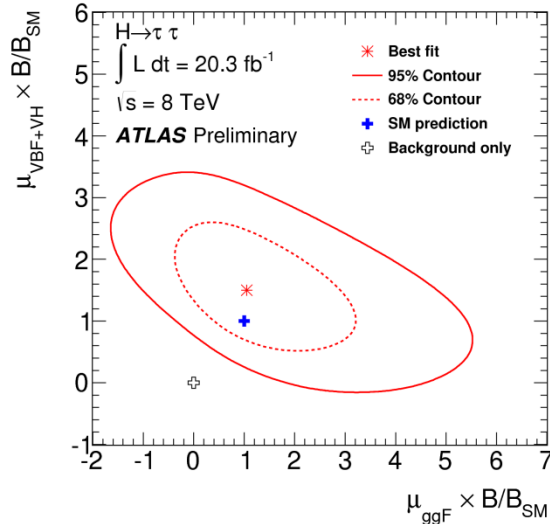
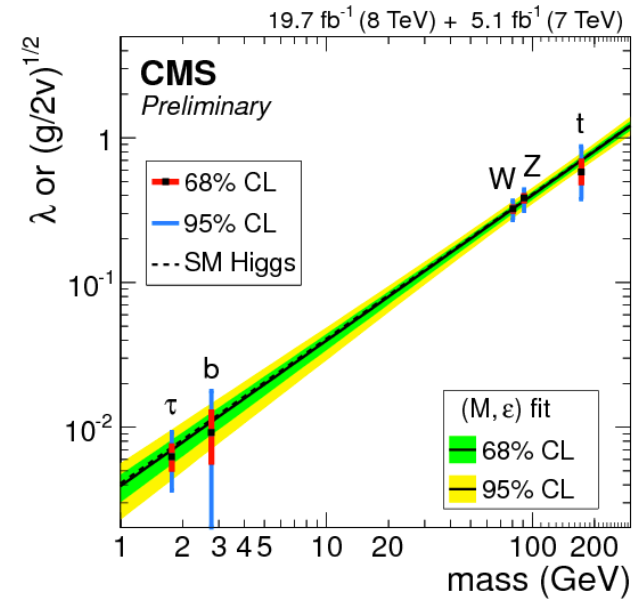
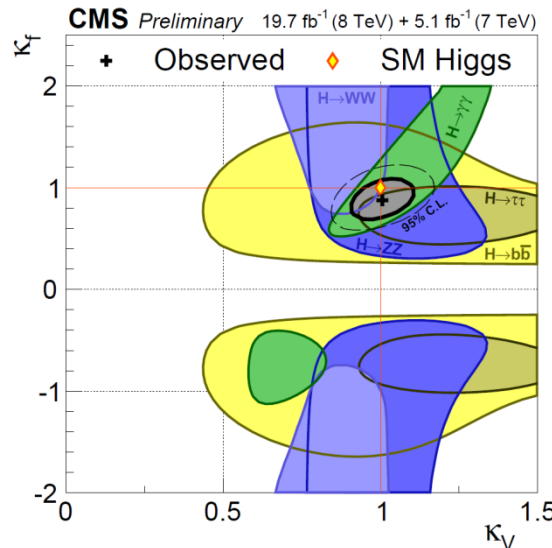
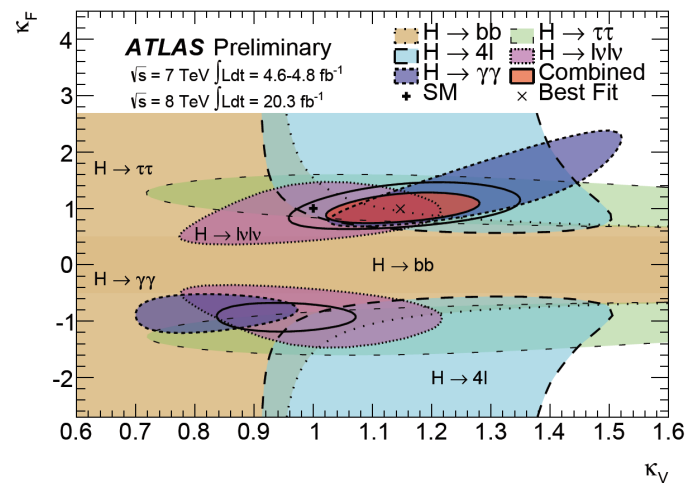
SM interactions only

Strong sensitivity to New Physics



H(125) Couplings are SM-like

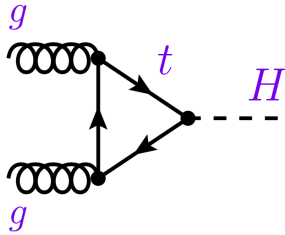
$$\kappa_i \equiv g_i / g_i^{\text{SM}}$$



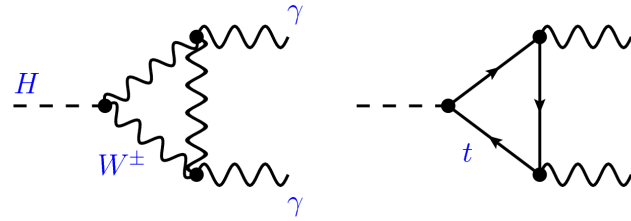
Strong evidence for H coupling to τ and b



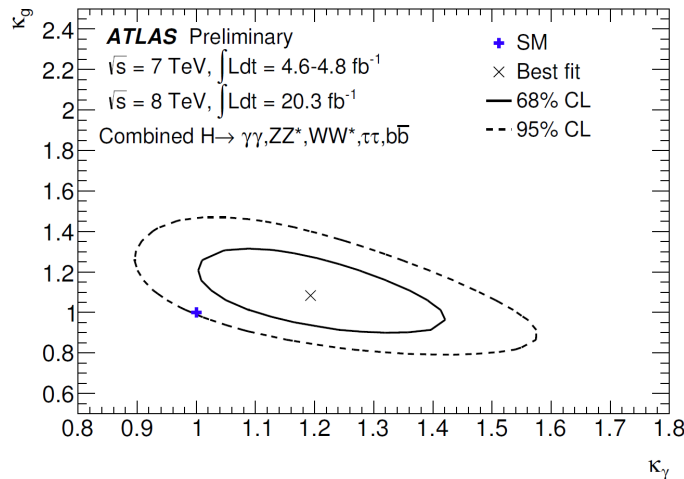
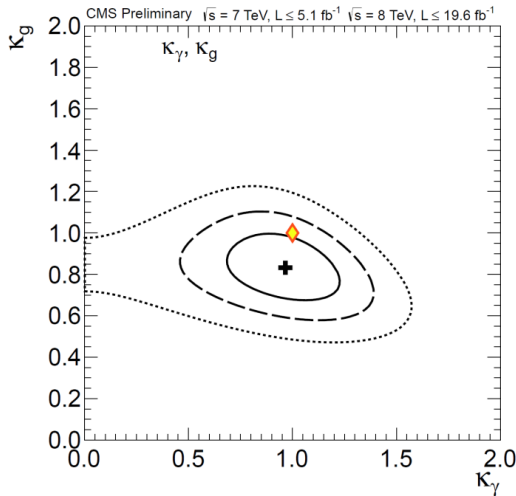
Strong (indirect) evidence for H coupling to t



Dominant Production Mechanism

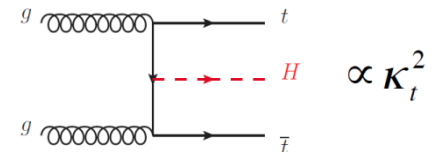


$$\Gamma \sim |1 - 0.21|^2$$

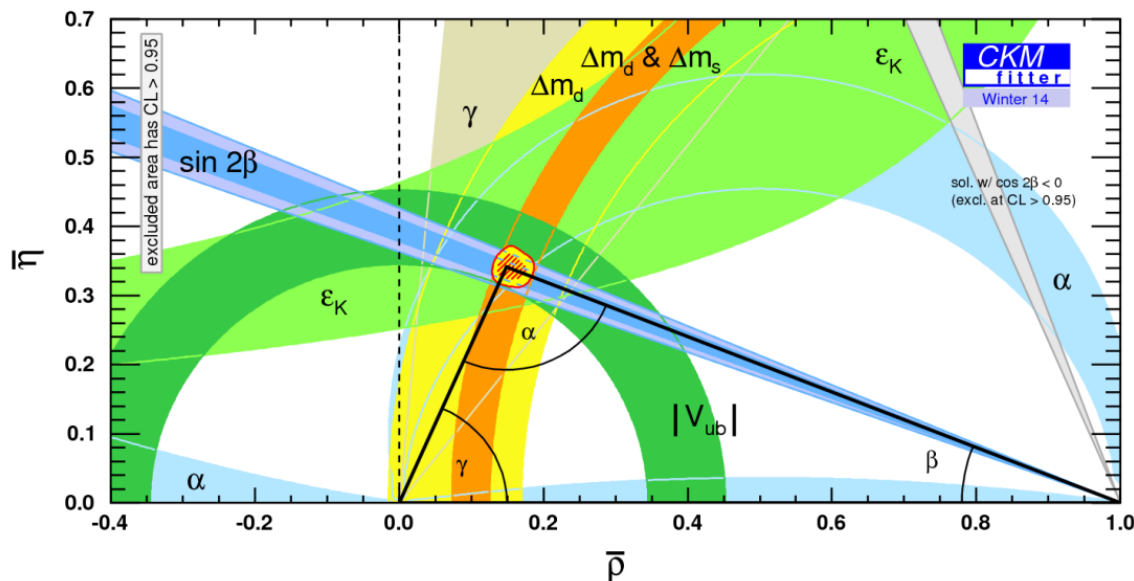


| $H \rightarrow \gamma\gamma$ | Signal Strength |
|------------------------------|-----------------|
| ATLAS | 1.17 ± 0.27 |
| CMS | 1.13 ± 0.24 |

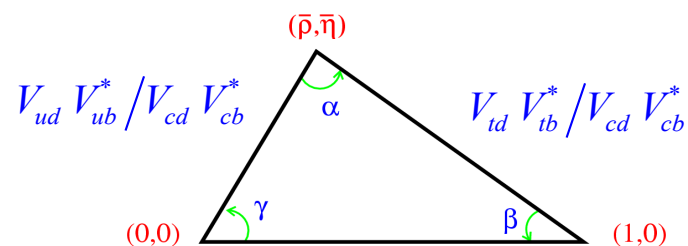
Direct (tree-level) sensitivity through $t\bar{t}H$



Quark Mixing



$$V_{ud} V_{ub}^* + V_{cd} V_{cb}^* + V_{td} V_{tb}^* = 0$$



$$V = \begin{bmatrix} 1 - \lambda^2/2 & \lambda & A\lambda^3(\rho - i\eta) \\ -\lambda & 1 - \lambda^2/2 & A\lambda^2 \\ A\lambda^3(1 - \rho - i\eta) & -A\lambda^2 & 1 \end{bmatrix} + \mathcal{O}(\lambda^4)$$

UT_{fit}

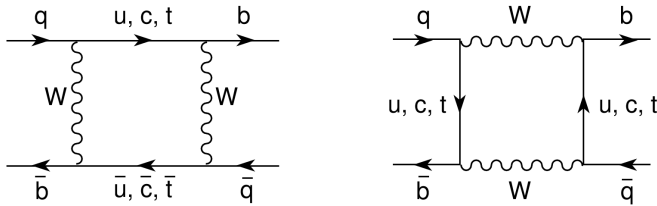
$$\bar{\eta} \equiv \eta \left(1 - \frac{1}{2}\lambda^2\right) = 0.351 \pm 0.014$$

$$\bar{\rho} \equiv \rho \left(1 - \frac{1}{2}\lambda^2\right) = 0.132 \pm 0.023$$

$$A = 0.821 \pm 0.012 \quad ; \quad \lambda = 0.2254 \pm 0.0006$$

Successful CKM Mechanism (Tree / Loop / CP-c / CP-v)

Bounds on New Flavour Physics



$$L_{\text{eff}} = L_{\text{SM}} + \sum_{D>4} \sum_k \frac{c_k^{(D)}}{\Lambda_{\text{NP}}^{D-4}} O_k^{(D)}$$

Isidori, 1302.0661

| Operator | Bounds on Λ in TeV ($c_{\text{NP}} = 1$) | | Bounds on c_{NP} ($\Lambda = 1$ TeV) | | Observables |
|----------------------------------|--|-------------------|--|-----------------------|---------------------------------|
| | Re | Im | Re | Im | |
| $(\bar{s}_L \gamma^\mu d_L)^2$ | 9.8×10^2 | 1.6×10^4 | 9.0×10^{-7} | 3.4×10^{-9} | $\Delta m_K; \epsilon_K$ |
| $(\bar{s}_R d_L)(\bar{s}_L d_R)$ | 1.8×10^4 | 3.2×10^5 | 6.9×10^{-9} | 2.6×10^{-11} | $\Delta m_K; \epsilon_K$ |
| $(\bar{c}_L \gamma^\mu u_L)^2$ | 1.2×10^3 | 2.9×10^3 | 5.6×10^{-7} | 1.0×10^{-7} | $\Delta m_D; q/p , \phi_D$ |
| $(\bar{c}_R u_L)(\bar{c}_L u_R)$ | 6.2×10^3 | 1.5×10^4 | 5.7×10^{-8} | 1.1×10^{-8} | $\Delta m_D; q/p , \phi_D$ |
| $(\bar{b}_L \gamma^\mu d_L)^2$ | 6.6×10^2 | 9.3×10^2 | 2.3×10^{-6} | 1.1×10^{-6} | $\Delta m_{B_d}; S_{\psi K_S}$ |
| $(\bar{b}_R d_L)(\bar{b}_L d_R)$ | 2.5×10^3 | 3.6×10^3 | 3.9×10^{-7} | 1.9×10^{-7} | $\Delta m_{B_d}; S_{\psi K_S}$ |
| $(\bar{b}_L \gamma^\mu s_L)^2$ | 1.4×10^2 | 2.5×10^2 | 5.0×10^{-5} | 1.7×10^{-5} | $\Delta m_{B_s}; S_{\psi \phi}$ |
| $(\bar{b}_R s_L)(\bar{b}_L s_R)$ | 4.8×10^2 | 8.3×10^2 | 8.8×10^{-6} | 2.9×10^{-6} | $\Delta m_{B_s}; S_{\psi \phi}$ |

- Generic flavour structure [$c_{\text{NP}} \sim \mathcal{O}(1)$] ruled out at the TeV scale
- $\Lambda_{\text{NP}} \sim 1$ TeV requires c_{NP} to inherit the strong SM suppressions (GIM)

Minimal Flavour Violation: The up and down Yukawa matrices are the only source of quark-flavour symmetry breaking

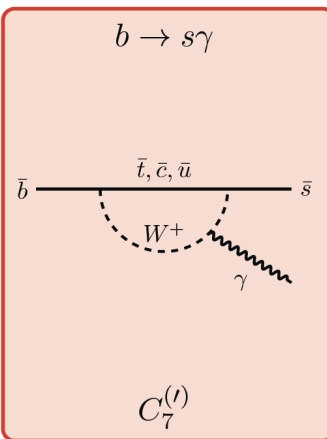
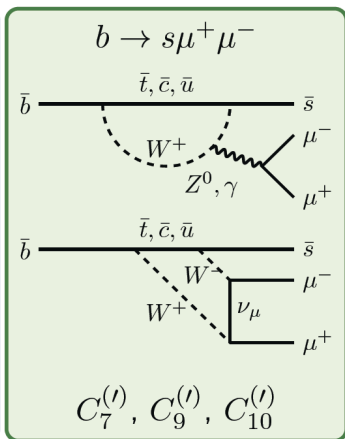
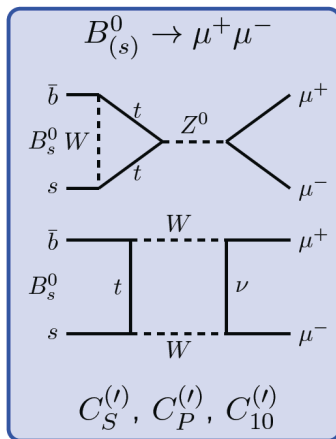
D'Ambrosio et al, Buras et al

Rare Decays

Albrecht

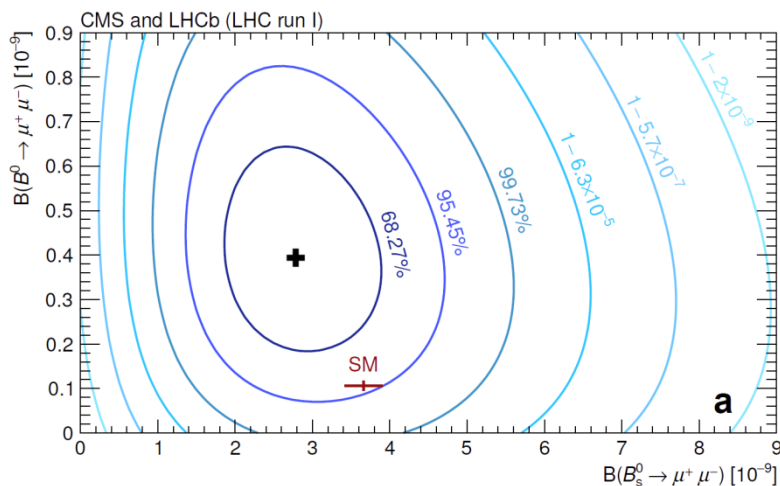
$$H_{\text{eff}} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i \left[\underbrace{C_i(\mu) O_i(\mu)}_{\text{left-handed part}} + \underbrace{C'_i(\mu) O'_i(\mu)}_{\substack{\text{right-handed part} \\ \text{suppressed in SM}}} \right]$$

- i=1,2 Tree
- i=3-6,8 Gluon penguin
- i=7 Photon penguin
- i=9,10 Electroweak penguin
- i=S Higgs (scalar) penguin
- i=P Pseudoscalar penguin



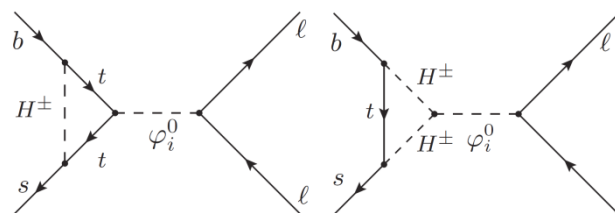
SM

- No tree-level contribution
- Strong CKM suppression



$B \rightarrow \mu^+ \mu^-$ sensitive to
(pseudo) scalar contributions
additional Higgs bosons

Li-Lu-A.P., 1404.5865



$B_d \rightarrow K^{*0} \mu^+ \mu^-$

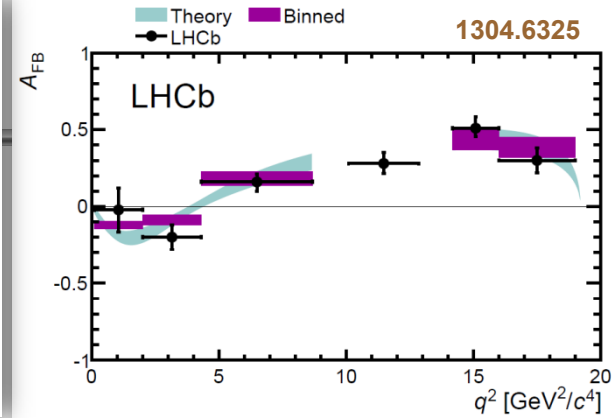
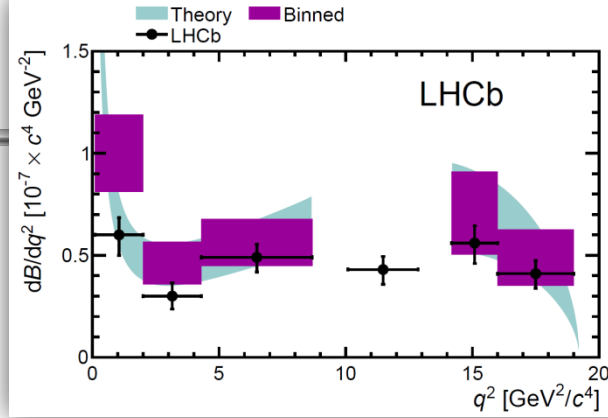
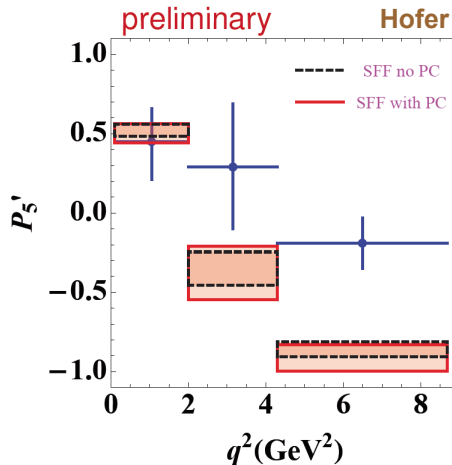
Phen. analysis with “clean observables”

(FF independent)

Descotes-Genon et al

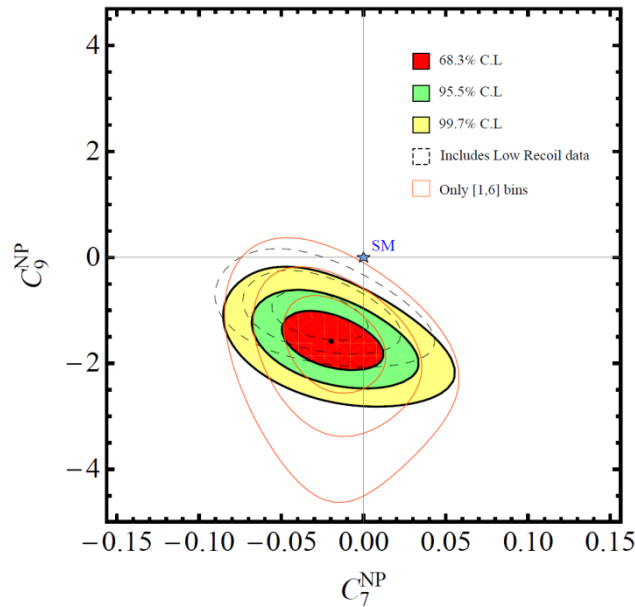
P'5 Anomaly

| Observable | Experiment | SM prediction | Pull |
|-------------------------------------|-------------------------|----------------------------|------|
| $\langle P_2 \rangle_{[0,1,2]}$ | $0.03^{+0.14}_{-0.15}$ | $0.172^{+0.020}_{-0.021}$ | -1.0 |
| $\langle P_2 \rangle_{[2,4,3]}$ | $0.50^{+0.00}_{-0.07}$ | $0.234^{+0.060}_{-0.086}$ | +2.9 |
| $\langle P_2 \rangle_{[4,3,8,68]}$ | $-0.25^{+0.07}_{-0.08}$ | $-0.407^{+0.049}_{-0.037}$ | +1.7 |
| $\langle P_2 \rangle_{[1,6]}$ | $0.33^{+0.11}_{-0.12}$ | $0.084^{+0.060}_{-0.078}$ | +1.8 |
| $\langle P'_5 \rangle_{[0,1,2]}$ | $0.45^{+0.21}_{-0.24}$ | $0.533^{+0.033}_{-0.041}$ | -0.4 |
| $\langle P'_5 \rangle_{[2,4,3]}$ | $0.29^{+0.40}_{-0.39}$ | $-0.334^{+0.097}_{-0.113}$ | +1.6 |
| $\langle P'_5 \rangle_{[4,3,8,68]}$ | $-0.19^{+0.16}_{-0.16}$ | $-0.872^{+0.053}_{-0.041}$ | +4.0 |
| $\langle P'_5 \rangle_{[1,6]}$ | $0.21^{+0.20}_{-0.21}$ | $-0.349^{+0.088}_{-0.100}$ | +2.5 |



Fit with “New Physics” effective operators

Descotes-Genon et al, 1307.5683



$$O_7 = \frac{e}{(16\pi)^2} m_b (\bar{s} \sigma_{\mu\nu} P_R b) F^{\mu\nu}$$

$$O_9 = \frac{e^2}{(16\pi)^2} m_b (\bar{s} \gamma_\mu P_L b) (\bar{\ell} \gamma^\mu \ell)$$

New Physics?

Altmannshofer-Straub, Beaujean et al,
Descotes-Genon et al, Horgan et al.

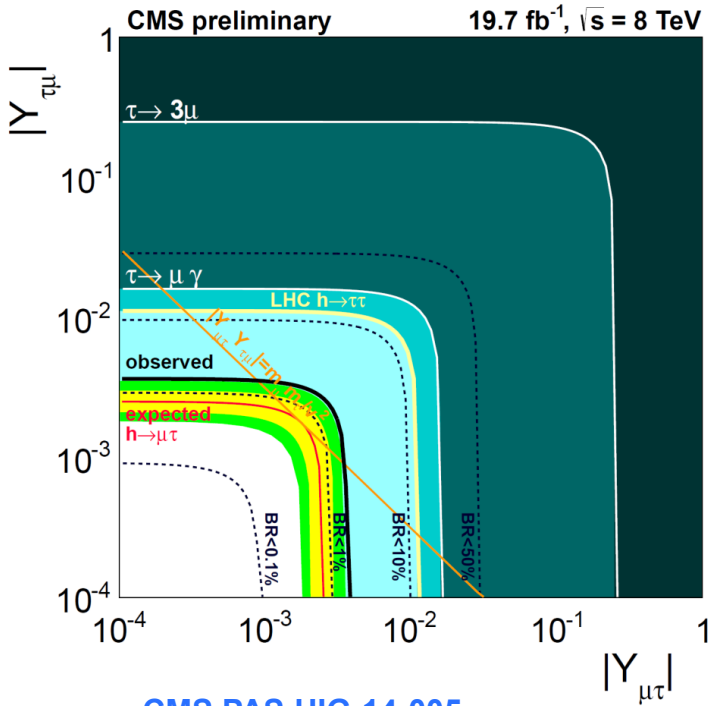
Hadronic uncertainties?

Jäger & Martín-Camalich, Zwicky et al

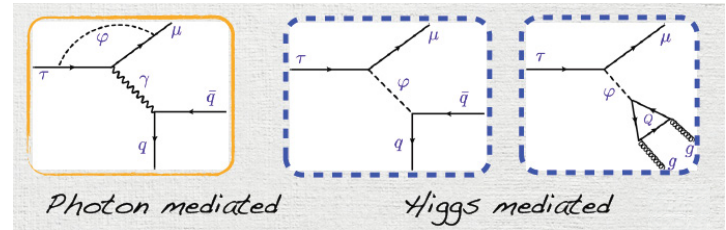
Flavour-Violating Higgs Couplings

Blankenburg et al, Celis et al,
Harnik et al, Davidson-Verdier,
Kopp-Nardecchia

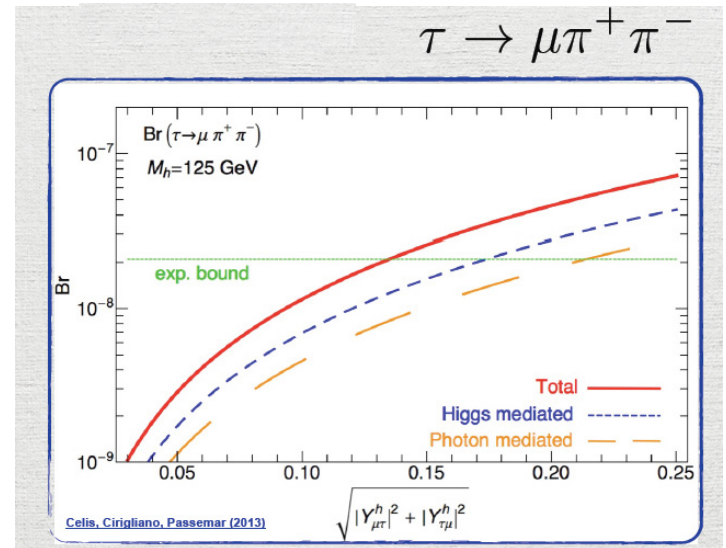
$$L = -h \left\{ Y_{e\mu} \bar{e}_L \mu_R + Y_{e\tau} \bar{e}_L \tau_R + Y_{\mu\tau} \bar{\mu}_L \tau_R + \dots \right\}$$



$$\text{Br}(H \rightarrow \mu\tau) < 1.57\% \quad (95\% \text{ CL})$$



Celis



Two-Higgs Doublet Models

5 scalar fields: $H^\pm, \phi_i^0 = (h, H, A)$ [3x3 mixing R_{ij}] $v = \sqrt{v_1^2 + v_2^2}$, $\tan \beta = v_2/v_1$

$$g_{hVV}^2 + g_{HVV}^2 + g_{AVV}^2 = \left(g_{hVV}^{\text{SM}}\right)^2$$

CP-conserving potential: $R = \begin{bmatrix} \cos \tilde{\alpha} & \sin \tilde{\alpha} & 0 \\ -\sin \tilde{\alpha} & \cos \tilde{\alpha} & 0 \\ 0 & 0 & 1 \end{bmatrix}$ $g_{\phi_i^0 VV} / g_{\phi_i^0 VV}^{\text{SM}} = R_{i1} = \cos \tilde{\alpha} \equiv \sin(\beta - \alpha)$

Yukawas: $L_Y = -\bar{Q}'_L (\Gamma_1 \phi_1 + \Gamma_2 \phi_2) d'_R + \dots$ $\xrightarrow{\text{EWSB}}$ $L_Y = -\frac{\sqrt{2}}{v} \bar{Q}'_L (M'_d \Phi_1 + Y'_d \Phi_2) d'_R + \dots$

M'_f & Y'_f **unrelated** (not simultaneously diagonal) $\xrightarrow{\hspace{2cm}}$ **FCNCs**

Solutions: (same for u_R and ℓ_R Yukawas)

- **Natural Flavour Conservation:** $\Gamma_1 = 0$ or $\Gamma_2 = 0$ (Z_2 models) Glashow-Weinberg...
- **Alignment:** $\Gamma_1 \propto \Gamma_2$ $\xrightarrow{\hspace{1cm}}$ $Y_{d,l} = \zeta_{d,l} M_{d,l}$, $Y_u = \zeta_u^* M_u$ AP-Tuzón
- **BGL Models:** “controlled” FCNC (symmetries) Branco et al

Z₂ 2HDMs

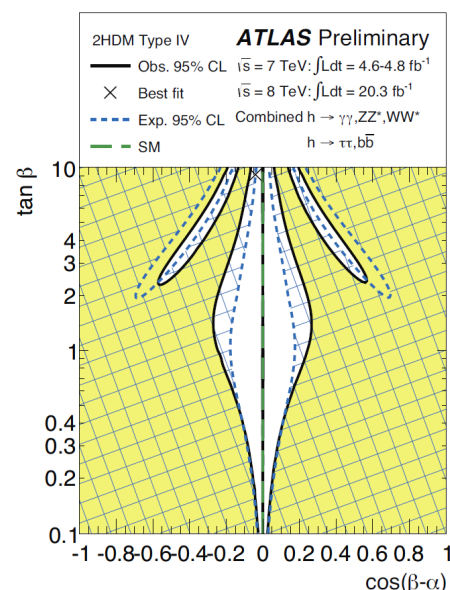
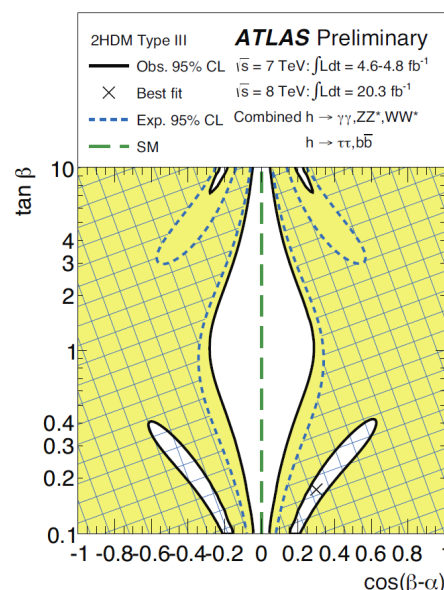
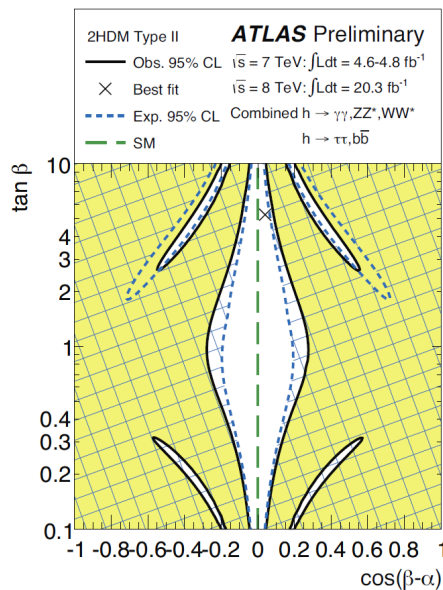
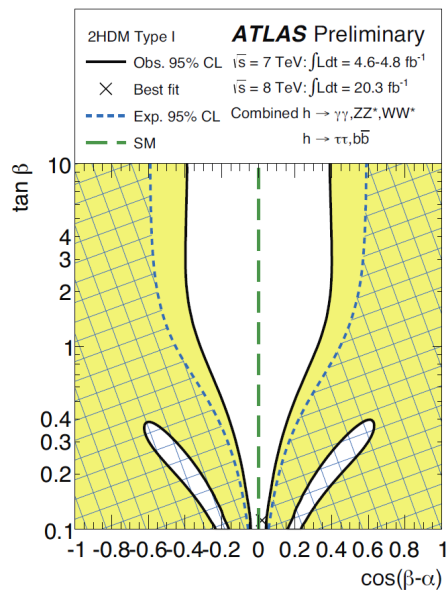
5 scalar fields: $H^\pm, \phi_i^0 = (h, H, A)$ [3x3 mixing R_{ij}]

$$v = \sqrt{v_1^2 + v_2^2}, \quad \tan \beta = v_2/v_1$$

CP-conserving potential:

$$R = \begin{bmatrix} \cos \tilde{\alpha} & \sin \tilde{\alpha} & 0 \\ -\sin \tilde{\alpha} & \cos \tilde{\alpha} & 0 \\ 0 & 0 & 1 \end{bmatrix}$$

$$g_{VV\phi_i^0} / g_{VV\phi_i^0}^{\text{SM}} = R_{i1} = \cos \tilde{\alpha} \equiv \sin(\beta - \alpha)$$



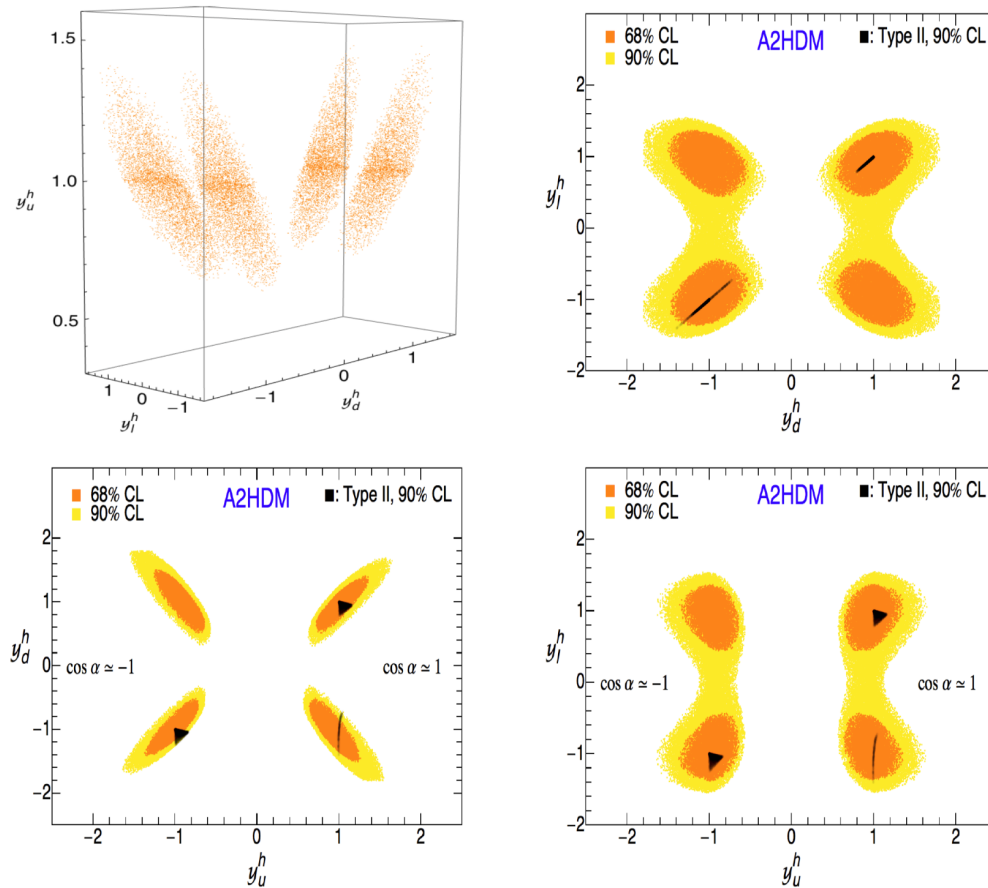
Yukawas: (SM units)

$$y_{d,l}^{\phi_i^0} = R_{i1} + (R_{i2} + i R_{i3}) \zeta_{d,l}$$

$$y_u^{\phi_i^0} = R_{i1} + (R_{i2} - i R_{i3}) \zeta_u^*$$

| Type | ζ_d | ζ_u | ζ_l |
|--------------------------|---------------|--------------|---------------|
| I | $\cot \beta$ | $\cot \beta$ | $\cot \beta$ |
| II | $-\tan \beta$ | $\cot \beta$ | $-\tan \beta$ |
| X (III, lepton specific) | $\cot \beta$ | $\cot \beta$ | $-\tan \beta$ |
| Y (IV, flipped) | $-\tan \beta$ | $\cot \beta$ | $\cot \beta$ |
| Inert | 0 | 0 | 0 |

Celis-Ilisie-AP, 1302.4022, 1310.7941



$$|\cos \tilde{\alpha}| > 0.80 \quad (90\% \text{ CL})$$

**General setting without FCNCs
& new sources of CP violation**

$$Y_{d,l} = \zeta_{d,l} M_{d,l} \quad , \quad Y_u = \zeta_u^* M_u$$

- Rich phenomenology @ LHC**

Altmannshofer et al, Barger et al, Celis et al, Cervero-Gerard, López-Val et al...

Many allowed possibilities

Search for light H^\pm, H, A

CP violation

- Flavour constraints fulfilled**

Celis et al, Jung et al, Li et al

- EDMs**

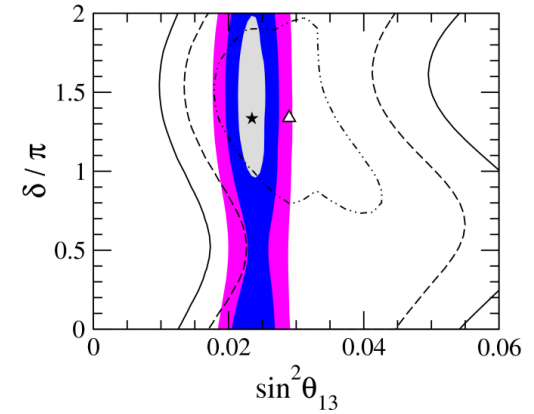
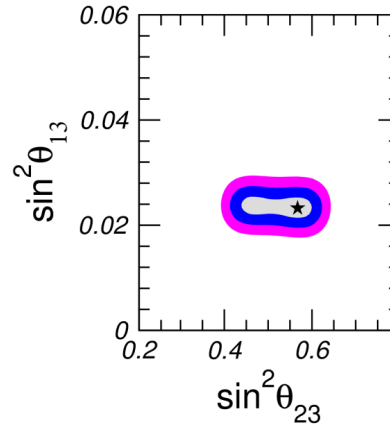
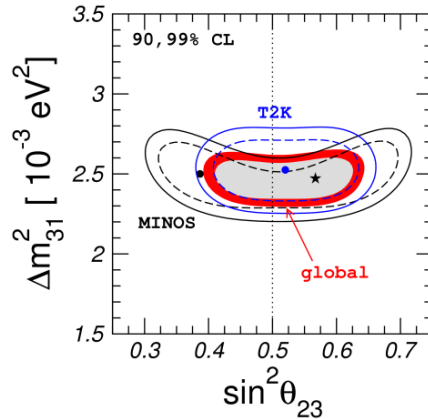
Jung-AP, 1308.6283

- Usual Z_2 models recovered in particular (CP-conserving) limits**

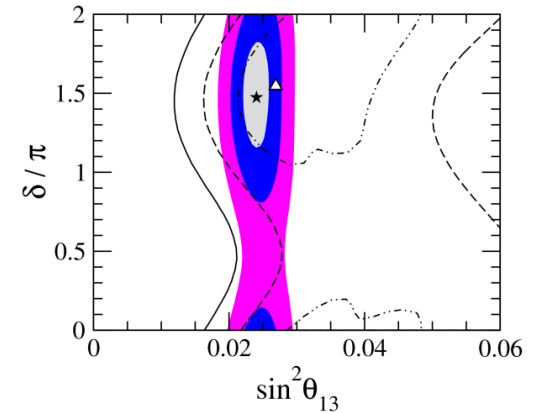
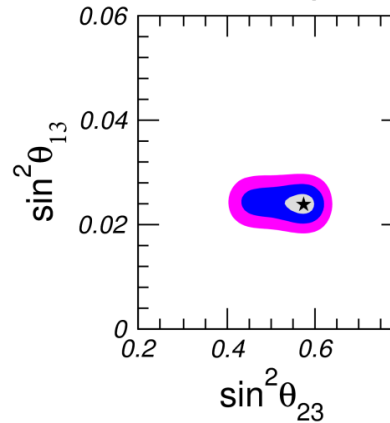
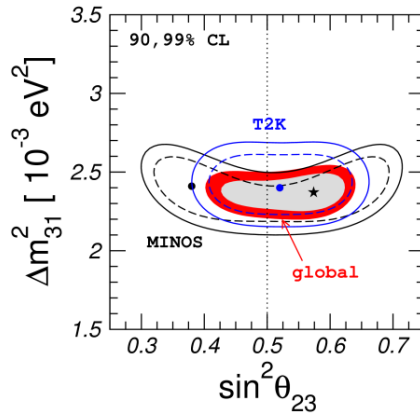
ν Oscillations

Forero et al, 1405.7540

NH



IH



González-García et al

NuFIT 2.0 (2014)

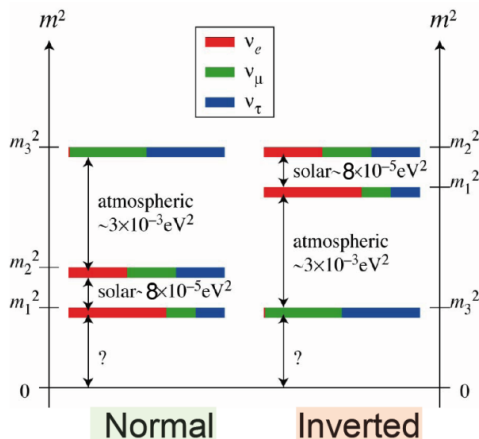
$$|U|_{3\sigma} = \begin{pmatrix} 0.801 \rightarrow 0.845 & 0.514 \rightarrow 0.580 & 0.137 \rightarrow 0.158 \\ 0.225 \rightarrow 0.517 & 0.441 \rightarrow 0.699 & 0.614 \rightarrow 0.793 \\ 0.246 \rightarrow 0.529 & 0.464 \rightarrow 0.713 & 0.590 \rightarrow 0.776 \end{pmatrix}$$

Flavour mixing is very different for quarks & leptons

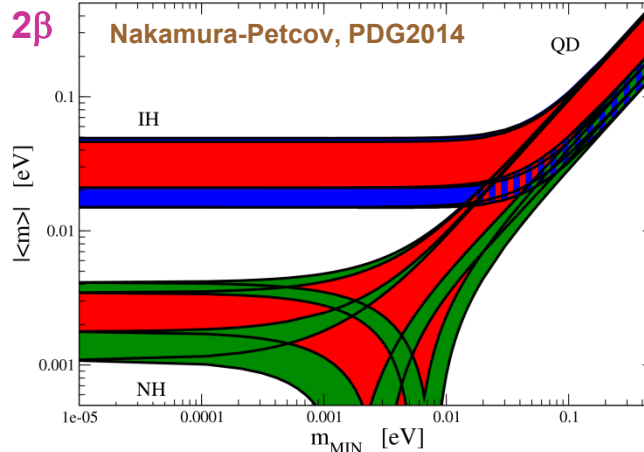
Open Questions in ν Physics

Mass Hierarchy

Blennow



Dirac / Majorana



$$\langle m \rangle = \left| m_1 U_{e1}^2 + m_2 U_{e2}^2 + m_3 U_{e3}^2 \right|$$

Mass Scale

Sterile ν_R ?

CP Violation

Flavour Symmetries

Leptogenesis

Low-E Effective Theory:

$$L = L_{SM} + \sum_d \frac{c_d}{\Lambda^{d-4}} O_d$$

1 $SU(2)_L \otimes U(1)_Y$ invariant operator with $d=5$

Weinberg

$$-\frac{c_{ij}}{\Lambda} \bar{L}_i \tilde{\phi} \tilde{\phi}^t L_j^c + \text{h.c.} \xrightarrow{\text{SSB}} -\frac{1}{2} \bar{\nu}_{iL} M_{ij} \nu_{jL}^c + \text{h.c.} \quad ; \quad M_{ij} \equiv \frac{c_{ij}}{\Lambda} v^2$$

Small Majorana Mass: $m_\nu > 0.05 \text{ eV}$ \rightarrow $\Lambda / c_{ij} < 10^{15} \text{ GeV}$



Desperately Seeking SUSY (Dulcinea)



In all the world there is no maiden fairer than the Empress of La Mancha, the peerless SUSY del Toboso

Your worship should bear in mind that SUSY is badly broken; got heavy through anomaly mediation



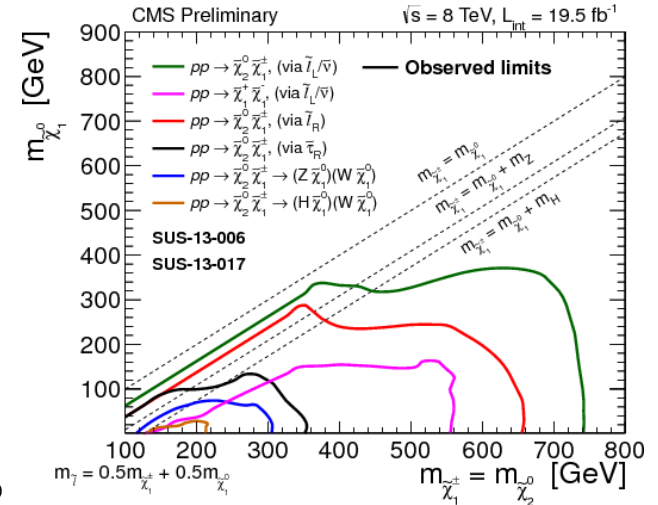
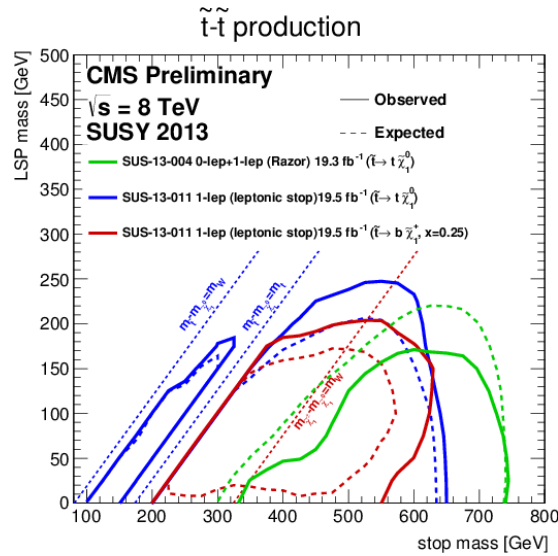
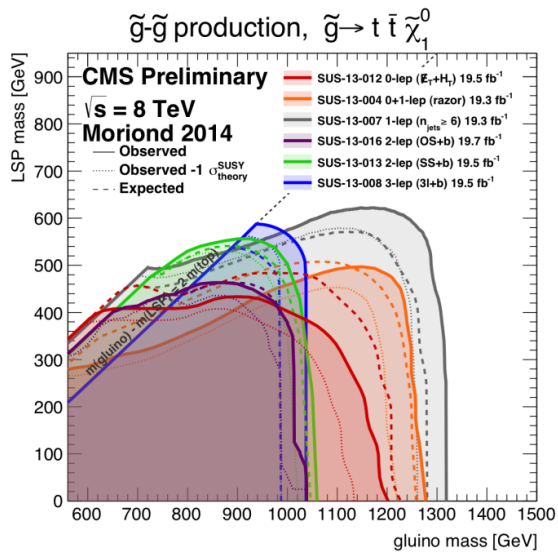
| Model | | e, μ, τ, γ | Jets | E_T^{miss} | $\int \mathcal{L} dt [\text{fb}^{-1}]$ | Mass limit | Reference | |
|---|---|---|----------------|---------------------|--|---|---|--|
| Inclusive Searches | MSUGRA/CMSSM | 0 | 2-6 jets | Yes | 20.3 | \tilde{q}, \tilde{g} 1.7 TeV | $m(\tilde{q})=m(\tilde{g})$ | |
| | MSUGRA/CMSSM | $1 e, \mu$ | 3-6 jets | Yes | 20.3 | \tilde{g} 1.2 TeV | any $m(\tilde{q})$ | |
| | MSUGRA/CMSSM | 0 | 7-10 jets | Yes | 20.3 | \tilde{g} 1.1 TeV | any $m(\tilde{q})$ | |
| | $\tilde{q}\tilde{q}, \tilde{q} \rightarrow q\tilde{\chi}_1^0$ | 0 | 2-6 jets | Yes | 20.3 | \tilde{q} 850 GeV | $m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(1^{\text{st}} \text{ gen. } \tilde{q})=m(2^{\text{nd}} \text{ gen. } \tilde{q})$ | |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow q\tilde{q}\tilde{\chi}_1^0$ | 0 | 2-6 jets | Yes | 20.3 | \tilde{g} 1.33 TeV | $m(\tilde{\chi}_1^0)=0 \text{ GeV}$ | |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq\tilde{\chi}_1^\pm \rightarrow qqW^\pm\tilde{\chi}_1^0$ | $1 e, \mu$ | 3-6 jets | Yes | 20.3 | \tilde{g} 1.18 TeV | $m(\tilde{\chi}_1^0)<200 \text{ GeV}, m(\tilde{\chi}^\pm)=0.5(m(\tilde{\chi}_1^0)+m(\tilde{g}))$ | |
| | $\tilde{g}\tilde{g}, \tilde{g} \rightarrow qq(\ell\ell/\ell\nu/\nu\nu)\tilde{\chi}_1^0$ | $2 e, \mu$ | 0-3 jets | - | 20.3 | \tilde{g} 1.12 TeV | $m(\tilde{\chi}_1^0)=0 \text{ GeV}$ | |
| | GMSB ($\tilde{\ell}$ NLSP) | $2 e, \mu$ | 2-4 jets | Yes | 4.7 | \tilde{g} 1.24 TeV | $\tan\beta<15$ | |
| | GMSB ($\tilde{\ell}$ NLSP) | $1-2 \tau + 0-1 \ell$ | 0-2 jets | Yes | 20.3 | \tilde{g} 1.6 TeV | $\tan\beta>20$ | |
| | GGM (bino NLSP) | 2γ | - | Yes | 20.3 | \tilde{g} 1.28 TeV | $m(\tilde{\chi}_1^0)>50 \text{ GeV}$ | |
| | GGM (wino NLSP) | $1 e, \mu + \gamma$ | - | Yes | 4.8 | \tilde{g} 619 GeV | $m(\tilde{\chi}_1^0)>50 \text{ GeV}$ | |
| | GGM (higgsino-bino NLSP) | γ | $1 b$ | Yes | 4.8 | \tilde{g} 900 GeV | $m(\tilde{\chi}_1^0)>220 \text{ GeV}$ | |
| | GGM (higgsino NLSP) | $2 e, \mu (Z)$ | 0-3 jets | Yes | 5.8 | \tilde{g} 690 GeV | $m(\text{NLSP})>200 \text{ GeV}$ | |
| | Gravitino LSP | 0 | mono-jet | Yes | 10.5 | \tilde{g} 645 GeV | $m(\tilde{G})>10^{-4} \text{ eV}$ | |
| | 3^{rd} gen. \tilde{g} med. | $\tilde{g} \rightarrow b\tilde{b}\tilde{\chi}_1^0$ | 0 | $3 b$ | Yes | 20.1 | \tilde{g} 1.25 TeV | $m(\tilde{\chi}_1^0)<400 \text{ GeV}$ |
| $\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^0$ | | 0 | 7-10 jets | Yes | 20.3 | \tilde{g} 1.1 TeV | $m(\tilde{\chi}_1^0)<350 \text{ GeV}$ | |
| $\tilde{g} \rightarrow t\tilde{t}\tilde{\chi}_1^\pm$ | | $0-1 e, \mu$ | $3 b$ | Yes | 20.1 | \tilde{g} 1.34 TeV | $m(\tilde{\chi}_1^0)<400 \text{ GeV}$ | |
| $\tilde{g} \rightarrow b\tilde{t}\tilde{\chi}_1^\pm$ | | $0-1 e, \mu$ | $3 b$ | Yes | 20.1 | \tilde{g} 1.3 TeV | $m(\tilde{\chi}_1^0)<300 \text{ GeV}$ | |
| 3^{rd} gen. squarks direct production | | $\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow b\tilde{\chi}_1^0$ | 0 | $2 b$ | Yes | 20.1 | \tilde{b}_1 100-620 GeV | $m(\tilde{\chi}_1^0)<90 \text{ GeV}$ |
| | $\tilde{b}_1\tilde{b}_1, \tilde{b}_1 \rightarrow t\tilde{\chi}_1^\pm$ | $2 e, \mu (SS)$ | $0-3 b$ | Yes | 20.3 | \tilde{b}_1 275-440 GeV | $m(\tilde{\chi}_1^\pm)=2 m(\tilde{\chi}_1^0)$ | |
| | $\tilde{t}_1\tilde{t}_1 (\text{light}), \tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$ | $1-2 e, \mu$ | $1-2 b$ | Yes | 4.7 | \tilde{t}_1 110-167 GeV | $m(\tilde{\chi}_1^0)=55 \text{ GeV}$ | |
| | $\tilde{t}_1\tilde{t}_1 (\text{light}), \tilde{t}_1 \rightarrow Wb\tilde{\chi}_1^0$ | $2 e, \mu$ | 0-2 jets | Yes | 20.3 | \tilde{t}_1 130-210 GeV | $m(\tilde{\chi}_1^0)=m(\tilde{t}_1)-m(W)-50 \text{ GeV}, m(\tilde{t}_1)<m(\tilde{\chi}_1^\pm)$ | |
| | $\tilde{t}_1\tilde{t}_1 (\text{medium}), \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$ | $2 e, \mu$ | 2 jets | Yes | 20.3 | \tilde{t}_1 215-530 GeV | $m(\tilde{\chi}_1^0)=1 \text{ GeV}$ | |
| | $\tilde{t}_1\tilde{t}_1 (\text{medium}), \tilde{t}_1 \rightarrow b\tilde{\chi}_1^\pm$ | 0 | $2 b$ | Yes | 20.1 | \tilde{t}_1 150-580 GeV | $m(\tilde{\chi}_1^0)<200 \text{ GeV}, m(\tilde{\chi}_1^\pm)-m(\tilde{\chi}_1^0)=5 \text{ GeV}$ | |
| | $\tilde{t}_1\tilde{t}_1 (\text{heavy}), \tilde{t}_1 \rightarrow t\tilde{\chi}_1^\pm$ | $1 e, \mu$ | $1 b$ | Yes | 20 | \tilde{t}_1 210-640 GeV | $m(\tilde{\chi}_1^0)=0 \text{ GeV}$ | |
| | $\tilde{t}_1\tilde{t}_1 (\text{heavy}), \tilde{t}_1 \rightarrow t\tilde{\chi}_1^0$ | 0 | $2 b$ | Yes | 20.1 | \tilde{t}_1 260-640 GeV | $m(\tilde{\chi}_1^0)=0 \text{ GeV}$ | |
| | $\tilde{t}_1\tilde{t}_1, \tilde{t}_1 \rightarrow c\tilde{\chi}_1^0$ | 0 | mono-jet/c-tag | Yes | 20.3 | \tilde{t}_1 90-240 GeV | $m(\tilde{t}_1)-m(\tilde{\chi}_1^0)<85 \text{ GeV}$ | |
| | $\tilde{t}_1\tilde{t}_1 (\text{natural GMSB})$ | $2 e, \mu (Z)$ | $1 b$ | Yes | 20.3 | \tilde{t}_1 150-580 GeV | $m(\tilde{\chi}_1^0)>150 \text{ GeV}$ | |
| | $\tilde{t}_2\tilde{t}_2, \tilde{t}_2 \rightarrow \tilde{t}_1 + Z$ | $3 e, \mu (Z)$ | $1 b$ | Yes | 20.3 | \tilde{t}_2 290-600 GeV | $m(\tilde{\chi}_1^0)<200 \text{ GeV}$ | |
| | EW direct | $\tilde{\tau}_{1,R}\tilde{\tau}_{1,L,R}, \tilde{\tau} \rightarrow \ell\tilde{\chi}_1^0$ | $2 e, \mu$ | 0 | Yes | 20.3 | $\tilde{\tau}$ 90-325 GeV | $m(\tilde{\chi}_1^0)=0 \text{ GeV}$ |
| | | $\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow \ell\nu(\ell\bar{\nu})$ | $2 e, \mu$ | 0 | Yes | 20.3 | $\tilde{\chi}_1^\pm$ 140-465 GeV | $m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$ |
| 2τ | | - | - | Yes | 20.3 | $\tilde{\chi}_1^\pm$ 100-350 GeV | $m(\tilde{\chi}_1^0)=0 \text{ GeV}, m(\tilde{\tau}, \tilde{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$ | |
| $\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \rightarrow \tilde{\ell}_1\nu\tilde{\ell}_1\ell(\tilde{\nu}\nu), \ell\tilde{\nu}\tilde{\ell}_1\ell(\tilde{\nu}\nu)$ | | $3 e, \mu$ | 0 | Yes | 20.3 | $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ 700 GeV | $m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)=0, m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{\chi}_1^\pm)+m(\tilde{\chi}_1^0))$ | |
| $\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0 Z$ | | $2-3 e, \mu$ | 0 | Yes | 20.3 | $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ 420 GeV | $m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)=0, \text{ sleptons decoupled}$ | |
| $\tilde{\chi}_1^\pm\tilde{\chi}_2^0 \rightarrow W\tilde{\chi}_1^0 h\tilde{\chi}_1^0$ | | $1 e, \mu$ | $2 b$ | Yes | 20.3 | $\tilde{\chi}_1^\pm, \tilde{\chi}_2^0$ 285 GeV | $m(\tilde{\chi}_1^\pm)=m(\tilde{\chi}_2^0), m(\tilde{\chi}_1^0)=0, \text{ sleptons decoupled}$ | |
| $\tilde{\chi}_2^0\tilde{\chi}_3^0, \tilde{\chi}_{2,3}^0 \rightarrow \tilde{\ell}_R\ell$ | | $4 e, \mu$ | 0 | Yes | 20.3 | $\tilde{\chi}_{2,3}^0$ 620 GeV | $m(\tilde{\chi}_2^0)=m(\tilde{\chi}_3^0), m(\tilde{\chi}_1^0)=0, m(\tilde{\ell}, \tilde{\nu})=0.5(m(\tilde{\chi}_2^0)+m(\tilde{\chi}_1^0))$ | |
| Long-lived particles | Direct $\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm$ prod., long-lived $\tilde{\chi}_1^\pm$ | Disapp. trk | 1 jet | Yes | 20.3 | $\tilde{\chi}_1^\pm$ 270 GeV | $m(\tilde{\chi}_1^\pm)-m(\tilde{\chi}_1^0)=160 \text{ MeV}, \tau(\tilde{\chi}_1^\pm)=0.2 \text{ ns}$ | |
| | Stable, stopped \tilde{g} R-hadron | 0 | 1-5 jets | Yes | 27.9 | \tilde{g} 832 GeV | $m(\tilde{\chi}_1^0)=100 \text{ GeV}, 10 \mu\text{s}<\tau(\tilde{g})<1000 \text{ s}$ | |
| | GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu})+\tau(e, \mu)$ | $1-2 \mu$ | - | - | 15.9 | $\tilde{\chi}_1^0$ 475 GeV | $10<\tan\beta<50$ | |
| | GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma\tilde{G}$, long-lived $\tilde{\chi}_1^0$ | 2γ | - | Yes | 4.7 | $\tilde{\chi}_1^0$ 230 GeV | $0.4<\tau(\tilde{\chi}_1^0)<2 \text{ ns}$ | |
| | $\tilde{q}\tilde{q}, \tilde{\chi}_1^0 \rightarrow qq\mu$ (RPV) | $1 \mu, \text{ displ. vtx}$ | - | - | 20.3 | \tilde{q} 1.0 TeV | $1.5<c\tau<156 \text{ mm}, \text{BR}(\mu)=1, m(\tilde{\chi}_1^0)=108 \text{ GeV}$ | |
| RPV | LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e + \mu$ | $2 e, \mu$ | - | - | 4.6 | $\tilde{\nu}_\tau$ 1.61 TeV | $\lambda'_{311}=0.10, \lambda_{132}=0.05$ | |
| | LFV $pp \rightarrow \tilde{\nu}_\tau + X, \tilde{\nu}_\tau \rightarrow e(\mu) + \tau$ | $1 e, \mu + \tau$ | - | - | 4.6 | $\tilde{\nu}_\tau$ 1.1 TeV | $\lambda'_{311}=0.10, \lambda_{1(2)33}=0.05$ | |
| | Bilinear RPV CMSSM | $2 e, \mu (SS)$ | 0-3 b | Yes | 20.3 | \tilde{q}, \tilde{g} 1.35 TeV | $m(\tilde{q})=m(\tilde{g}), c\tau_{LSP}<0.1 \text{ mm}$ | |
| | $\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow ee\tilde{\nu}_\mu, e\mu\tilde{\nu}_e$ | $4 e, \mu$ | - | Yes | 20.3 | $\tilde{\chi}_1^\pm$ 750 GeV | $m(\tilde{\chi}_1^0)>0.2 \times m(\tilde{\chi}_1^\pm), \lambda_{121} \neq 0$ | |
| | $\tilde{\chi}_1^\pm\tilde{\chi}_1^\pm, \tilde{\chi}_1^\pm \rightarrow W\tilde{\chi}_1^0, \tilde{\chi}_1^0 \rightarrow \tau\tau\tilde{\nu}_e, e\tau\tilde{\nu}_\tau$ | $3 e, \mu + \tau$ | - | Yes | 20.3 | $\tilde{\chi}_1^\pm$ 450 GeV | $m(\tilde{\chi}_1^0)>0.2 \times m(\tilde{\chi}_1^\pm), \lambda_{133} \neq 0$ | |
| | $\tilde{g} \rightarrow qq\tilde{q}$ | 0 | 6-7 jets | - | 20.3 | \tilde{g} 916 GeV | $\text{BR}(\ell)=\text{BR}(b)=\text{BR}(c)=0\%$ | |
| | $\tilde{g} \rightarrow \tilde{t}_1 t, \tilde{t}_1 \rightarrow bs$ | $2 e, \mu (SS)$ | 0-3 b | Yes | 20.3 | \tilde{g} 850 GeV | | |
| Other | Scalar gluon pair, $\text{sgluon} \rightarrow q\tilde{q}$ | 0 | 4 jets | - | 4.6 | sgluon 100-287 GeV | incl. limit from 1110.2693 | |
| | Scalar gluon pair, $\text{sgluon} \rightarrow t\tilde{t}$ | $2 e, \mu (SS)$ | $2 b$ | Yes | 14.3 | sgluon 350-800 GeV | | |
| | WIMP interaction (D5, Dirac χ) | 0 | mono-jet | Yes | 10.5 | M^* scale 704 GeV | $m(\chi)<80 \text{ GeV}$, limit of <687 GeV for D8 | |

$\sqrt{s} = 7 \text{ TeV}$ full data
 $\sqrt{s} = 8 \text{ TeV}$ partial data
 $\sqrt{s} = 8 \text{ TeV}$ full data

10⁻¹ 1 Mass scale [TeV]

*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

Strong limits on SUSY partners



Tension with Higgs mass:

$$M_h^2 \leq M_Z^2 \cos^2(2\beta) + \epsilon$$

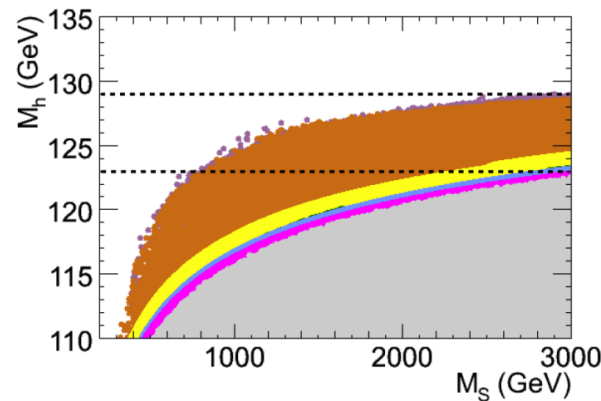
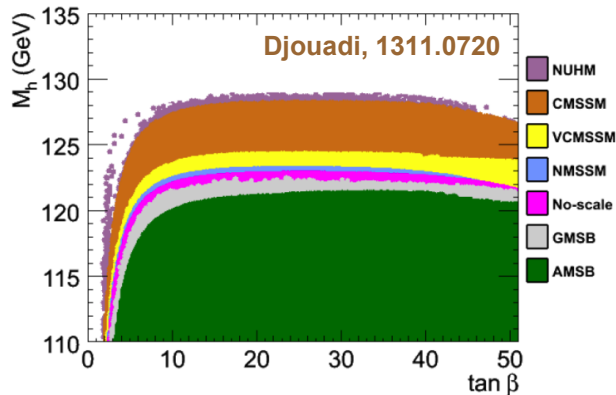
$$\epsilon \approx \frac{3 m_t^4}{2\pi^2 v^2 \sin^2(\beta)} \left[\log \left(\frac{M_S^2}{m_t^2} \right) + \frac{X_t^2}{M_S^2} \left(1 - \frac{X_t^2}{12M_S^2} \right) \right]$$

$$M_S^2 = m_{\tilde{t}_1} m_{\tilde{t}_2}$$

Decoupling ($M_A \gg M_Z$)

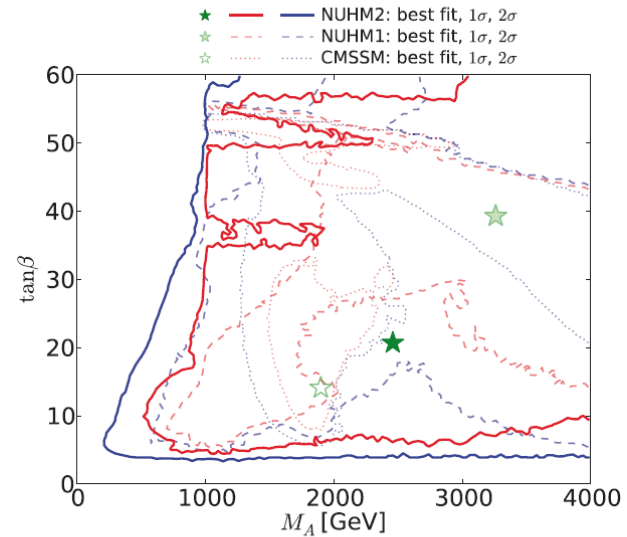
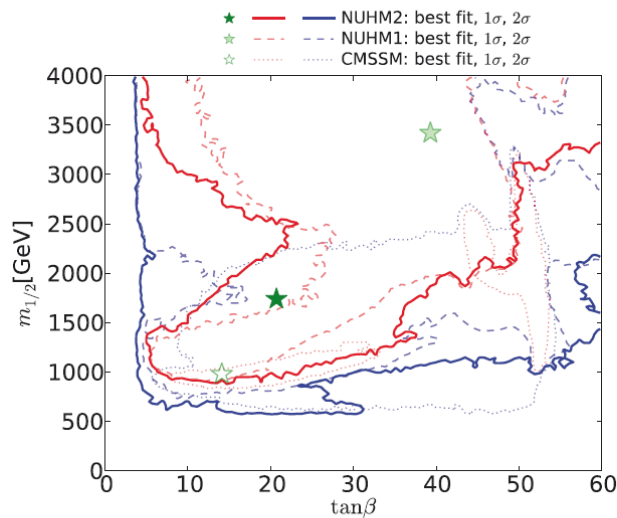
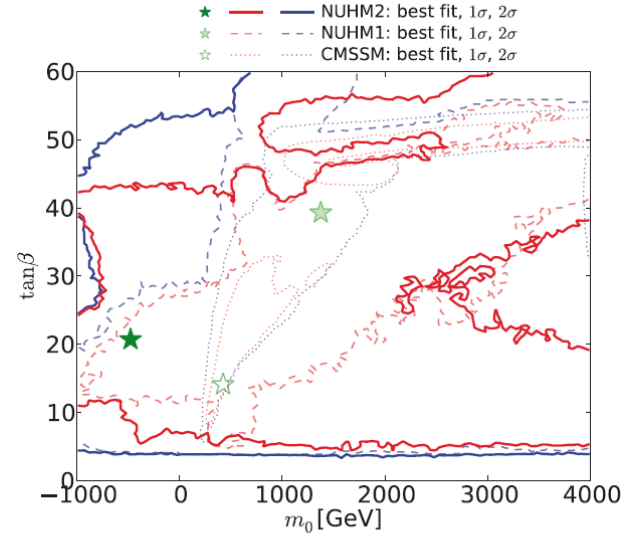
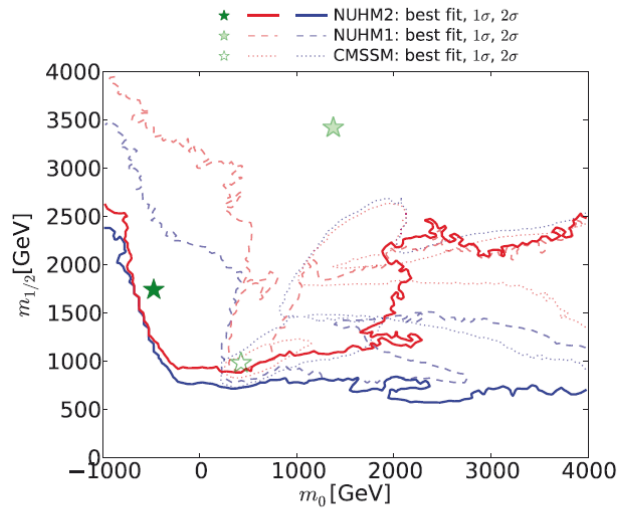
$\cos^2(2\beta) \rightarrow 1$

Maximal stop mixing X_t



Improved higher-order calculations allow slightly larger values of M_h

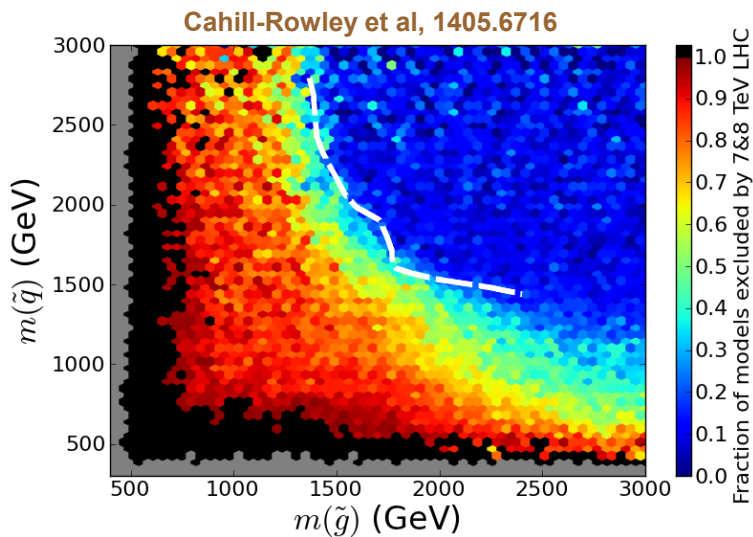
Heinemeyer et al



$(g - 2)_\mu$ cannot be explained (not included in the fit)

Which SUSY?

- **Looks bad in CMSSM** (120 MSSM parameters reduced to 4 + 1 sign)
- **More freedom in the Phenomenological MSSM**



120 MSSM parameters reduced to 19-20

Many “models” consistent with data

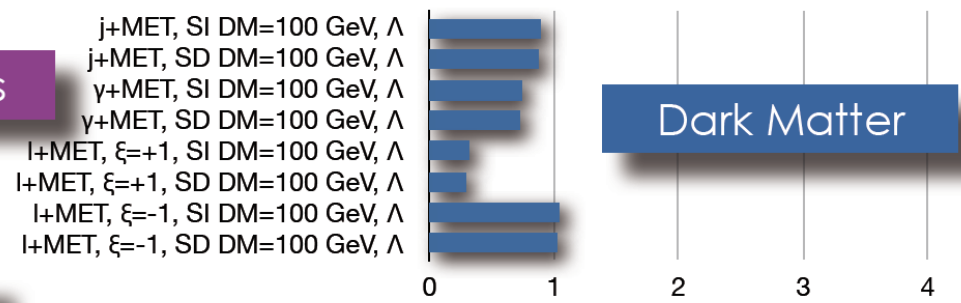
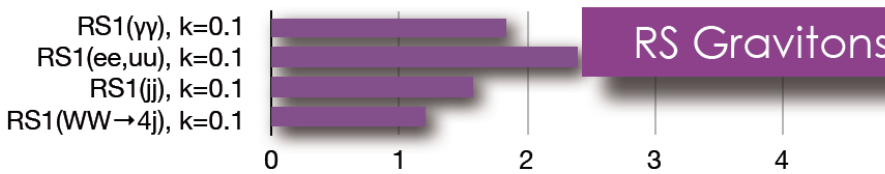
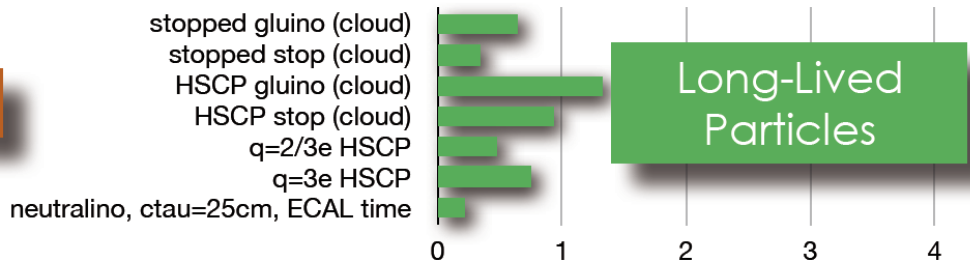
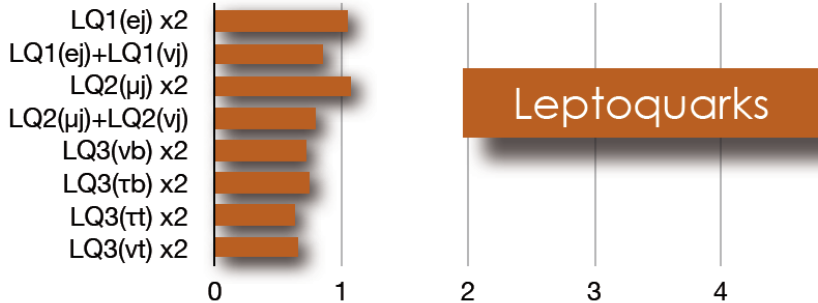
Data-driven search

- **Many SUSY variants:** NMSSM, Split, High-Scale, Stealth, 5D, Natural, Folded, Twin ...

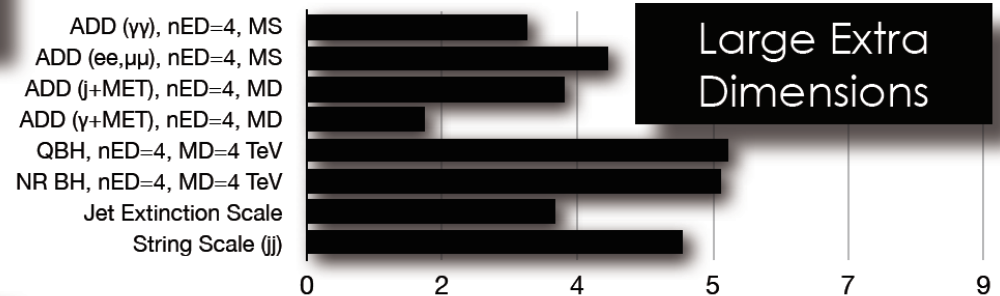
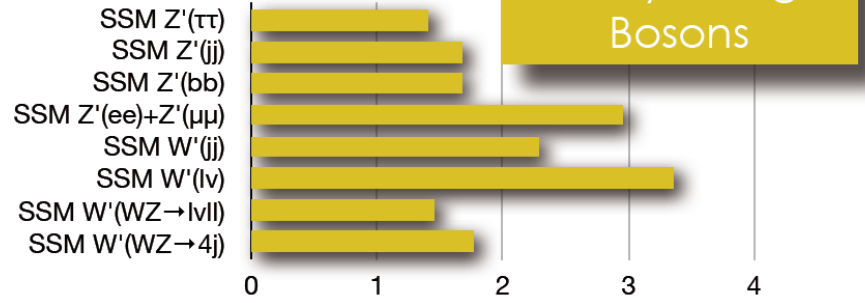
Naturalness?

$$\Delta M_h^2 \propto M_{\text{SUSY}}^2$$

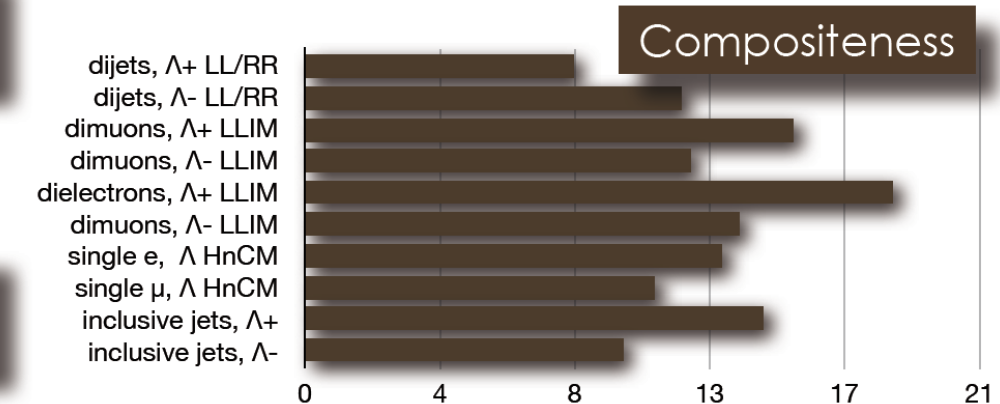
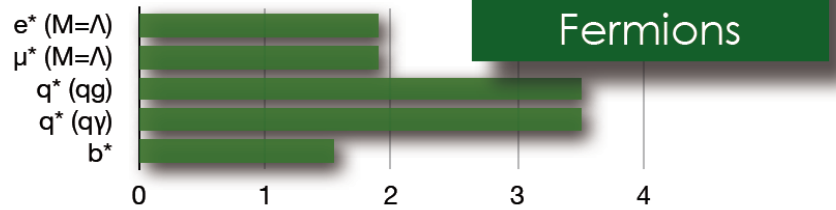
CMS Preliminary



Heavy Gauge Bosons



Excited Fermions



Multijet Resonances





Massive & dark SUSY states show up through a hidden portal from a warped dimension

Look, your worship, it's just the spectrum of the Standard Model

Effective Field Theory

$$L_{\text{eff}} = L^{(4)} + \sum_{D>4} \sum_k \frac{c_k^{(D)}}{\Lambda_{\text{NP}}^{D-4}} O_k^{(D)}$$

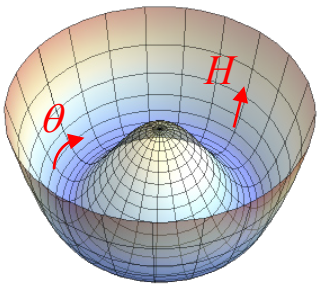
- Most general Lagrangian with the SM gauge symmetries
- Light ($m \ll \Lambda_{\text{NP}}$) fields only
- The SM Lagrangian corresponds to $D=4$
- $c_k^{(D)}$ contain information on the underlying dynamics:

$$L_{\text{NP}} \doteq g_X (\bar{q}_L \gamma^\mu q_L) X_\mu \quad \rightarrow \quad \frac{g_X^2}{M_X^2} (\bar{q}_L \gamma^\mu q_L) (\bar{q}_L \gamma_\mu q_L)$$

- Options for **H(126)**:
 - $\text{SU}(2)_L$ doublet (SM)
 - Scalar singlet
 - Additional light scalars

Custodial Symmetry

$$\Sigma \equiv (\phi_C, \phi) = \begin{pmatrix} \phi^{(0)*} & \phi^{(+)} \\ -\phi^{(-)} & \phi^{(0)} \end{pmatrix} = \frac{1}{\sqrt{2}} [v + H(x)] U(\varphi) \quad , \quad U(\varphi) = \exp \left\{ \frac{i}{v} \vec{\sigma} \vec{\varphi} \right\}$$



$$\begin{aligned} \mathcal{L}(\phi) &= (D_\mu \phi)^\dagger D^\mu \phi - \lambda \left(\phi^\dagger \phi - \frac{v^2}{2} \right)^2 \\ &= \frac{1}{2} \text{Tr} \left[(D^\mu \Sigma)^\dagger D^\mu \Sigma \right] - \frac{\lambda}{4} \left(\text{Tr} [\Sigma^\dagger \Sigma] - v^2 \right)^2 \\ &= \frac{v^2}{4} \text{Tr} \left[(D^\mu U)^\dagger D^\mu U \right] + O(H/v) \end{aligned}$$

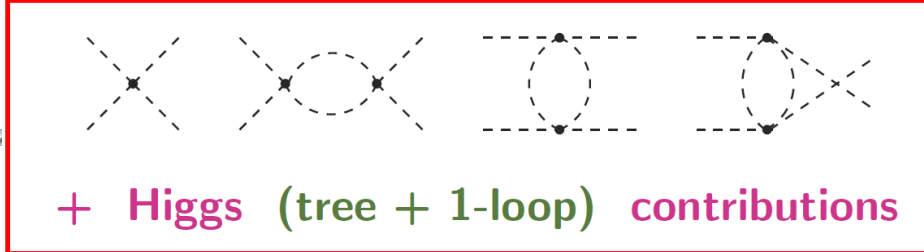
- Invariant under global $SU(2)_L \otimes SU(2)_R \supset SU(2)_L \otimes U(1)_Y$

$$\Sigma \rightarrow g_L \cdot \Sigma \cdot g_R^\dagger \quad , \quad g_X \in SU(2)_X$$

- Same Lagrangian than QCD pions: $f_\pi \rightarrow v$, $\pi^\pm, \pi^0 \rightarrow \varphi^\pm, \varphi^0 \rightarrow W_L^\pm, W_L^0$

Chiral Goldstone Bosons: $SU(2)_L \otimes SU(2)_R \rightarrow SU(2)_C$

$$W_L W_L \rightarrow W_L W_L$$



$$\mathcal{L} = \frac{v^2}{4} \langle D^\mu U^\dagger D_\mu U \rangle \left[1 + 2a \frac{H}{v} + b \frac{H^2}{v^2} \right]$$

(Espriu–Mescia–Yencho, Delgado–Dobado–Llanes-Estrada)

$$A(s, t, u) = \frac{s}{v^2} (1 - a^2) + \frac{4}{v^4} \left[a_4^r(\mu) (t^2 + u^2) + 2 a_5^r(\mu) s^2 \right] \\ + \frac{1}{16\pi^2 v^4} \left\{ \frac{1}{9} (14a^4 - 10a^2 - 18a^2b + 9b^2 + 5) s^2 + \frac{13}{18} (1 - a^2)^2 (t^2 + u^2) \right. \\ - \frac{1}{2} (2a^4 - 2a^2 - 2a^2b + b^2 + 1) s^2 \log\left(\frac{-s}{\mu^2}\right) \\ \left. + \frac{1}{12} (1 - a^2)^2 \left[(s^2 - 3t^2 - u^2) \log\left(\frac{-t}{\mu^2}\right) + (s^2 - t^2 - 3u^2) \log\left(\frac{-u}{\mu^2}\right) \right] \right\}$$

SM: $a = b = 1$, $a_4 = a_5 = 0$



$$A(s, t, u) \sim \mathcal{O}(M_H^2/v^2)$$

$$A(\gamma\gamma \rightarrow W_L^+ W_L^-)_{\text{NLO}} \sim \frac{(a^2 - 1)}{8\pi^2 v^2}$$

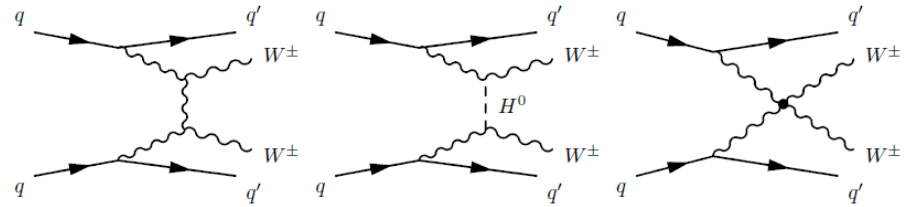
Delgado et al

Deviations of the SM gauge couplings imply bad UV behaviours

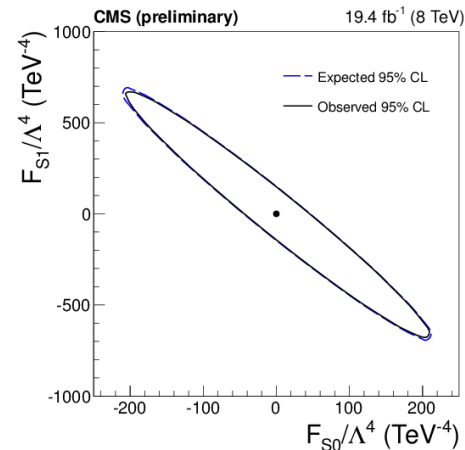
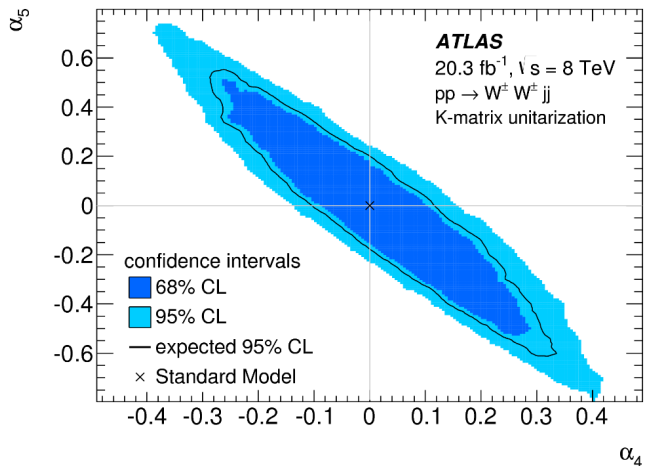
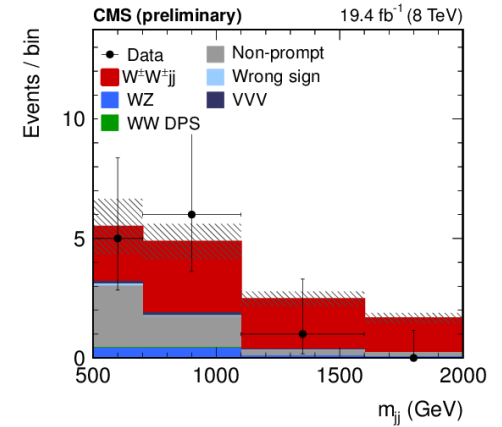
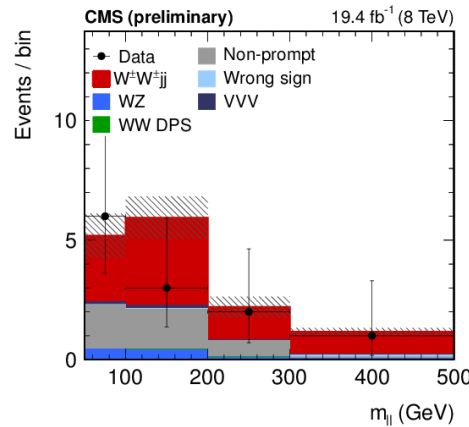
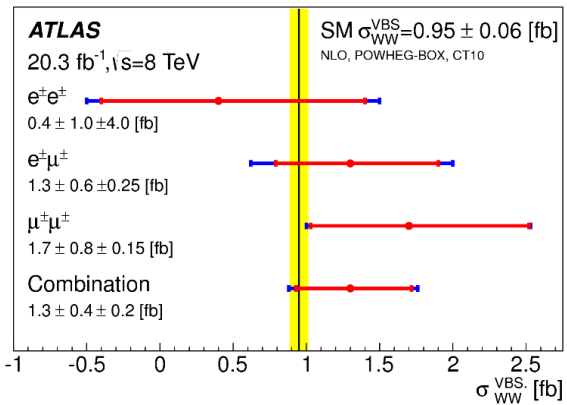
New states needed to restore unitarity

WW Scattering @ LHC

First evidence of $W^\pm W^\pm$ scattering (3.6σ)



ATLAS, arxiv:1405.6241



CMS-PAS-SMP-13-015



Strongly-Coupled Scenarios

- **Symmetry Breaking:** $SU(2)_L \otimes SU(2)_R \rightarrow SU(2)_C$
- **Goldstone Dynamics** \rightarrow **Electroweak Effective Theory**
- **Strong Electroweak Dynamics** \rightarrow **Heavy Resonances**
- **Many possibilities:** (Walking, Conformal) Technicolour, CFT, 5D ...
- **Light Scalar Resonance H(125)**
 \rightarrow **Pseudo-Goldstone (composite) Higgs, Dilaton ...**

The power of the dark side

Holds the Universe together and makes *85% of all the matter in it!*

Interacts very weakly
(not charged)

- Gravity ✓
- Higgs-like Interactions ?

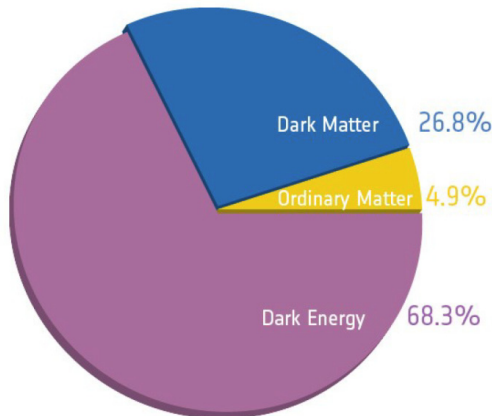


SUSY and the WIMP Miracle ?

- If the LSP is the lightest neutralino it will behave as WIMP dark matter
- In the MSSM the lightest neutralino is generically a mixture of the Bino, Wino, and the two Higgsinos
- If you are more ambitious, can try to require that the LSP is a thermal relic with the correct abundance to explain all ALL dark matter

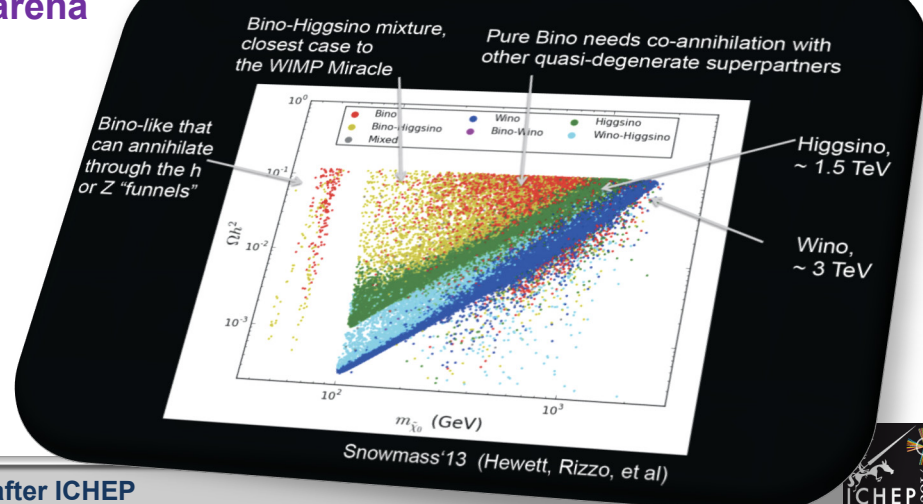


Dark Matter



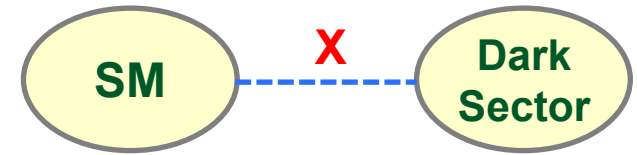
M. Carena

SUSY and the WIMP "Miracle"



Hidden Portals

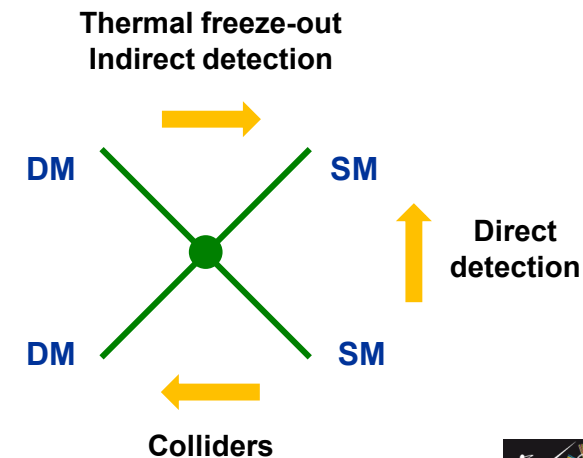
Coupling to a hidden Dark Sector through new SM-singlet particles



- **Higgs Portal:** $\chi H^\dagger H$, $\chi^2 H^\dagger H$
- **Vector Portal:** $V_{\mu\nu} F^{\mu\nu}$
- **Neutrino Portal:** $\bar{L}_L H N_R$
- **Axion Portal:** $a \tilde{G}_{\mu\nu} G^{\mu\nu}$, $\partial^\mu a \bar{\psi} \gamma_\mu \gamma_5 \psi$

DM candidates in many BSMs

Complementary experimental information



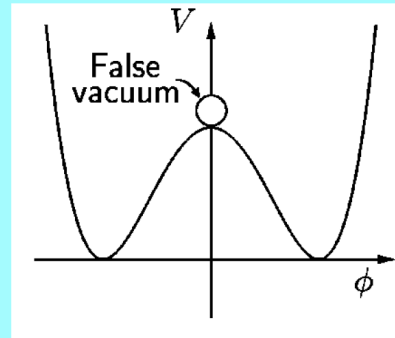
Inflation Paradigm

Unsolved issues in the standard model

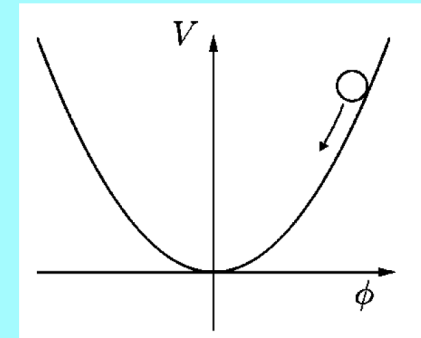
- Horizon problem
Why is the CMB so smooth?
- The flatness problem
Why is the Universe flat? Why is $\Omega \sim 1$?
- The structure problem
Where do the fluctuations in the CMB come from?
- The relic problem
Why aren't there magnetic monopoles?

★ **Negative Pressure** \implies **Repulsive Gravity.** Guth

★ State dominated by scalar field potential energy \implies Negative Pressure.



New (Small Field) Inflation
Linde; Albrecht & Steinhardt (1982)



Chaotic (Large Field) Inflation
Linde (1983)

Planck



$$\Omega = 1.0010 \pm 0.0065$$

BICEP2



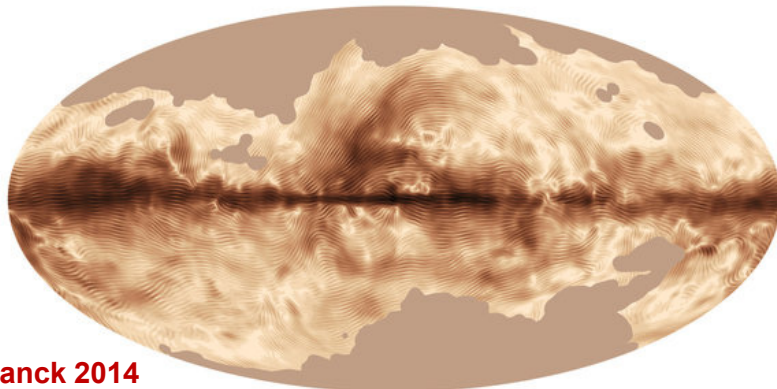
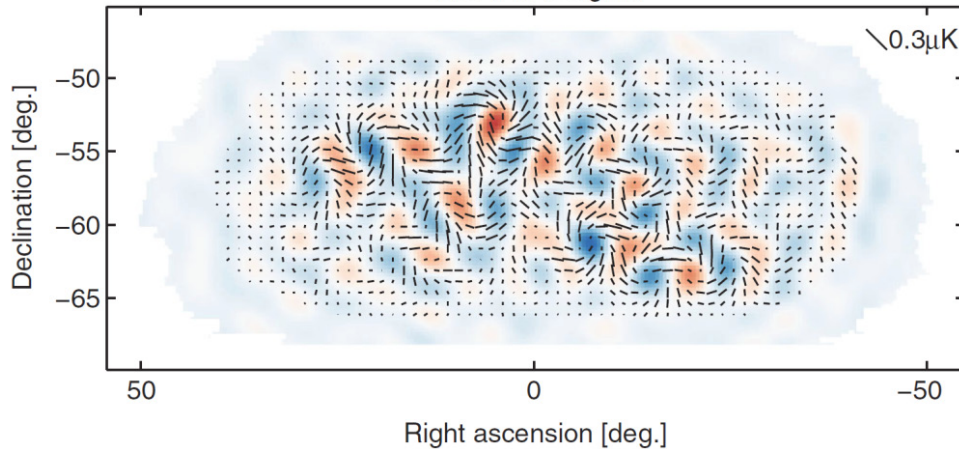
$$\rho_{\text{inf}}^{1/4} = 2.2 \times 10^{16} \text{ GeV} \left(\frac{r}{0.2} \right)^{1/4}$$

BICEP2 Data & Inflation Paradigm

$$r = 0.020^{+0.07}_{-0.05} \quad (\text{tensor / scalar})$$

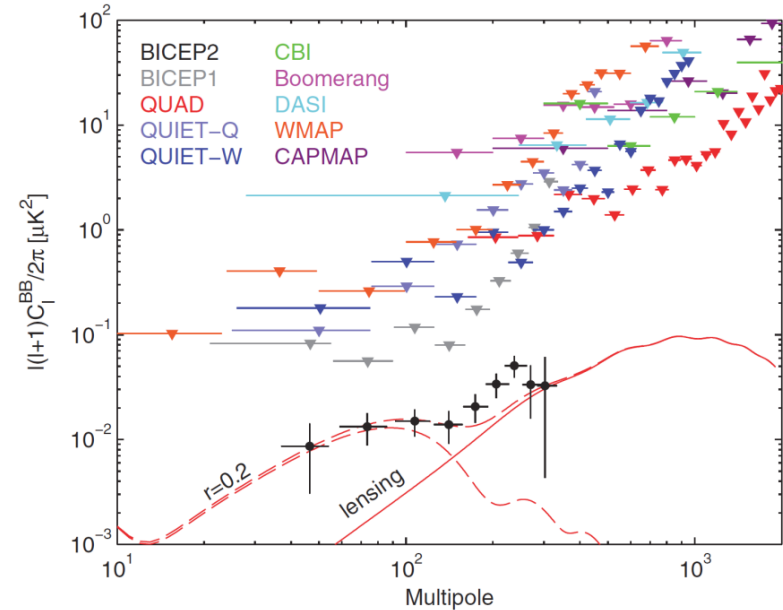
PRL 112, 241101 (2014)

BICEP2: B signal



Planck 2014

Milky Way's (dust) magnetic fingerprint



- Evidence of inflationary gravitational waves?
- Foreground polarized dust emission?

Flauger et al 1405.7351, Mortonson-Seljak 1405.5857

Inflaton

□ Another scalar field?

□ Could “Higgs Inflation” work?

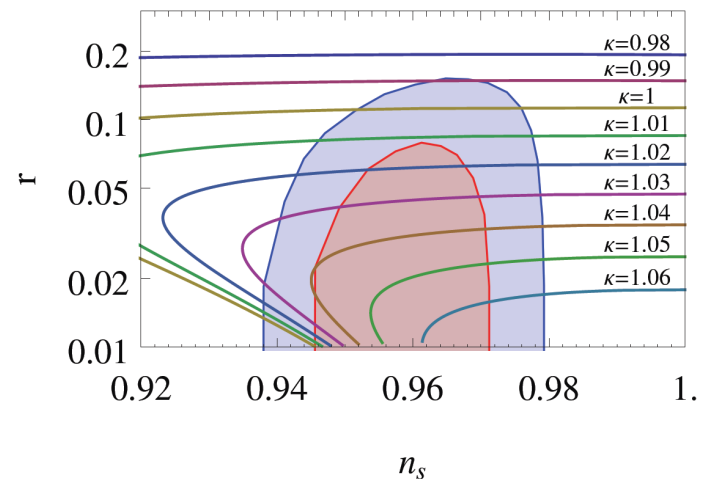
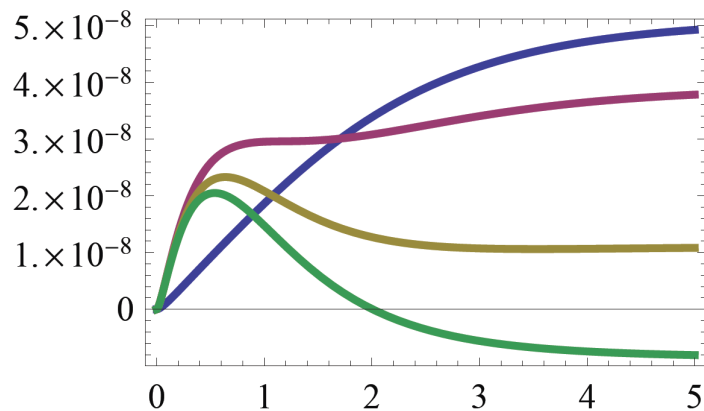
$$S_G = -\frac{1}{2} \int d^4x \sqrt{-g} \{M_{\text{Pl}}^2 R + \xi H^2 R\}$$

$$\xi \approx 47000 \sqrt{\lambda} \quad (\text{COBE}) \quad \rightarrow \quad n_s \approx 0.97 \quad , \quad r \approx 0.003$$

Quantum effects: $M_H > M_{\text{crit}} \sim 129.6 \text{ GeV}$

Bezrukov-Shaposhnikov, 1403.6078

Close to M_{crit} , n and r strongly depend on M_H and m_t



Status & Outlook

- The **SM** appears to be the right theory at the EW scale
- The **H(125)** behaves as the SM scalar boson
- The **CKM** mechanism works very well
- Neutrinos do have (**tiny**) masses. Lepton flavour is violated
- Different **flavour structure** for quarks & leptons
- **New physics needed** to explain many pending questions:
Flavour, CP, baryogenesis, dark matter, cosmology...



- **How far is the Scale of New-Physics Λ_{NP} ?**
- **Which symmetry keeps M_{H} away from Λ_{NP} ?**
Supersymmetry, scale/conformal symmetry...
- **Which kind of New Physics?**

Awaiting great discoveries @ LHC

This, no doubt, Sancho, will be a most mighty and perilous adventure, in which it will be needful for me to put forth all my valour and resolution

Let your worship be calm, senior. Maybe it's all enchantment, like the phantoms last night

